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TECHNICAL REPORT RK-CR-84-3

A COMPUTER PROGRAM FOR THE PERFORMANCE ANALYSIS OF SCARFED NOZZLES

Prepared by:

Joe D. Hoffman
Battelle Columbus Research Laboratories
Research Triangle Park, NC 27709

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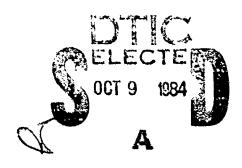
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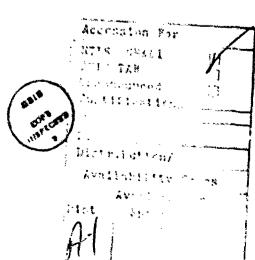
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An analysis is presented for predicting the performance of scarfed propulsive nozzles. The model assumes that the scarfing is not severe, so that the flow within the nozzle is axisymmetric. The flowfield within the nozzle and the scarfed nozzle extension is calculated by the method of characteristics. The oblique shock wave emaninating from the junction of the nozzle and the scarfed nozzle extension is fitted discretely and tracked through the flowfield. The nozzle forces are determined by integrating the wall pressure distribution.										

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SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered) Nozzle axial and side forces are calculated, and missile axial and side forces are determined. A computer program implementing the analysis is discussed. Fifteen sample cases are presented to illustrate the use of the computer

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SECTION 1

INTRODUCTION

The conventional propulsive nozzle is an axisymmetric converging-diverging nozzle terminating in a plane perpendicular to the axis of the nozzle. If the flow entering the nozzle is axisymmetric, then the flow within the nozzle is also axisymmetric, and the resultant thrust vector lies along the nozzle axis.

In some applications, such as a nozzle exiting through the side of a missile where the nozzle axis is not aligned with the missile axis, the nozzle terminates along the line of intersection of the nozzle contour and the missile outer skin. Such a nozzle is called a scarfed nozzle. Figure I illustrates several possible scarfed nozzle configurations.

Unbalanced side forces are generated in scarfed nozzles, and the resultant thrust vector does not lie along the nozzle axis or the missile axis. When two or more identical scarfed mozzles are located symmetrically around the missile axis, the resultant thrust vector does lie along the missile axis. However, even in that case the missile axial thrust is reduced below the nozzle axial thrust due to the misalignment of the nozzle axis and the missile axis.

Expansion waves or shock waves, depending on the pressure ratio, emanate from the nozzle exit lip contour as the internal flow leaves the protection of the splid wall and flows into the surrounding

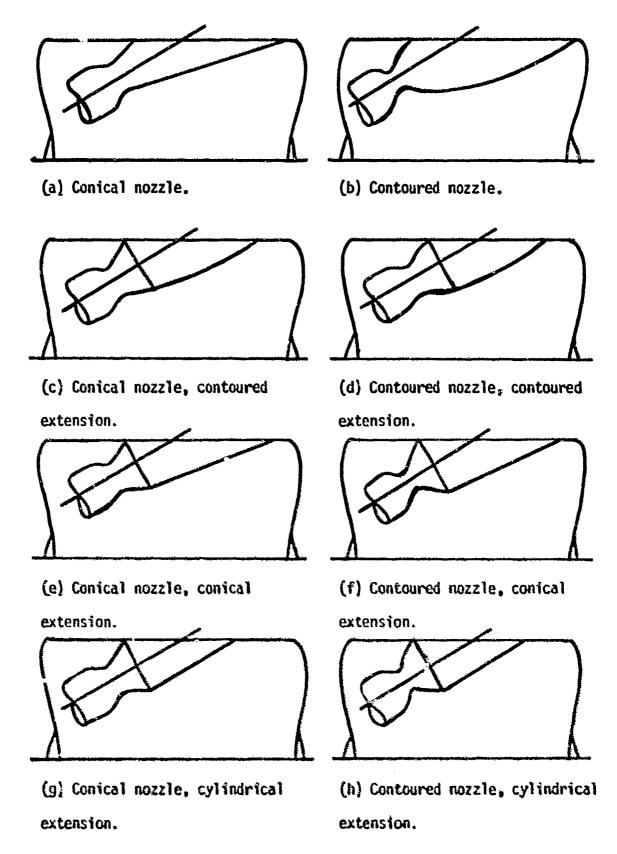


Figure 1. Several possible scarfed nozzle configurations.

atmosphere. When scarfing is not too severe, the waves emanating from the top of the nozzle at the initial point of scarfing do not intersect the opposite wall, and the portion of the flowfield affecting the pressure on the nozzle wall remains axisymmetric. That is the case considered in the present analysis. When scarfing is severe, the aforementioned waves intersect the opposite wall, and the flow downstream of that intersection is fully three-dimensional. That case is not considered in the present analysis.

The present report presents an analysis, and a computer program for implementing that analysis, for determining the performance of scarfed axisymmetric propulsive nozzles when the scarfing is small enough so that the portion of the flowfield affecting the pressure on the nozzle wall remains axisymmetric.

SECTION II

NOZZLE GEOMETRIC MODEL

1. INTRODUCTION

Several scarfed nozzle configurations are illustrated in Figure 1. Six specific nozzle configurations are considered in the present study. These configurations are illustrated in Figure 2.

Figure 2(a) illustrates a scarfed conical nozzle, where the nozzle contour is specified analytically as a straight line.

$$y = a + bx (1)$$

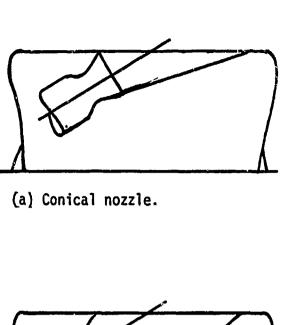
Figure 2(b) illustrates a conical nozzle, specified by equation (1), followed by a scarfed conical extension starting at the end of the conical nozzle, where the nozzle extension is specified analytically as a straight line.

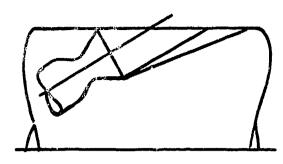
$$y = e + fx \tag{2}$$

When the slope f is zero, a cylindrical extension is obtained.

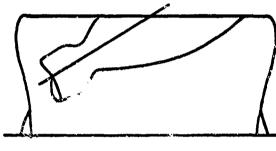
Figure 2(c) illustrates a scarfed quadratic nozzle, where the nozzle contour is specified analytically as a second-order (i.e., quadratic) polynomial.

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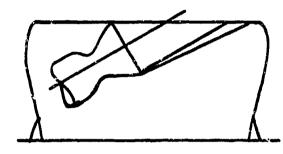




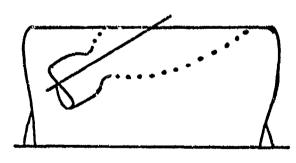
(b) Conical nozzle followed by a conical (or cylindrical) extension.



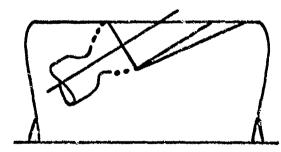
(c) Quadratic nozzle.



(d) Quadratic nozzle followed by a conical (or cylindrical) extension.



(e) Tabular nozzle.



(f) Tabular nozzle followed by a conical (or cylindrical) extension.

Figure 2. Nozzle geometric models considered.

$$y = a + bx + cx^2 \tag{3}$$

Figure 2(d) illustrates a quadratic nozzle, specified by equation (3), followed by a scarfed conical extension starting at the end of the quadratic nozzle, where the nozzle extension is specified by equation (2).

Figure 2(e) illustrates a scarfed tabular nozzle, where the basic nozzle contour is specified in tabular form.

$$y_i = f(x_i), (i = 1,..., n) tabular$$
 (4)

Figure 2(f) illustrates a tabular nozzle, specified by equation (4), followed by a scarfed conical extension starting at the end of the tabular nozzle, where the nozzle extension is specified by equation (2).

Figure 3 presents a meridional plane view of the geometric model employed for the scarfed nozzle flowfield analysis considered in the present investigation. The nozzle consists of a conventional axisymmetric throat and supersonic expansion contour up to point E where scarfing begins. That portion of the scarfed nozzle is called the basic nozzle. The scarfed section from point E to point F is a conical extension to the basic nozzle. That portion of the scarfed nozzle is called the nozzle extension. In the present analysis, the scarfed section always starts at point E.

The specification of the scarfed nozzle geometry is discussed in in the following paragraphs.

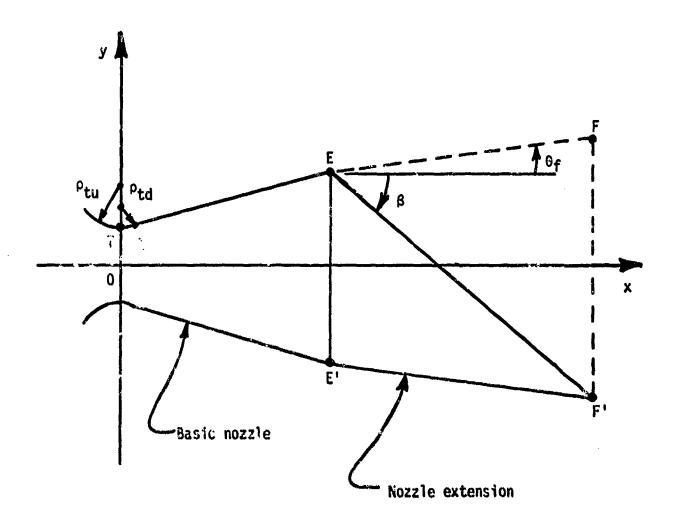


Figure 3. Scarfed nozzle geometric model.

2. BASIC NOZZLE GEOMETRY

The basic nozzle geometry consists of a double circular arc throat contour joined smoothly to a supersonic expansion contour. The supersonic expansion contour can be conical, quadratic, or tabular. The specification of the basic nozzle geometry is discussed in the following paragraphs.

Throat geometry. The throat of the Lasic nozzle consists of a double circular arc contour that attaches smoothly to the supersonic expansion contour at point A. The throat geometry is completely specified by the throat radius y_t , the upstream throat radius of curvature ρ_{td} , the downstream throat radius of curvature ρ_{td} , and the throat attachment angle θ_a where the supersonic expansion contour attaches smoothly to the circular arc throat contour. The location of point A is determined from the following two equations.

$$x_a = \rho_{td} \sin \theta_a$$
 (5)

$$y_a = y_t + \rho_{td} (1 - \cos \theta_a)$$
 (6)

The upstream throat radius of curvature ρ_{td} is used in the calculation of the supersonic initial-value line which spans the nozzle throat from point T to the point on the nozzle centerline where the Mach number is 1.0.

Conical nozzle. When the supersonic expansion contour is conical, the cone angle α must equal the throat attachment angle θ_a . Five

options exist for specifying the conical nozzle contour.

1. Specified throat attachment angle θ_a and nozzle length x_e . In this case, the nozzle exit lip radius y_e is determined from

$$y_e = y_a + (x_e - x_a) \tan \theta_a \tag{7}$$

2. Specified throat attachment angle θ_a and exit lip radius y_e . In this case, the nozzle length x_e is determined from

$$x_e = x_a + (y_e - y_a) / \tan \theta_a$$
 (8)

3. Specified throat attachment angle θ_a and nozzle area ratio ϵ = $(y_e/y_t)^2$. In this case, the nozzle exit lip radius y_e is determined from

$$y_e = y_t \sqrt{\varepsilon}$$
 (9)

The nozzle length x_e is then determined from equation (8).

4. Specified nozzle length \mathbf{x}_e and exit lip radius $\mathbf{y}_e.$ The throat attachment angle θ_a is determined from

$$\theta_a = \tan^{-1} \left(\frac{y_e - y_a}{x_e - x_a} \right) \tag{10}$$

5. Specified nozzle length x_e and nozzle area ratio ϵ . The nozzle exit lip radius y_e is determined from equation (9) and the throat attachment angle θ_a is determined from equation (10).

In all five cases, point A (x_a, y_a) , point E (x_e, y_e) , and the throat attachment angle θ_a are known. The equation for the supersonic expansion contour is given by

$$y(x) = a + bx \tag{11}$$

where

$$b = \frac{dy(x)}{dx} = \tan \alpha$$
 and $a = y_a - bx_a$ (12)

Quadratic nozzle. When the supersonic expansion contour is quadratic, the supersonic contour must attach smoothly to the circular arc throat contour, and the nozzle length x_e must be specified. The equation for the supersonic expansion contour is

$$y(x) = a + bx + cx^2 \tag{13}$$

Three options exist for specifying the supersonic expansion contour.

1. Specified throat attachment angle θ_a , nozzle exit lip angle θ_e , and nozzle length x_e . Substituting the three known values (x_a,y_a) , (x_a,θ_a) , and (x_e,θ_e) into equation (13) and its derivative (i.e., dy/dx = b + 2cx = tan θ) yields

$$c = \frac{\tan \theta_e - \tan \theta_a}{2(x_e - x_a)} \tag{14}$$

$$b = \tan \theta_a - 2cx_a \tag{15}$$

$$a = y_a - bx_a - cx_a^2$$
 (16)

2. Specified throat attachment angle θ_a , nozzle length x_e , and exit lip radius y_e . Substituting the three known values (x_a,y_a) , (x_a,θ_a) , and (x_e,y_e) into equation (13) yields

$$c = \frac{\left(\frac{y_e - y_a}{x_e - x_a}\right) - \tan \theta_a}{x_e - x_a}$$
 (17)

$$b = \tan \theta_a - 2cx_a \tag{18}$$

$$a = y_a - bx_a - cx_a^2$$
 (19)

3. Specified throat attachment angle θ_a , nozzle length x_e , and nozzle area ratio ϵ . In this case, the nozzle exit lip radius y_e is given by equation (9). The coefficients a, b, and c are given by equations (17) to (19).

Tabular nozzle. If the supersonic expansion contour is tabular, a set of (x,y) pairs spanning the region from point A to point E must be specified. These points must be chosen so that a smooth transition to the circular arc throat contour is obtained. The first point in the table is a point just downstream of the throat attachment point, point A. The last point in the table specifies the nozzle exit lip point, point E.

3. NOZZLE EXTENSION GEOMETRY

The nozzle extension consists of a conical section beginning at the end of the basic nozzle, point E in Figure 3, and ending at point F. The scarfing angle ß is the angle, in the meridional plane, between the nozzle axis and the missile axis. For a conventional propulsive nozzle, the scarfing angle ß is zero. In general, the intersection of a body of revolution such as the nozzle contour with a second body of revolution such as the missile outer envelope will not yield a curve of intersection that lies within a plane. However, if the diameter of the missile is much larger than the diameter of the nozzle, the curve of intersection lies in a surface that approaches a plane as the diameter of the missile increases. In the present analysis, the curve of intersection is assumed to lie in a plane.

The geometry of the nozzle extension is completely specified by the attachment point (x_e, y_e) , the angle of the conical extension θ_f , and the length of the conical extension x_f .

If the nozzle extension attaches smoothly to the basic nozzle (i.e., $\theta_{\rm e} = \theta_{\rm f}$), then a continous flow occurs across the transition region. If $\theta_{\rm f}$ is less than $\theta_{\rm e}$, then an oblique shock wave emanates from point E and propagates downstream in the nozzle extension. Both possibilities are accounted for in the present analysis. If $\theta_{\rm f}$ is greater than $\theta_{\rm e}$, then a centered expansion wave emanates from point E and propagates downstream in the nozzle extension. That possibility is not considered in the present analysis.

SECTION III

FLOWFIELD MODEL

INTRODUCTION

The flowfield model consists of four parts.

- (a) The transonic flow analysis in the throat region of the basic nozzle.
- (b) The supersonic flow analysis in the basic nozzle.
- (c) The determination of the shock wave (or Mach line if $\theta_e = \theta_f$) which emanates from the point of attachment of the nozzle extension to the basic nozzle (point E in Figure 3) and propagates downstream into the nozzle extension.
- (d) The supersonic flow analysis in the nozzle extension downstream of the oblique shock wave.

The specification of these four flowfield models is presented in the following paragraphs.

The flow is assumed to originate upstream of the nozzle in a uniform flow region having constant stagnation pressure and temperature P_t and T_t , respectively. The flowing fluid is assumed to be a thermally and calorically perfect gas (i.e., constant molecular weight and specific heats). The presence of condensed phases and chemical reactions are neglected, and boundary layer effects are considered negligible. Consequently, the flowfield is isentropic everywhere except within the oblique shock wave, which is described by the standard oblique shock wave relationships.

The flowfield in the transonic region and in the basic nozzle is irrotational (i.e., constant entropy and stagnation enthalpy throughout) since the flow originates in a uniform flow region and is isentropic. The flowfield is rotational downstream of the oblique shock wave due to the entropy gradient produced by the curved oblique shock wave.

The overall numerical algorithm is discussed in detail in Section V.

A brief introduction to the overall numerical algorithm is presented here to identify the various flowfield models that are required to determine the performance of a scarfed propulsive nozzle.

The flow in the throat region of the nozzle is assumed to be completely specified by a perturbation analysis that depends only on the geometry of the nozzle throat. From this analysis, the flow properties along an initial-value line spanning the nozzle throat are determined, as illustrated in Figure 4.

Right-running Mach lines are then emanated from each point along the initial-value line, starting with the points adjacent to the axis, and continued until they intersect the nozzle axis. The last such right-running Hach line emanating from the nozzle throat wall point defines the extent of the initial-value problem. This region is illustrated in Figure 5.

Point locations are prespecified along the circular arc initial expansion contour downstream of the nozzle throat wall point. A right-running Mach line is emanated from each prespecified wall point and continued until it intersects the nozzle axis. The last such right-running Mach line defines the extent of the initial expansion flowfield. This region is illustrated in Figure 6.

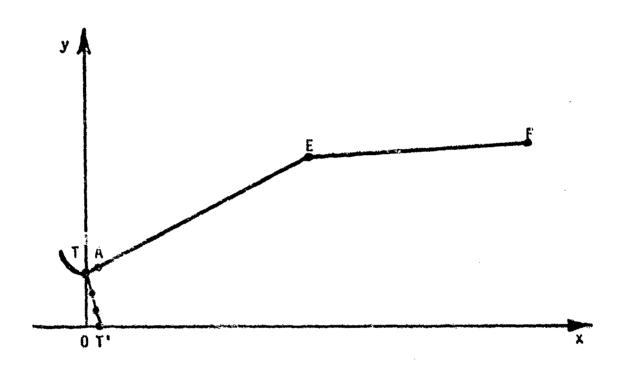


Figure 4. Initial-value line.

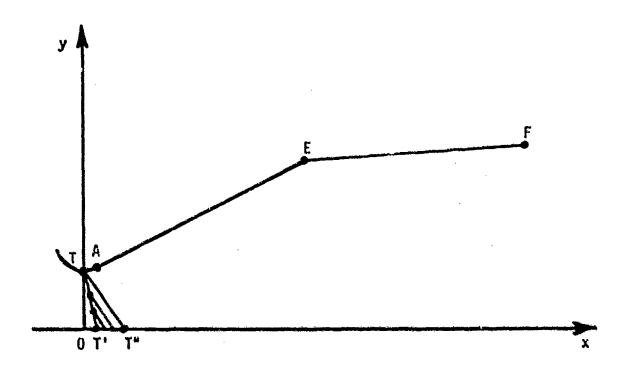


Figure 5. Extent of the initial-value problem.

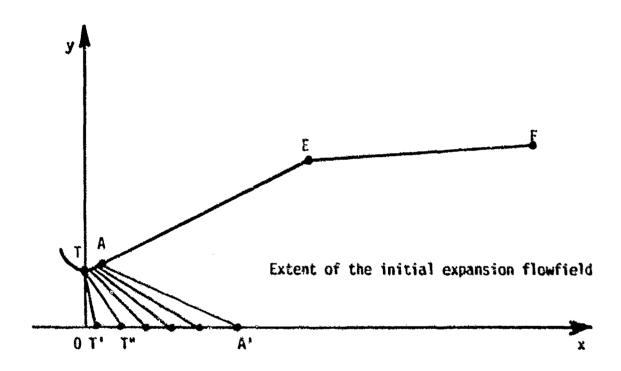


Figure 6. Flowfield from the initial expansion contour.

Left-running Mach lines are then emanated from each point along the last right-running Mach line in the initial expansion flowfield and continued until they intersect the nozzle wall. The first such left-running Mach line begins at the first interior point on the aforementioned right-running Mach line, the second left-running Mach line begins from the second interior point, and so on. This process is continued until either the end of the basic nozzle is reached, or the last point on the right-running Mach line, which lies on the nozzle axis, has been reached. In this latter case, left-running Mach lines are emanated from the nozzle axis and continued until they intersect the nozzle wall. In either case, a left-running Mach line will eventually reach the end of the basic nozzle contour, point E. This flowfield is illustrated in Figure 7.

Due to the finite change in wall slope at point E, an oblique shock wave forms at that point. That right-running shock wave propagates into the flowfield and intersects the left-running Mach lines as they propagate toward the wall. The flowfield up to the shock wave is irrotational, and the flowfield downstream of the shock wave is rotational. The flow property changes across the shock wave depend on the strength of the shock wave, which is determined by both the upstream flowfield and the downstream flowfield. After passing through the oblique shock wave, the left-running Mach lines continue until they intersect the scarfed nozzle excension. This process is continued until the end of the scarfed nozzle extension, point F, is reached. This flowfield is illustrated in Figure 8.

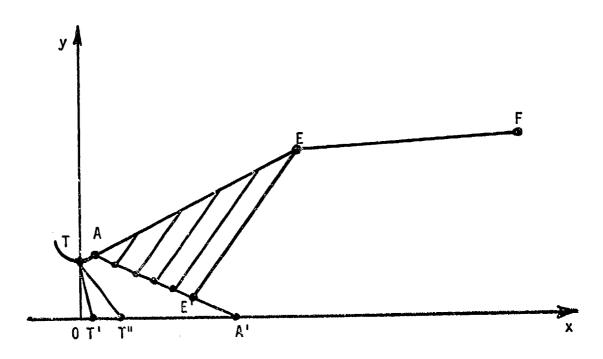


Figure 7. Flowfield in the basic nozzle.

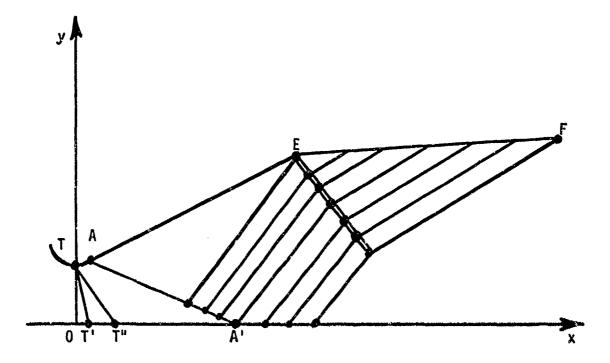


Figure 8. Flowfield in the scarfed nozzle extension.

Once the entire scarfed nozzle flowfield has been calculated as described above, the performance of the scarfed nozzle can be computed. That computation is described in Section IV.

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2. TRANSONIC FLOW MODEL

A. FLOWFIELD

The flowfield in the transonic throat region of the basic nozzle is assumed to be completely determined by the geometry of the nozzle throat. The influence of the upstream subsonic geometry is assumed to be negligible. This is a good approximation when the nozzle inlet is "well behaved." Figure 9 illustrates the throat geometry.

As illustrated in Figure 9, the sonic line originates on the nozzle wall upstream of the throat (point T), spans the throat region, and intersects the nozzle axis downstream of the throat. When the throat upstream radius of curvature ρ_{tu} is large compared to the throat radius y_t , linearized flow analyses give reasonable predictions for the transonic flowfield. Zucrow and Hoffman (1) present a discussion of such analyses.

In the present case, the analysis developed by Kliegel and Levine (2) is employed. The velocity components u and v in the x and y directions, respectively, are given by

$$\frac{u(x,y)}{a^*} = 1 \div \frac{u_1(\bar{x},\bar{y})}{(R+1)} + \frac{1}{(R+1)^2} [u_1(\bar{x},\bar{y}) + u_2(\bar{x},\bar{y})] + \frac{1}{(R+1)^3} [u_1(\bar{x},\bar{y}) + 2u_2(\bar{x},\bar{y}) + u_3(\bar{x},\bar{y})]$$
(20)

$$\frac{v(\bar{x},\bar{y})}{a^*} = \left[\frac{\gamma+1}{2(R+1)}\right]^{1/2} \left\{ \frac{v_1(\bar{x},\bar{y})}{(R+1)} + \frac{1}{(R+1)^2} \left[\frac{3}{2} v_1(\bar{x},\bar{y}) + v_2(\bar{x},\bar{y}) \right] + \frac{1}{(R+1)^3} \left[\frac{15}{8} v_1(\bar{x},\bar{y}) + \frac{5}{2} v_2(\bar{x},\bar{y}) + v_3(\bar{x},\bar{y}) \right] \right\} (21)$$

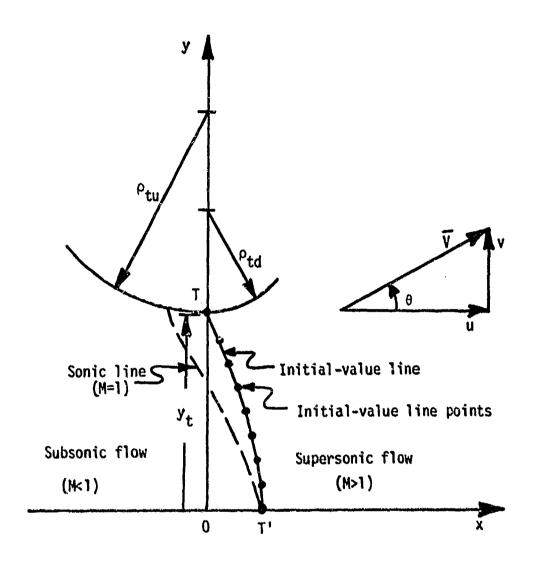


Figure 9. Nozzle throat geometry

where

$$a^* = (\gamma RT^*)^{1/2} = \left(\frac{2\gamma RT_t}{\gamma + 1}\right)^{1/2}$$
 (22)

$$R = \frac{\rho_{tu}}{y_t}, \ \bar{x} = \left(\frac{2R}{\gamma+1}\right)^{1/2} \frac{x}{y_t}, \ \bar{y} = \frac{y}{y_t}$$
 (23)

 T_t is the stagnation temperature, γ and R are the gas specific heat ratio and gas constant, respectively, and $u_1(\bar{x},\bar{y})$, etc., are polynomial expressions in \bar{x} and \bar{y} given by equations (15.100) to (15.105) in Reference (1).

In the present analysis, the initial-value line is a parabola extending from the nozzle throat wall point, point T in Figure 9, to the point on the nozzle axis where M = 1.0, as illustrated in Figure 9. The particular parabola employed is the curve along which the flow angle θ [θ = tan⁻¹(v/u)] is equal to zero in Sauer's linearized flow analysis [Reference (3), see also Reference (1)]. That parabola is specified by

$$x = \varepsilon - \frac{(\gamma+1)\alpha y^2}{2(3+\delta)}$$
 (24)

where $\delta = 0$ denotes planar flow, $\delta = 1$ denotes axisymmetric flow, and

$$\varepsilon = -\frac{y_t}{2(3+\delta)} \left[\frac{(\gamma+1)(1+\delta)}{(\rho_{t+1}/y_t)} \right]^{1/2}$$
 (25)

$$\alpha = \left[\frac{1+\delta}{(\gamma+1)\rho_{+1}y_{+}} \right]^{1/2} \tag{26}$$

The initial-value line is determined by specifying the y coordinates of NI points across the throat, and calculating the corresponding values of x from equation (24). Figure 10illustrates a typical result. These sets of (x,y) pairs are substituted into equations (20) and (21) to determine the velocity components u(x,y) and v(x,y). Then the velocity magnitude V is determined from

$$V = (u^2 + v^2)^{1/2} \tag{27}$$

From the energy equation for the adiabatic flow of a thermally and calorically perfect gas [see Reference (4)],

$$T = T_{t} - \frac{V^2}{2c_{p}} \tag{28}$$

where $\mathbf{c}_{\mathbf{p}}$ is the constant pressure specific heat. The speed of sound a is given by

$$a = (\gamma RT)^{1/2} \tag{29}$$

and the Mach number M is defined as

$$N = \frac{V}{a} \tag{30}$$

From the definition of the stagnation pressure P_{t} for a thermally and calorically perfect gas [see Reference (4)].

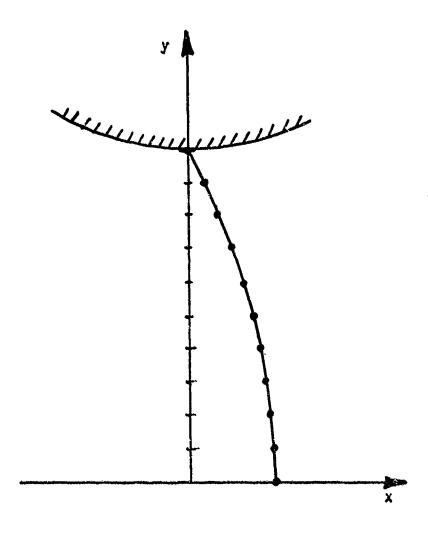


Figure 10. Supersonic initial-value line.

$$P = P_{t} \left(\frac{T}{T_{t}} \right) \frac{\gamma}{\gamma - 1}$$
 (31)

From the thermal equation of state, the density ρ is given by

$$\rho = \frac{P}{RT} \tag{32}$$

Thus, all the flow properties along the initial-value line have been determined.

B. PERFORMANCE PARAMETERS

The nozzle mass flow rate $\hat{\mathbf{m}}$ is obtained by numerically integrating the differential mass flow rate $d\hat{\mathbf{m}}$ across the initial-value line (IVL). Thus,

$$dh = \rho u 2\pi y dy - \rho v 2\pi y dx \tag{33}$$

$$\hat{m} = \int_{IVL} d\hat{m}$$
 (34)

The nozzle discharge coefficient $\mathbf{C}_{\boldsymbol{D}}$ is defined as

$$C_{D} = \frac{\dot{m}}{m_{1-D}} \tag{35}$$

where the reference ideal one-dimensional isentropic choked mass flow rate \hat{m}_{1-D} is given by [see Reference (4)]

$$\hat{\mathbf{m}}_{1-D} = \frac{\mathbf{r}_{t} \mathbf{A}_{t}}{(\gamma R \mathbf{T}_{t})^{1/2}} \tag{36}$$

where the parameter Γ is defined as

$$\Gamma = \gamma \left(\frac{2}{\gamma+1}\right)^{\frac{\gamma+1}{2(\gamma-1)}} \tag{37}$$

The initial-value line thrust $F_{\mbox{IVL}}$ is the sum of the axial components of the pressure forces and the momentum flux across the initial-value line. Thus,

$$dF = (P - P_a) 2\pi y dy + u d\hat{m}$$
 (38)

$$F = \int_{IVL} dF \tag{39}$$

where P_{a} is the atmospheric pressure. The thrust efficiency η_{F} is defined as

$$\eta_{\mathsf{F}} = \frac{\mathsf{F}}{\mathsf{F}_{\mathsf{1-0}}} \tag{40}$$

where the reference ideal one-dimensional isentropic choked thrust is given by

$$F_{1-D} = (P^* - P_a) A_t + a^* m_{1-D}$$
 (41)

where the throat sonic velocity a* is given by equation (22) and the

throat choking pressure P* is given by [see Reference (4)]

$$P* = P_{t} \left(\frac{2}{\gamma+1}\right)^{\frac{\gamma}{\gamma-1}} \tag{42}$$

The initial-value line specific impulse \mathbf{I}_{sp} is defined as

$$I_{SP} = \frac{F}{m} \tag{43}$$

The specific impulse efficiency $\boldsymbol{n}_{\boldsymbol{I}}$ is defined as

$$\eta_{I} = \frac{I_{Sp}}{I_{Sp,1-D}} \tag{44}$$

where the reference ideal one-dimensional specific impulse is defined as

$$I_{sp,1-D} = \frac{F_{1-D}}{\dot{m}_{1-D}} \tag{45}$$

By combining equations (35), (40), and (44), it is obvious that

$$\eta_{F} = C_{D} \eta_{I} \tag{46}$$

The above analysis completely determines all of the properties of interest along the supersonic initial-value line. Some numerical examples are presented in Reference (1).

3. IRROTATIONAL FLOW MODEL

A. GOVERNING EQUATIONS

The portion of the supersonic flowfield between the supersonic initial-value line and the oblique shock wave emanating from the wall discontinuity where the basic nozzle and the nozzle extension meet (i.e., point E) is an irrotational flowfield. The basic equations governing an adiabatic inviscid flow are the continuity equation, the Euler momentum equations, and the energy equation. These equations are derived in Reference (5). For steady two-dimensional flow, those equations are

$$\rho u_{x} + \rho v_{y} + u \rho_{x} + v \rho_{y} + \frac{\delta \rho v}{y} = 0$$
 (47)

$$\rho_{\text{UU}_{X}} + \rho_{\text{VU}_{Y}} + P_{X} = 0 \tag{48}$$

$$\rho u v_{\chi} + \rho v v_{y} + P_{y} = 0 \tag{49}$$

$$uh_x + vh_y + u(\frac{v^2}{2})_x + v(\frac{v^2}{2})_y = 0$$
 (50)

where $\delta = 0$ for planar flow and $\delta = 1$ for axisymmetric flow, and h is the enthalpy of the fluid. For isentropic (i.e., adiabatic and frictionless) flows, it is convenient to replace the energy equation, equation (50), by the speed of sound equation [see Reference (5)] to eliminate the enthalpy h from the system of equations. That equation is

$$u_X^p + v_y^p - a^2(u_{\rho_X} + v_{\rho_y}) = 0$$
 (51)

Combining equations (47) and (51) to eliminate the derivatives of density yields

$$\rho a^{2}(u_{x} + v_{y} + \frac{\delta v}{y}) + uP_{x} + vP_{y} = 0$$
 (52)

Equations (47), (48), (49), and (52) comprise a set of four coupled partial differential equations for determining the four flow properties u, v, P, and ρ . Those equations are the basis of the rotational flow model employed downstream of the oblique shock wave. That flow model is developed further in Section III.5.

If the flowfield originates in a region of parallel uniform flow and remains isentropic, then the vorticity of the flowfield is everywhere zero [see Reference (5)]. Thus,

$$\bar{\zeta} = \nabla x \bar{V} = 0 \tag{53}$$

which, for two-dimensional flow, gives

$$u_y - v_v = 0 \tag{54}$$

A single equation containing derivatives only of u and v can be obtained by multiplying equations (48) and (49) by u and v, respectively, adding them together, and subtracting equation (52). The result

$$\rho(u^2 - a^2)u_x + \rho(v^2 - a^2)v_y + \rho u v(u_y + v_x) - \frac{\delta \rho a^2 v}{y} = 0$$
 (55)

Dividing by the density ρ and introducing equation (54) gives

$$(u^2 - a^2)u_x + (v^2 - a^2)v_y + 2uvu_y - \frac{\delta a^2 v}{v} = 0$$
 (56)

Equation (56) is sometimes called the gasdynamic equation. When considered in conjunction with equation (54), the irrotationality condition, a set of two partial differential equations is obtained for determining the two velocity components u and v. The remaining flow properties (i.e., V, T, P, ρ , and M) can be determined as discussed in Section III.2 for transonic flow.

The flow model employed for the irrotational flowfield between the nozzle throat and the oblique shock wave consists of equations (36) and (54), which are repeated and renumbered below.

$$(u^2 - a^2)u_x + (v^2 - a^2)v_y + 2uvu_y - \frac{\delta a^2 v}{y} = 0$$
 (57)

$$\mathbf{u}_{\mathbf{v}} - \mathbf{v}_{\mathbf{x}} = 0 \tag{58}$$

B. METHOD OF CHARACTERISTICS

In the present analysis, equations (57) and (58) are solved by the numerical method of characteristics. Characteristics are curves in the solution space (i.e., the xy plane) along which the partial differential

equations combine to form ordinary differential equations. That is, the partial differential equations become directional derivatives along the characteristic curves. The basic principles of the method of characteristics are discussed in Reference (6). Application of the method of characteristics to steady two-dimensional irrotational supersonic flow is discussed in Reference (7). A brief summary of those results is presented below.

C. CHARACTERISTIC EQUATIONS

For steady two-dimensional irrotational supersonic flow, two families of characteristic curves exist; the left-running and right-running Mach lines. The slope λ of a characteristic curve is defined as

$$\frac{\mathrm{d}y}{\mathrm{d}x} = \lambda \tag{59}$$

A quadratic equation can be found for determining λ . As shown in Reference (6), that equation is

$$(u^2 - a^2)\lambda^2 - 2uv\lambda + (v^2 - a^2) = 0$$
 (60)

Solving equation (60) for λ yields

$$\lambda_{\pm} = (\frac{dy}{dx})_{\pm} = \frac{uv \pm a^2 \sqrt{M^2 - 1}}{u^2 - a^2}$$
 (61)

Equation (61) can be simplified using the definitions

$$u = V \cos \theta$$
, $v = V \sin \theta$, $\theta = \tan^{-1}(v/u)$, $\alpha = \sin^{-1}(1/M)$ (62)

where θ is the flow angle and α is the Mach angle. After considerable manipulation, the following result is obtained.

$$\left(\frac{\mathrm{d}y}{\mathrm{d}x}\right)_{\pm} = \lambda_{\pm} = \tan \left(\theta \pm \alpha\right) \tag{63}$$

Equation (63), which defines the characteristic curves, is called the characteristic equation. The subscript + denotes the C₊ characteristic, or Teft-running Mach line, so called because it runs off to the left of the streamline when looking in the downstream direction. Similarly, the subscript - denotes the C₋ characteristic, or right-running Mach line. Figure II illustrates these characteristic curves, or Mach lines, at a point in a flowfield. Two Mach lines pass through every point in a flowfield, resulting in two infinite families of left-running and right-running Mach lines.

D. COMPATIBILITY RELATIONS

Equations (57) and (58), when combined and written as directional derivatives along the Mach lines, yield a single ordinary differential equation, called the compatibility relation, which is valid along each family of Mach lines. Consider any continuous function f(x,y). The total derivative of f is given by

$$df = f_X dx + f_V dy ag{64}$$

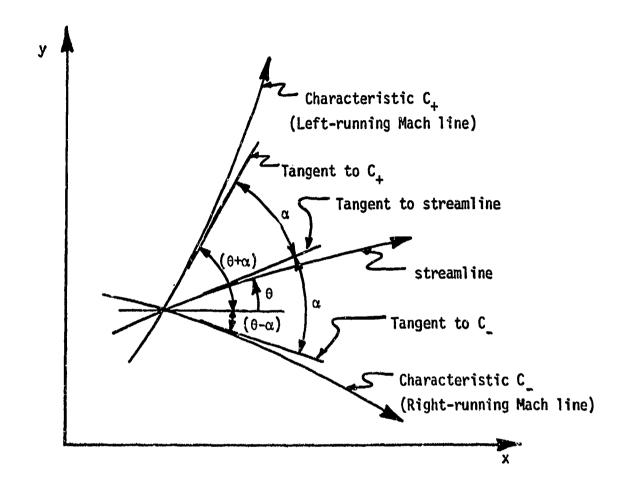


Figure 11. Characteristic curves for irrotational flow.

which may be written as

$$\frac{df}{dx} = f_x + f_y \frac{dy}{dx} = f_x + \lambda f_y$$
 (65)

The partial derivatives in equations (57) and (58) can be put in the form of equation (65) by multiplying equation (58) by the term

$$(u^2 - a^2) \lambda - 2uv$$
 (66)

and adding the result to equation (57) to obtain

$$(u^{2} - a^{2})u_{x} + (v^{2} - a^{2})v_{y} + 2uvu_{y} - \frac{\delta a^{2}v}{y} +$$

$$[(u^{2} - a^{2})_{\lambda} - 2uv](u_{y} - v_{x}) = 0$$
(67)

Rearranging equation (67) yields

$$(u^{2} - a^{2})[u_{x} + \lambda u_{y}] + [2uv - (u^{2} - a^{2})\lambda] v_{x} + (v^{2} - a^{2})v_{y} - \frac{6a^{2}v}{y} = 0$$
(68)

From equation (60),

$$(v^2-a^2) = [2uv - (u^2-a^2)\lambda] \lambda$$
 (69)

Substituting equation (69) into equation (68) gives

The partial derivatives in equation (70) are in the form of equation (65). Consequently, equation (70) may be written as

$$(u^2 - a^2) du + [2uv - (u^2 - a^2)\lambda] dv - \frac{\delta a^2 v}{y} dx = 0$$
 (71)

Equation (71) is a total differential equation which is valid along the Mach lines. It is called the compatibility relation. When applied along both left-running and right-running Mach lines, equation (71) yields two equations for determining changes in u and v along the two families of Mach lines.

E. SUMMARY OF RESULTS

In summary, for steady two-dimensional irrotational supersonic flow, the characteristic equations are

$$\left(\frac{\mathrm{d}y}{\mathrm{d}x}\right)_{\pm} = \lambda_{\pm} = \tan \left(\theta \pm \alpha\right) \tag{72}$$

and the compatibility relations are

$$(u^2-a^2) du_{\pm} + [2uv - (u^2-a^2) \lambda_{\pm}] dv_{\pm} - \frac{\delta a^2 v}{y} dx_{\pm} = 0$$
 (73)

F. NUMERICAL METHOD OF CHARACTERISTICS

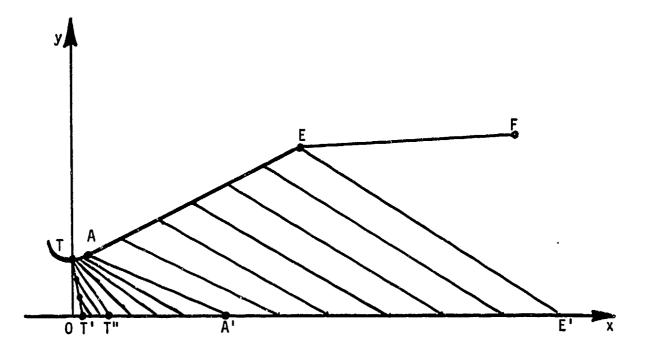
The characteristic equations and compatibility relations for steady

two-dimensional irrotational supersonic flow are presented in the preceding section. Those equations are ordinary differential equations. They can be integrated numerically by the second-order accurate modified-Euler predictor-corrector method. Such a procedure is called the numerical method of characteristics. The details of its implementation are presented in References (6) and (7). That is the procedure employed to calculate the flowfield between the throat of the basic nozzle and the oblique shock wave at the entrance to the scarfed extension.

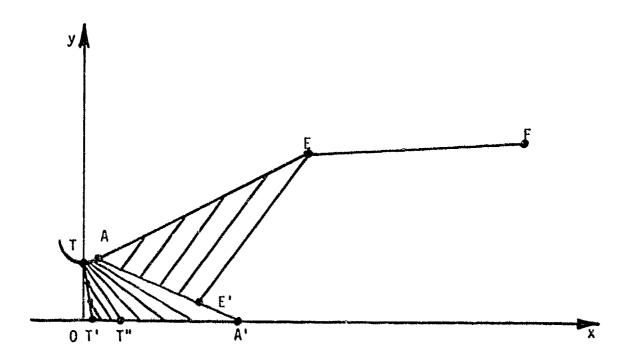
The basic procedure is to start from a supersonic initial-value line across the nozzle throat, such as illustrated in Figure 10. The left-running and right-running Mach lines emanating from the initialvalue line points are constructed by solving equation (72). The velocity components u and v are determined at the intersections of the Mach lines by solving equation (73). Since equations (72) and (73) are coupled and nonlinear, they must be solved numerically and simultaneously. This procedure defines the interior point unit process. When a left-running Mach line impinges on the nozzie wall, the boundary conditions of known wall location and slope must be applied. A rightrunning Mach line is reflected into the flowfield from that point of impingement. That procedure defines the wall point unit process. Similarly, when a right-running Mach line reaches the nozzle axis, the boundary conditions of y = v = 0 must be applied. A left-running Mach line is reflected into the flowfield from that point of impingement. That procedure defines the axis point unit process.

This procedure is continued until the entire flowfield between the

nozzle throat and the oblique shock wave has been crisscrossed with left-running and right-running Mach lines and the flow properties have been calculated at all the intersections of the two families of Mach lines. A coarse schematic of the resulting Mach line network is illustrated in Figure 12. Figure 12(a) illustrates the Mach line network when right-running Mach lines are used to construct the flow-field. Figure 12(b) illustrates the Mach line network when left-running Mach lines are used to construct the flow-field.



(a) Right-running Mach line network.



(b) Left-running Mach line network.

Figure 12. Mach line network in the basic nozzle.

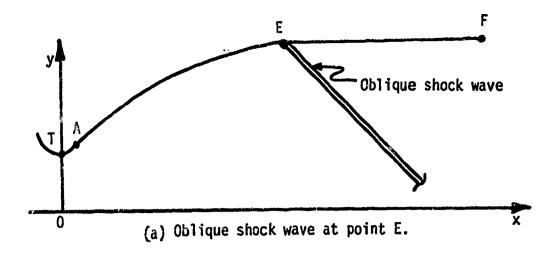
4. SHOCK WAVE MODEL

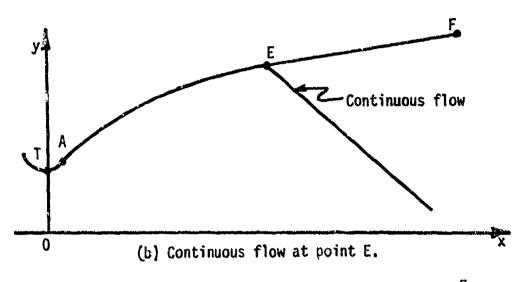
A. PHYSICAL MODEL

The nozzle geometric model considered in the present investigation consists of a basic nozzle followed by a nozzle extension, as illustrated in Figure 13. As discussed in Section II.1, the basic nozzle may be conical, quadratic, or tabular. The nozzle extension is conical.

At the junction of the basic nozzle and the nozzle extension, the slope of the wall of the basic nozzle may be larger than, equal to, or less than the slope of the nozzle extension, resulting in the generation of an oblique shock wave, a continuous flow, or a centered expansion fan, respectively, as illustrated in Figure 13. The third case, where the slope of the basic nozzle is less than the slope of the nozzle extension, is not considered in the present study. The second case, where the slope of the basic nozzle is equal to the slope of the nozzle extension and the flow is continuous, is treated as the special case of an oblique shock wave of infinitesimal strength. Consequently, in the present analysis, it is always assumed that an oblique shock wave emanates from the junction of the basic nozzle and the nozzle extension. The flow model for determining the location of and the flow properties across the oblique shock wave are discussed in this section.

The strength of an oblique shock wave is always greater than a continuous flow. Consequently, the oblique shock emanating from the junction point will propagate upstream into the flowfield determined by the basic nozzle, as illustrated in Figure 14. The strength of the oblique shock wave is influenced by both the upstream flowfield





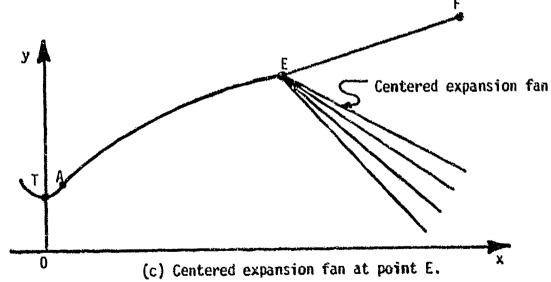


Figure 13. Flowfield at point E.

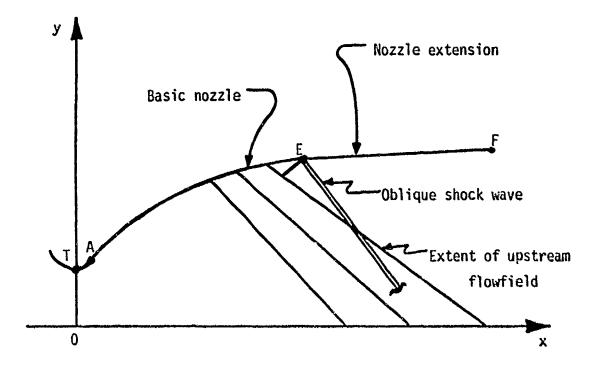


Figure 14. Oblique shock wave.

determined by the basic nozzle and the downstream flowfield determined by the nozzle extension.

B. GOVERNING EQUATIONS

The basic geometry of an oblique shock wave is illustrated in Figure 15. The upstream velocity is denoted by V_1 , the flow turning angle between the upstream and downstream flow directions is denoted by δ , and the downstream velocity is denoted by V_2 . The shock wave angle is denoted by ϵ . All of the flow properties (i.e., V, M, P, T, ρ , P_t , and T_t) are known upstream of the oblique shock wave. The problem is to determine the corresponding flow properties downstream of the shock wave for a specified value of the flow turning angle δ . A detailed analysis of the oblique shock wave is presented in Reference (8).

A coordinate system normal and tangential to the oblique shock wave is employed. The equations of continuity, momentum, and energy are applied in that coordinate system. A major conclusion of that analysis is that the tangential component of velocity does not change across the oblique shock wave (i.e., $V_{T1} = V_{T2} = V_{T}$), and that the components of velocity normal to the shock wave (i.e., V_{N1} and V_{N2}) are governed by the equations for a normal shock wave.

If the wave angle ε is known, the upstream normal Mach number is given by

$$H_{R1} = H_1 \sin \varepsilon$$
 (74)

From Figure 15, the downstream normal Nach number $N_{\rm N2}$ is given by [see Reference (8)]

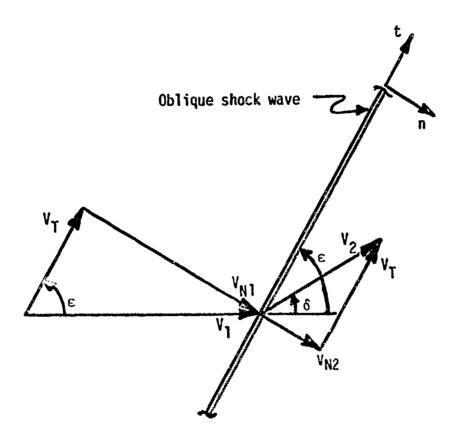


Figure 15. Oblique shock wave geometry.

$$M_{N2} = \left[\frac{1 + \frac{\gamma - 1}{2} M_{N1}^2}{\gamma M_{N1}^2 - \frac{\gamma - 1}{2}} \right]^{1/2}$$
 (75)

From Figure 10, the downstream Mach number is

$$M_2 = \frac{M_{N2}}{\sin (\varepsilon - \delta)} \tag{76}$$

For a given wave angle ε , the corresponding flow turning angle δ is given by [see Reference (8)]

$$\frac{1}{\tan \delta} = \left[\frac{\Upsilon + 1}{2} - \frac{M_1^2}{M_{N1}^2 - 1} - 1 \right] \tan \varepsilon \tag{77}$$

If the flow turning angle δ is specified instead of the wave angle ϵ , then equations (74) to (77) must be solved by iteration (e.g., using the secant method). In either case, for a given M₁ and either ϵ or δ , the corresponding M_{N1}, M_{N2}, M₂, and either ϵ or δ can be determined.

The remaining flow property ratios are obtained from the corresponding results for a normal shock wave [see Reference (8)].

$$\frac{P_2}{P_1} = \frac{2\gamma}{\gamma - 1} M_{N1}^2 - \frac{\gamma - 1}{\gamma + 1} \tag{78}$$

$$\frac{\rho_2}{\rho_1} = \frac{(\gamma+1) M_{N1}^2}{2 + (\gamma-1)M_{N1}^2} \tag{79}$$

$$\frac{V_{N2}}{V_{N1}} = \frac{\rho_1}{\rho_2} \tag{80}$$

$$V_{T} = V_{1} \cos \varepsilon \tag{81}$$

$$V_2 = (V_{N2}^2 + V_T^2)^{1/2}$$
 (82)

C. NUMERICAL SOLUTION PROCEDURE

Two different oblique shock wave situations occur in a scarfed nozzle. The first situation is at the junction of the basic nozzle and the nozzle extension, point E in Figure 14, where the oblique shock wave emanates. At that point, the flow turning angle δ is specified as the difference between the flow angle at the end of the basic nozzle and the flow angle at the beginning of the nozzle extension. The procedure described in the preceding paragraphs can be applied directly to determine the properties of the oblique shock wave at that point.

The second situation is the general point on the oblique shock wave where both the upstream and downstream flowfield affect the shock wave. Since the strength of an oblique shock wave is greater than the strength of a Mach line, the oblique shock wave emanating from point E propagates into the upstream flowfield at a steeper angle than the angle of the right-running Mach lines emanating from the basic nozzle contour. Since a right-running oblique shock wave turns the flow toward the axial direction (i.e., toward the shock wave) and the Mach number and Mach angle are smaller downstream of an oblique shock wave, the slopes of the right-running Mach lines on the downstream side of the shock wave are greater that the slope of the shock wave. Consequently, the location and stength of the oblique shock wave are also influenced by the downstream flowfield. Thus, the portion of the downstream flowfield that interacts with the oblique shock wave must be determined along with

the determination of the shock wave itself. Figure 16 illustrates the propagation of the oblique shock wave into the upstream flowfield and the overtaking of the shock wave on the downstream side by a Mach line from the nozzle extension.

The region of the flowfield affecting the solution at a point on the oblique shock wave, denoted as point 4, is outlined in Figure 16.

The upstream flowfield is assumed to be specified, and the propagation of the oblique shock wave from point E to point 5 has been calculated, as has the right-running Mach line on the downstream side of the shock wave from the nozzle extension to the left-running Mach line from point 5.

An enlarged illustration of this region is presented in Figure 17.

The determination of the location of point 4 and the flow properties on either side of the oblique shock wave at that point requires an iterative procedure having five major steps. The first step is to find the location of point 4 as the intersection of the oblique shock propagated from point 5 at the average wave angle ε with the left-running Mach line propagated from point 2 in the known upstream flowfield. This is accomplished in an iterative manner by first assuming a value for ε_4 (note that ε_5 is known), performing the remaining four major steps outlined below, and then varying ε_4 iteratively until the overall procedure converges.

The second step is the application of the method of characteristics for irrotational flow (see Section III.3) to determine the flow properties at point 4 on the upstream side of the oblique shock wave. The third step is the determination of the flow properties at point 4 on the downstream side of the oblique shock wave by applying the oblique shock

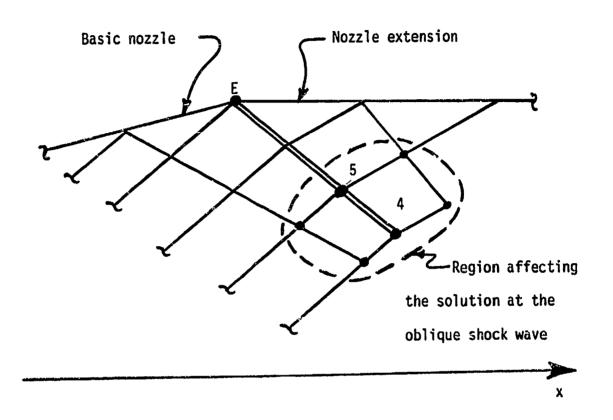
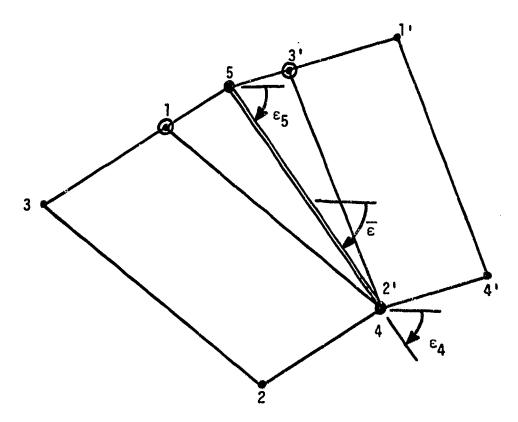


Figure 16. Propagation of the oblique shock wave.



Known points:

1', 2, 3, and 5

Solution points:

4, 2', and 4'

Interpolated points: 1 and 3'

Figure 17. Oblique shock wave point unit process.

wave. analysis presented at the beginning of this section. The fourth step is the application of the method of characteristics for rotational flow (see Section III.5) to determine the flow properties at point 4', which is the intersection of the left-running Mach line from point 4 on the downstream side of the oblique shock wave with the right-running Mach line from point 1'.

All of the above four steps are performed for a specified value of ϵ_4 , the oblique shock wave angle at point 4, which together with the known wave angle ϵ_5 at point 5, defines the average wave angle $\bar{\epsilon}$ and thus the location of point 4. A check on the correctness of ϵ_4 is made as the fifth step in the overall procedure. The pressure P_{4d} on the downstream side of the oblique shock wave is computed in step 3 above. A right-running Mach line is projected rearward from point 4 to intersect a previous Mach line (or the wall of the nozzle extension for the first point on the oblique shock wave downstream of point E) at point 3'. The compatibility relation valid along that Mach line is solved to obtain a second value for P_{4d} . The discrepancy detween these two values of P_{4d} is a measure of the error in ϵ_4 . The secant method can be used to vary ϵ_4 to drive the difference between the two values of P_{4d} to within any desired convergence limit.

The procedure described above is repeated at successive points along the oblique shock wave until the shock wave approaches the nozzle axis.

5. ROTATIONAL FLOW MODEL

A. GOVERNING EQUATIONS

The portion of the flowfield downstream of the oblique shock wave emanating from the junction of the basic nozzle and the nozzle extension is a rotational flowfield due to the transverse entropy gradient caused by the curved oblique shock wave. The basic equations governing a rotational isentropic flowfield are presented in Section III.3 [equations (47), (48), (49), and (51)]. Those equations are repeated and renumbered below.

$$\rho u_{X} + \rho v_{y} + u \rho_{X} + v \rho_{y} + \frac{\delta \rho v}{y} = 0$$
 (83)

$$\rho u u_{x} + \rho v u_{y} + P_{x} = 0$$
 (84)

$$\rho u v_{x} + \rho v v_{y} + P_{y} = 0$$
 (85)

$$uP_x + vP_y - a^2(u\rho_x + v\rho_y) = 0$$
 (86)

Due to the presence of an entropy gradient ∇s , the vorticity $\ddot{\zeta}$ is not zero, and equation (53) is not valid. Consequently, the simplifications obtained in equations (57) and (58) are not applicable, and the full set of equations presented above, equations (83) to (86), must be solved. Those equations are solved by the numerical method of characteristics (see Section III.3.8).

B. CHARACTERISTIC EQUATIONS

For steady two-dimensional rotational supersonic flow, three

families of characteristic curves exist; the left-running and right-running Mach lines and the streamlines. The slope λ of a characteristic curve is defined as

$$\frac{dy}{dx} = \lambda \tag{87}$$

A fourth-order equation can be found for determining λ . As shown in Reference (9), that equation is

$$(u\lambda - v)^{2} [(u\lambda - v)^{2} - a^{2}(1 + \lambda^{2})] = 0$$
 (88)

The fourth-order equation is the product of two second-order equations. Solving either of those two second-order equations gives two of the characteristic slopes, and solving the other second-order equation gives the other two characteristic slopes.

Solving the first second-order term yields

$$(u\lambda_0 - v)^2 = 0 \tag{89}$$

$$\lambda_{0} = \left(\frac{dy}{dx}\right)_{0} = \frac{v}{u} \tag{90}$$

which defines the streamline. Since λ_0 appears two times, the streamline is a repeated characteristic. The subscript $_0$ will be used henceforth to denote the streamline.

Rearranging the second second-order term yields

$$(u^2 - a^2)\lambda^2 - 2uv\lambda + (v^2 - a^2) = 0 (91)$$

Equation (91) is identical to equation (60), which defines the Mach lines. As shown in Section III.3.C, the Mach lines are given by

$$\left(\frac{\mathrm{d}y}{\mathrm{d}x}\right)_{\pm} = \lambda_{\pm} = \tan \left(\theta \pm \alpha\right) \tag{92}$$

The subscripts ± are used to denote the left-running and right-running Mach lines, respectively.

Consequently, for steady two-dimensional supersonic rotational flow, four characteristic curves are obtained: the streamline repeated twice and the left-running and right-running Mach lines. Figure 18 illustrates these characteristic curves at a point in a flowfield. A streamline and two Mach lines pass through every point in a flowfield, resulting in three infinite families of characteristic curves.

C. COMPATIBILITY RELATIONS

Equations (83) to (86), when combined and written as directional derivatives along the streamlines and Mach lines, yield two ordinary differential equations valid along the streamlines and one ordinary differential equation valid along each Mach line. Thus, four compatibility relations are determined for steady two-dimensional supersonic rotational flow.

As shown in Section III.3.D, the total derivative of any continuous function f(x,y) must satisfy the relation [see equation (65)]

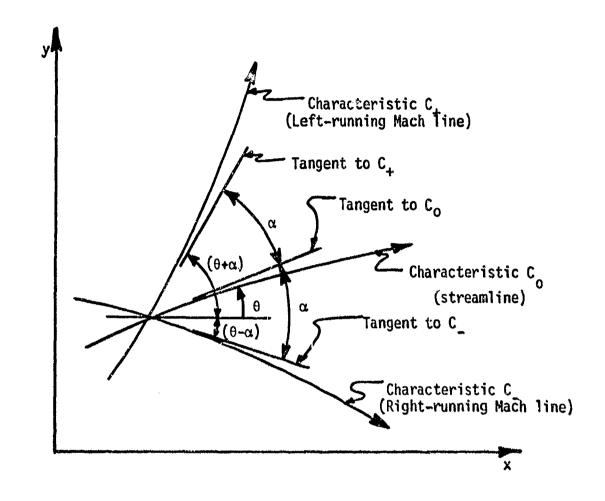


Figure 18. Characteristic curves for rotational flow.

$$\frac{df}{dx} = f_x + \lambda f_y \tag{93}$$

The partial derivatives in equations (83) to (86) can be put in the form of equation (93) by forming appropriate linear combinations of those equations.

First, consider the streamline where λ_0 satisfies equations (89) and (90). Equation (86) may be rearranged as

$$u(P_x + \frac{v}{u}P_y) - ua^2(\rho_x + \frac{v}{u}\rho_y) = 0$$
 (94)

The derivatives in equation (94) are in the form of equation (93), where $\lambda = \lambda_0 = v/u$. Consequently, along streamlines, equation (94) may be written as

$$dP - a^2 d\rho = 0$$
 (on streamlines) (95)

Equation (95) is the speed of sound equation.

The second compatibility relation valid along streamlines can be determined by multiplying equation (85) by λ and adding the result to equation (84).

$$\rho u \left(u_{x} + \frac{v}{u} u_{y} \right) + \lambda \rho u \left(v_{x} + \frac{v}{u} v_{y} \right) + \left(P_{x} + \lambda P_{y} \right) = 0$$
 (96)

The derivatives in equation (96) are in the form of equation (93), where $\lambda = \lambda_0 = v/u$. Consequently, noting that along streamlines $\lambda u = v$, equation (96) may be written as

$$\rho u du + \rho v dv + dP = 0 (97)$$

Equation (97) is Bernoulli's equation for steady two-dimensional isentropic flow, which may be written as

$$\rho V dV + dP = 0$$
 (on streamlines) (98)

Second, consider the Mach lines where λ_{\pm} satisfies equations (91) and (92). The compatibility relation valid along Mach lines is obtained by forming the following sum:

$$(u\lambda_{\pm} - v)$$
 [Eq. (82)] + $(-\lambda_{\pm})$ [Eq. (80)] + [Eq. (81)] + $\frac{(u\lambda_{\pm} - v)}{a^2}$ [Eq. (82)] = 0

After considerable manipulation, equation (99) reduces to

$$(\rho v)du_{\pm} - (\rho u)dv_{\pm} + [\lambda_{\pm} - u(u\lambda_{\pm} - v)/a^{2}]dP_{\pm} - \delta[\rho v(u\lambda_{\pm} - v)/y]dx_{\pm} = 0$$
 (100)

Equation (100) can be simplified by using the definitions

$$u = V \cos\theta$$
, $v = V \sin\theta$, $\theta = \tan^{-1}(v/u)$, $\alpha = \sin^{-1}(1/N)$ (101)

where θ is the flow angle and α is the Nach angle. After considerable manipulation, the following result is obtained.

$$\frac{\sqrt{M^2-1}}{\rho V^2} dP_{\pm} \pm d\theta_{\pm} + S \left[\frac{\sin \theta}{y N \cos (\theta \pm \alpha)} \right] dx_{\pm} = 0$$
 (102)

where the upper subscripts on dP, d0, and dx and the upper signs in \pm d0 and cos (0 \pm α) correspond to the left-running Mach line, and vice versa.

D. SUMMARY OF RESULTS

In summary, for steady two-dimensional rotational supersonic flow, three infinite families of characteristic curves exist; the streamlines and the left-running and right-running Nach lines. Along the streamlines,

$$\left(\frac{\mathrm{d}y}{\mathrm{d}x}\right)_{0} = \lambda_{0} = \frac{v}{u} \tag{103}$$

$$dP_0 - a^2 d\rho_0 = 0 ag{104}$$

$$\rho V dV_{o} + dP_{o} = 0 \tag{105}$$

Along the Hach lines,

$$\left(\frac{\mathrm{d}v}{\mathrm{d}x}\right)_{\pm} = \lambda_{\pm} = \tan\left(\theta \pm \alpha\right) \tag{106}$$

$$\frac{\sqrt{N^2-1}}{\rho V^2} dP_{\pm} \pm d\theta + \delta \left[\frac{\sin \theta}{yN \cos (\theta \pm \alpha)} \right] dx_{\pm} = 0$$
 (107)

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E. NUMERICAL METHOD OF CHARACTERISTICS

The characteristic equations and compatibility relations for steady two-dimensional rotational supersonic flow are presented in the

preceding section. Those equations are ordinary differential equations. They can be integrated numerically by the second-order accurate modified-Euler predictor-corrector method. Such a procedure is called the numerical method of characteristics. The details of its implementation are presented in Reference (9). That is the procedure employed to calculate the flowfield in the nozzle extension downstream of the oblique shock wave emanating from the junction of the basic nozzle and the nozzle extension.

The basic procedure is to start from the downstream side of the oblique shock wave. The left-running and right-running Mach lines are constructed by solving equation (106). The pressure P and flow angle θ are determined at the intersections of the Mach lines by solving equation (107). The streamlines passing through these intersection points are projected upstream to intersect diagonal lines joining upstream intersection points by solving equation (103). The density ρ and velocity V are determined at the intersection points by solving equations (104) and (105). Since equations (103) to (107) are coupled and nonlinear, they must be solved numerically and simultaneously. This procedure defines the interior point unit process.

When a left-running Mach line impinges on the wall of the nozzle extension, the boundary conditions of known wall location and slope must be applied. A right-running Mach line is reflected into the flowfield from that point of impingement. That procedure derines the wall point unit process.

This procedure is continued until the entire flowfield in the nozzle extension has been crisscrossed with left-running and right-

running Mach lines and the flow properties have been calculated at all of the intersections of the two families of Mach lines. A coarse schematic of the resulting flowfield is illustrated in Figure 19.

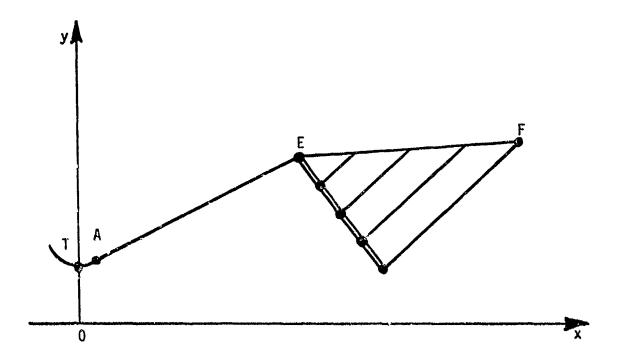


Figure 19. Mach line network in the nezzle extension.

SECTION IV

SCARFED NOZZLE PERFORMANCE MODEL

1. INTRODUCTION

The performance of a propulsive nozzle is specified by its mass flow rate, thrust vector, and moment vector. For conventional nozzles, the thrust vector lies along the nozzle axis, and the only thrust component is the axial component. In that case, all of the moments are zero. For a scarfed nozzle, side forces and moments are present. The present section presents the performance model for the scarfed nozzle configurations considered in this study (see Section III.2).

2. NOZZLE GEOMETRIC MODEL

The nozzle geometric model considered in the present investigation is illustrated schematically in Figure 20. The basic nozzle consists of a double circular arc throat following by a supersonic expansion contour. The supersonic expansion contour may be conical, quadratic, or specified by tabular data. The basic nozzle geometry and flowfield are axisymmetric.

The nozzle extension is a conical contour attached to the basic nozzle. The flow in the nozzle extension is assumed to be axisymmetric. That is, the wave (i.e., Mach line or shock wave) emanating from point E is assumed to fall downstream of the exit of the scarfed conical extension, as illustrated in Figure 20. In that case, the flowfield in the scarfed nozzle extension may be computed as though the nozzle extension were not scarfed (i.e., as if the dashed protion of the contour were present). The flowfield downstream of line EG, which is outside of the solid nozzle boundary, is obviously three-dimensional. However, the flowfield upstream of line EF is axisymmetric.

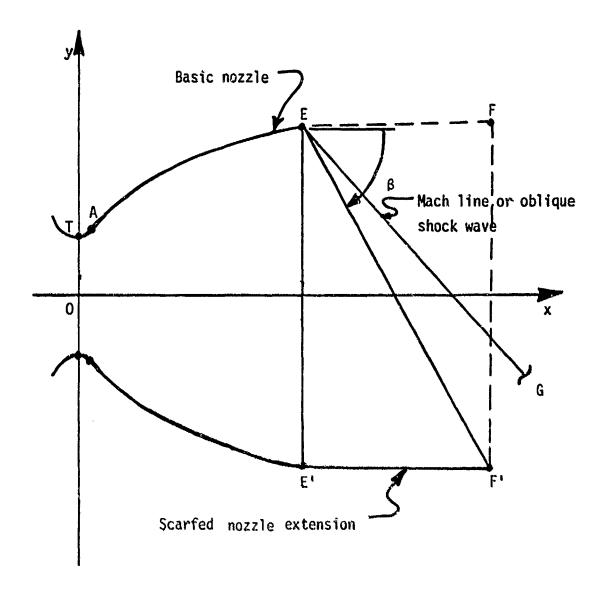


Figure 20. Scarfed nozzle geometric model.

3. BASIC NOZZLE PERFORMANCE EVALUATION

The basic nozzle is a conventional axisymmetric nozzle. The flowfield in the throat region is predicted by the linearized flow analysis presented in Section III.2. From that analysis, an initial-value line is obtained from which the supersonic flowfield can be calculated. In addition, (see Figure 10), the nozzle mass flow rate \mathfrak{m} , discharge coefficient C_D , and initial-value line thrust F_{IVL} are obtained.

The flowfield in the supersonic region is calculated by the method of characteristics for steady two-dimensional irrotational supersonic flow, as discussed in Section III.3. A schematic of the resulting Mach line network is presented in Figure 12. From that analysis, the pressure acting along the wall is known at each point where a Mach line intersects the wall. The thrust developed along the supersonic expansion contour is obtained by integrating (numerically) the axial component of the force developed by the pressure acting on the wall. Thus,

$$F_{SS} = \int_{y_{+}}^{y_{e}} (P - P_{a}) 2\pi y \, dy$$
 (108)

The total thrust developed by the basic nozzle, F_{N^*} is the sum of the thrust developed across the initial-value line, F_{IVL} , and the thrust developed along the supersonic contour, F_{SS} . Thus,

$$F_{N} = F_{IVL} + F_{SS} \tag{109}$$

4. SCARFED NOZZLE EXTENSION PERFORMANCE EVALUATION

The flowfield in the nozzle extension is calculated by the method of characteristics for steady two-dimensional rotational supersonic flow, as discussed in Section III.4. A schematic of the resulting Mach line network is presented in Figure 19. From that analysis, the pressure acting along the wall of the nozzle extension is known at each point where a Mach line intersects the wall. The thrust and moments developed by the scarfed nozzle extension are obtained by integrating (numerically) the differential force and moment components developed by the pressure acting on the wall.

In the present analysis, it is assumed that the nozzle extension is axisymmetric (in fact, conical) and that the exit of the nozzle extension is scarfed by a plane perpendicular to the xy plane which passes through point E on the top of the nozzle where scarfing begins and through point F on the bottom of the nozzle where scarfing ends. Consequently, the flowfield is symmetrical about the xy plane. Thus, as illustrated in Figure 21, only two force components exist, F_x and F_y , and only one moment component exists, M_z . The force components and moment are assumed to act at the center of the throat of the nozzle.

The geometric model employed to determine the force components and moment developed by the scarfed nozzle extension is illustrated in Figure 22. The geometry of the scarfed nozzle extension is specified by the location of point E (x_e, y_e) , the angle of the scarfed conical extension θ_f , and either the scarfing angle β or the location of point F (x_f, y_f) .

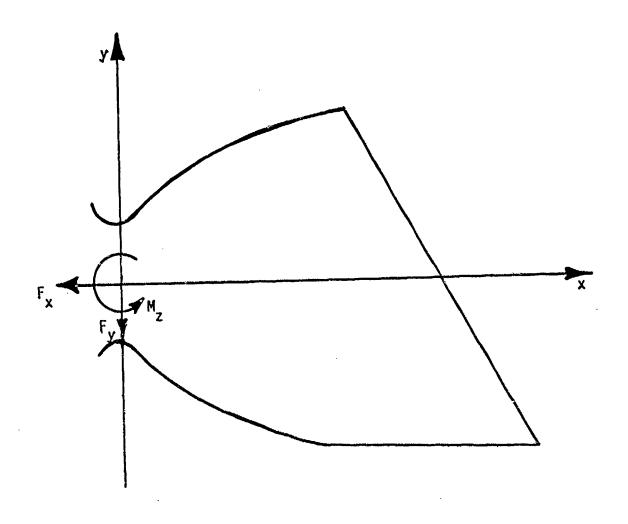
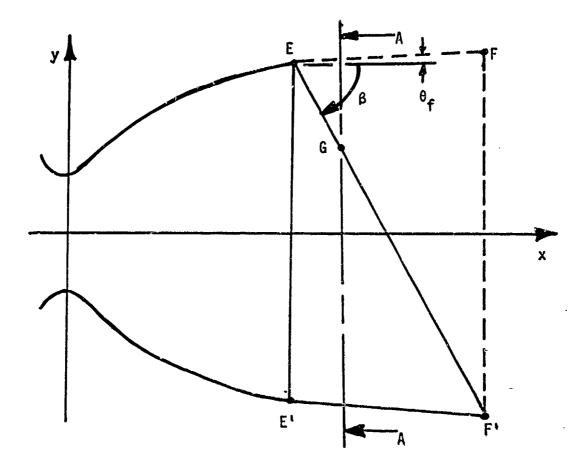
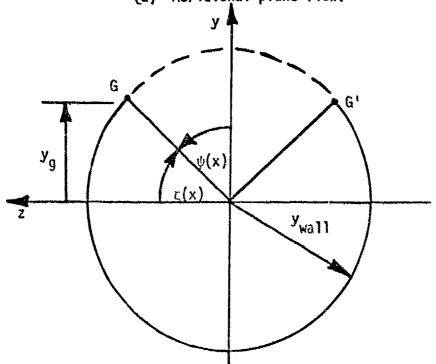


Figure 21. Scarfed nozzle performance model.



(a) Meridional plane view.



(b) Cross section A-A view.

Figure 22. Geometric model for force and moment evaluation.

The locations of points F and F are determined as follows. The equation of the scarfing line, line EF' is given by

$$\frac{y - y_e}{x - x_e} = -\tan \beta \tag{110}$$

where the angle β is in the range $0 \le \beta \le 90$ deg. Solving equation (110) for y gives

$$y = y_e + x_e \tan \beta - x \tan \beta = a - x \tan \beta$$
 (111)

where

$$a = y_e + x_e \tan \beta \tag{112}$$

The equation of line E'F' is given by

$$\frac{y - y_{e^1}}{x - x_{e^1}} = -\tan \theta_f \tag{113}$$

$$y = y_e^i + x_e^i \tan \theta_f - x \tan \theta_f = b - x \tan \theta_f$$
 (114)

where

$$b = y_{e'} + x_{e'} \tan \theta_{f'}$$
 (115)

Substituting (x_f, y_f) into equations (111) and (113) gives

$$y_{f'} = a - x_f \tan \beta \tag{116}$$

$$y_{f^{\dagger}} = b - x_{f} \tan \theta_{f} \tag{117}$$

Solving equations (116) and (117) simultaneously yields

$$x_{f} = \frac{a - b}{\tan \beta - \tan \theta_{f}}$$
 (118)

Equation (117) may then be solved for y_f and $y_f = -y_f$.

Point G illustrated in Figure 22 is the intersection point of a plane perpendicular to the nozzle axis with the edge of the scarfed extension, line EF. The location of point G corresponding to a specified value of $\mathbf{x}_{\mathbf{q}}$ is obtained by substituting $\mathbf{x}_{\mathbf{q}}$ into equation (111). Thus,

$$y_g = a - x_g \tan \beta \tag{119}$$

The angle $\zeta(x)$ illustrated in Figure 22 (b) is determined as follows.

$$\zeta(x) = \sin^{-1}(y_g/y_{\text{wall}}) \tag{120}$$

where y_{wall} is the radius of the scarfed conical extension at point G. The angle $\psi(x)$ illustrated in Figure 22(b) is determined from

$$\psi(x) = \frac{\pi}{2} - \zeta(x) \tag{121}$$

All of the geometric properties of the scarfed conical extension are now determined. Those properties are the coordinates (x_e, y_e) and (x_f, y_f) , the equation of the scarfing line EF' [equation (114)], and the angle $\psi(x)$.

The performance of the scarfed conical extension is determined by integrating the pressure forces acting along the wall from the point where the scafing begins, point E, to the end of the scarfed conical extension, point F. The geometric model for integrating the pressure forces is presented in Figure 23. The scarfed conical extension is assumed to be symmetrical about the xy plane. Consequently, the pressure forces acting on only one-half of the nozzle must be determined by integration. The total pressure force is then found by symmetry. For symmetry about the xy plane, no net component of the pressure force acts in the z-direction.

The components of the pressure force acting in the x and y directions are determined by defining a differential element of area $d\overline{A}$ having a magnitude dA and acting in the direction of the outward unit normal \overline{n} . The unit vector system \overline{i} , \overline{j} , \overline{k} in the Cartesian coordinate system xyz is illustrated in Figure 23. The radial unit vector \overline{i}_r , which is directed radially outward from the x-axis and lies in a plane perpendicular to the x-axis, is also illustrated in Figure 23. The element of area $d\overline{A}$ is given by

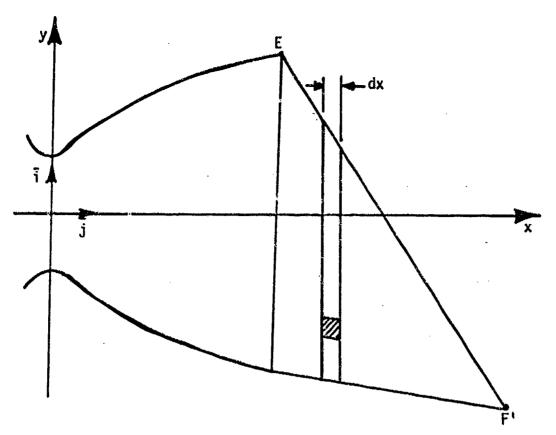
$$d\vec{A} = \vec{i}_r dA_r - \vec{i}_r dA_x \qquad (122)$$

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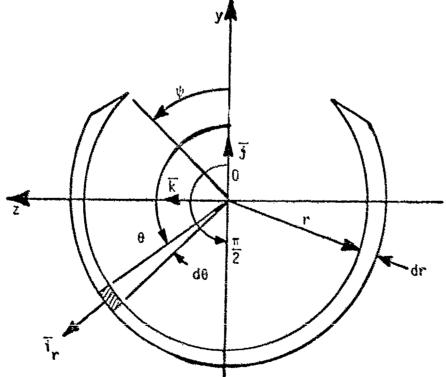
where

$$dA_{r} = r d\theta dx \tag{123}$$

$$dA_{v} = r d\theta dr \tag{124}$$



(a) Meridional plane view.



(b) Cross-section view.

Figure 23. Geometric model for thrust calculations.

In the Cartesian coordinate system, $\overline{\mathbf{I}}_{\mathbf{r}}$ is given by

$$\vec{T}_r = \vec{J} \cos \theta + \vec{k} \sin \theta$$
 (125)

where θ is measured counterclockwise from the positive y-axis. Combining equations (122) to (125) gives

$$d\overline{A} = -\overline{1} r d\theta dr + \overline{j} r \cos \theta d\theta dx + \overline{k} r \sin \theta d\theta dx \qquad (126)$$

The pressure force is given by

$$d\overline{F}_{p} = (P - P_{a}) d\overline{A}$$
 (127)

Due to the symmetry about the xy plane, the net component of the pressure force perpendicular to that plane is zero. Consequently, the third term on the right-hand side of equation (126) may be dropped from consideration, and equations (126) and (127) may be combined to yield

$$d\overline{F}_{p} = -\overline{i}(P - P_{a})r d\theta dr + \overline{j}(P - P_{a})r \cos\theta d\theta dx \qquad (128)$$

The total pressure force acting on the scarfed conical extension is obtained by integrating equation (128) between plane EE' and plane FF'. Thus,

$$\overline{F}_{p} = \overline{1} F_{x} + \overline{J} F_{y}$$
 (129)

where

$$F_{x} = -\int_{0}^{y} y_{e} \int_{0}^{\theta_{e}} (P - P_{a}) r d\theta dr$$
 (130)

$$F_{y} = \int_{x_{e}}^{x_{f}} \int_{\theta_{e}}^{\theta_{f}} (P - P_{a}) r \cos \theta \ d\theta \ dx$$
 (131)

The pressure P acting on the nozzle wall is a function of the local value of x, but it is not a function ω , the angle θ at a given axial location. Consequently, at a given axial location, equations (130) and (131) may be integrated over the range $\psi \leq \theta \leq \pi$, where ψ is a function of the axial location. Thus,

$$F_{x} = -2 \int_{y_{e}}^{y_{f}} (\pi - \psi) (P - P_{a}) r dr$$
 (132)

$$F_y = -2 \int_{x_e}^{x_f} (P - P_a) r \sin \psi dx$$
 (133)

The factor 2 appearing in equations (132) and (133) accounts for the two halves of the nozzle, since θ is integrated from ψ to π .

In equations (132) and (133), the pressure P is a function of the axial location. Consequently, those two equations must be integrated numerically. At any given axial location, the integral of equation (132) over the radial increment y_{i-1} to y_i and the integral of equation (133) over the axial increment x_{i-1} to x_i are given by

$$(\Delta F_{\mathbf{x}})_{\mathbf{i}} = -2\overline{P} \overline{y} (\pi - \overline{\psi}) (y_{\mathbf{i}} - y_{\mathbf{i}-1})$$
 (134)

$$(\Delta F_y)_i = -2\overline{P} \, y \, \sin \overline{\psi} \, (x_i - x_{i-1}) \tag{135}$$

where the overbar denotes average values over the increment i-1 to i, defined as

$$\overline{P} = 0.5(P_i + P_{i-1}) - P_a$$
 (136)

$$\overline{y} = 0.5(y_i + y_{i-1})$$
 (137)

$$\overline{\Psi} = 0.5(\psi_i + \psi_{i-1}) \tag{138}$$

The negative signs in equations (134) and (135) indicate that F_χ and F_y act in the negative directions relative to the coordinate system illustrated in Figure 23.

The total pressure force acting on the scarfed conical extension, $F_{SCE}, \mbox{ is obtained by summing equations (134) and (135) over the range} \\ x_e \mbox{ to } x_f. \mbox{ Thus,}$

$$F_{x, SCE} = \sum_{i=2}^{N} (\Delta F_{x})_{i}$$
 (139)

$$F_{y, SCE} = \sum_{i=2}^{N} (\Delta F_{y})_{i}$$
 (140)

where i = 1 denotes point E, i = 2 denotes the first wall point downstream of point E, and i = N denotes point F'.

5. SCARFED NOZZLE PERFORMANCE EVALUATION

The total axial force acting on the entire nozzle is obtained by adding the force given by equation (139) to the axial force F_N [see equation (109)] developed by the portion of the nozzle upstream of the conical scarfed extension. Thus,

$$F_{x} = F_{IVL} + F_{SS} + F_{x, SCE}$$
 (141)

The total side force acting on the entire nozzle is simply the side force acting on the scarfed conical extension, given by equation (140). Thus,

$$F_{y} = F_{y, SCE}$$
 (142)

The scarfed nozzle axial and side specific impulses are given by

$$(I_{sp})_{x} = \frac{|F_{x}|}{m} \tag{143}$$

$$\left(I_{sp}\right)_{y} = \frac{|F_{y}|}{m} \tag{144}$$

The geometric relationship between the nozzle axis and the missile axis is illustrated in Figure 24. The missile coordinate system is denoted by X,Y. The nozzle coordinate system x,y is assumed to lie in the meridional plane through the missile axis, that is, in the XY plane. The forces acting on the missile are thus

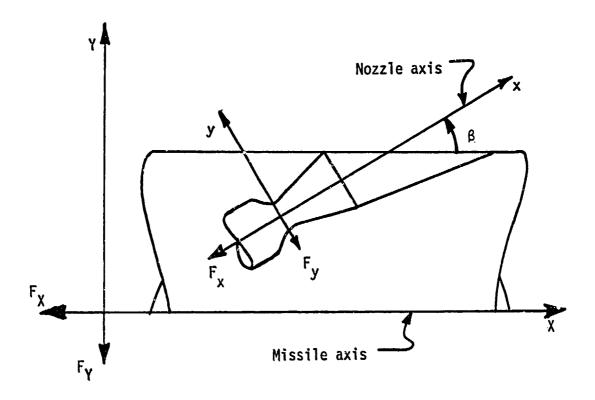


Figure 24. Relationship between nozzle axis and missile axis.

$$F_{\chi} = F_{\chi} \cos \beta - F_{y} \sin \beta \tag{145}$$

$$F_{Y} = F_{X} \sin \beta + F_{y} \cos \beta \tag{146}$$

The missile specific impulse is given by

$$(I_{sp})_{\chi} = \frac{|F_{\chi}|}{m} \tag{147}$$

$$(I_{sp})_{\gamma} = \frac{|F_{\gamma}|}{ii} \tag{148}$$

In the present analysis, the major item of interest is the axial thrust delivered by the rocket motor to the missile, F_{χ} , and how that thrust is related to the axial thrust developed by the nozzle, F_{χ} . It is assumed that several rocket motors (at least two) are arranged symmetrically around the circumference of the missile, so that all side forces and turning moments exactly cancel. If that is not the case, the analysis may be easily extended to determine the turning moments associated with the scarfed nozzle. That is not done, however, in the present analysis.

The final performance parameter of interest is the ratio η of the axial thrust delivered to the missile, F_χ , to the axial thrust that would be generated by the nozzle if the scarfed conical extension were not scarfed and if it were aligned along the axis of the missile. That value of thrust is determined by calculating the axial force developed by the conical extension without scaring, F_{CE} , and adding that value to the axial force acting on the basic nozzle. The axial force F_{CE} is

determined by applying equation (108) from point E to point F. Thus,

$$F_{CE} = \int_{y_e}^{y_f} (P - P_a) 2\pi y \, dy$$
 (149)

Note that if the angle θ_f of the scarfed conical extension is zero, then F_{CE} is zero. The total reference axial thrust of the unscarfed scarfed nozzle F_{Ref} is given by

$$F_{Ref} = F_{IVL} + F_{SS} + F_{CE}$$
 (150)

Thus, the ratio n is given by

$$\eta = \frac{F_{x}}{F_{Ref}} \tag{151}$$

The factor η may be regarded as the efficiency of the scarfed nozzle.

From Figure 24, it is obvious that the scarfed nozzle exit flow area lies on the missile skin. Thus, the pressure forces acting on that area are normal to the missile axis and do not contribute to the missile axial thrust. The missile axial thrust depends on the momentum flux crossing the nozzle exit area, which depends on the mozzle mass flow rate and nozzle exit velocity. The exit velocity is independent of the pressure level; it depends on the mozzle geometry, the gas thermodynamic model (i.e., γ and R), and the stagnation temperature $\Gamma_{\rm t}$. The mass flow rate depends linearly on the pressure level, that is, the stagnation pressure. Consequently, the missile axial thrust and nozzle mass flow rate both depend linearly on the stagnation pressure. However, the

ratio of those two quantities, the missile axial specific impulse, is independent of both the stagnation pressure and the atmospheric pressure. Consequently, missile axial specific impulse is the most meaningful performance parameter for a scarfed propulsive nozzle.

SECTION V

OVERALL NUMERICAL ALGORITHM

1. INTRODUCTION

The objective of the present investigation was to develop a procedure for predicting the performance of a particular class of scarfed propulsive nozzles. That class of nozzles consists of a conventional basic nozzle which has a double circular arc throat contour followed by a conical, quadratic, or tabular supersonic contour. At the end of the basic nozzle, a scarfed conical extension is attached. Figure 25 illustrates this type of scarfed propulsive nozzle. A detailed discussion of the nozzle geometric model is presented in Section II.

The flowfield in the transonic throat region of the nozzle is predicted by a linearized flow analysis. The flowfield in the supersonic portion of the basic nozzle is predicted by the method of characteristics for steady two-dimensional irrotational supersonic flow. The transition between the basic nozzle flowfield and the nozzle extension flowfield occurs across an oblique shock wave that emanates from the junction of the basic nozzle wall and the nozzle extension. The flowfield in the nozzle extension is oredicted by the method of characteristics for steady two-dimensional rotational supersonic flow. A detailed discussion of the aforementioned flow mode's is presented in Section III.

The performance of the scarfed propulsive nozzle is obtained by integrating the flowfield properties to determine the overall

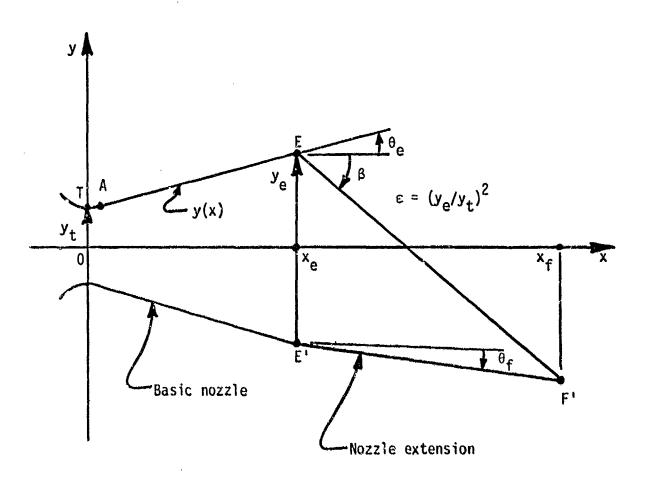


Figure 25. Scarfed nozzle geometric model.

performance parameters (i.e., mass flow rate, force components, and turning moments). The mass flow rate is obtained by integrating across the transonic initial-value line. The thrust of the basic nozzle is determined by integrating the momentum flux and pressure forces across the initial-value line, and adding to that value the integral of the pressure forces along the nozzle wall. The thrust of the nozzle extension is obtained by integrating the pressure forces along the wall of the nozzle extension. The total nozzle thrust components are determined by adding the thrust components of the basic nozzle and the nozzle extension. The thrust delivered to the missile is found by calculating the component of the nozzle thrust components in the direction of the missile axis. Details of the performance model are discussed in Section IV.

The objective of the present section is to present the logic employed to combine all of the individual analyses presented in Sections II to IV into an overall numerical algorithm for predicting the performance of scarfed propulsive nozzles.

A computer program has been written for implementing the analysis developed in this investigation. A discussion of that program is presented in Section VI. In the present section describing the overall numerical algorithm, references are made to the particular program routines that implement the various aspects of the overall numerical algorithm. Thus, program MAINOO controls the overall logic flow by calling, in sequence, program MAINOO to read in the input data.

program MAIN20 to construct the flowfield in the basic nozzle, program MAIN30 to construct the flowfield in the nozzle extension, and program MAIN40 to calculate the performance of the scarfed nozzle.

2. FLOWFIELD CONSTRUCTION LOGIC

Five modes of flowfield Mach line network construction are considered in the program. The first four modes all use the right-running Mach line initial expansion flowfield illustrated in Figure 26. The fifth mode uses the left-running Mach line initial expansion flowfield illustrated in Figure 27. In both figures, line TT' is a supersonic initial-value line obtained from a transonic flow analysis. Right-running Mach lines are initiated from each point on line TT', starting adjacent to the nozzle axis, and propagated downward until they intersect the nozzle axis. Right-running Mach line TT', emanating from the nozzle throat point, point T, defines the downstream extent of the flowfield from the initial-value line. This flowfield is identical in Figures 26 and 27.

The initial expansion flowfield illustrated in Figure 26 is obtained by initiating right-running Mach lines from prespecified points along the nozzle throat downstream circular arc contour, contour TA. These right-running Mach lines are continued until they intersect the nozzle axis. Right-running Mach line AA' emanating from the throat attachment point, point A, defines the downstream extent of the flowfield from the initial expansion contour.

The initial expansion flowfield illustrated in Figure 27 is obtained by initiating left-running Mach lines from right-running Mach line TT'.

These left-running Mach lines are continued until they intersect the nozzle throat downstream circular arc contour, contour TA. Left-running Mach line A'A emanating from point A' on right-running Mach line TT'

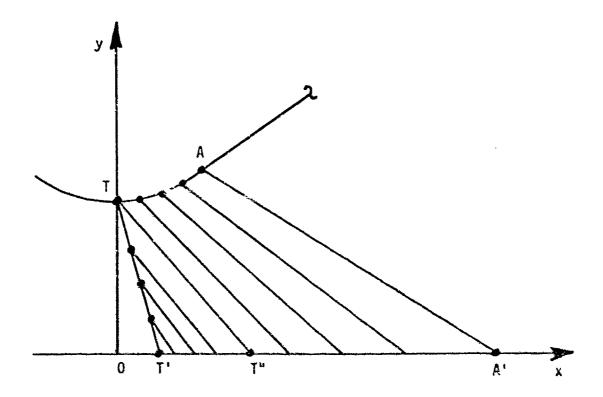


Figure 26. Right-running Mach line initial-expansion flowfield.

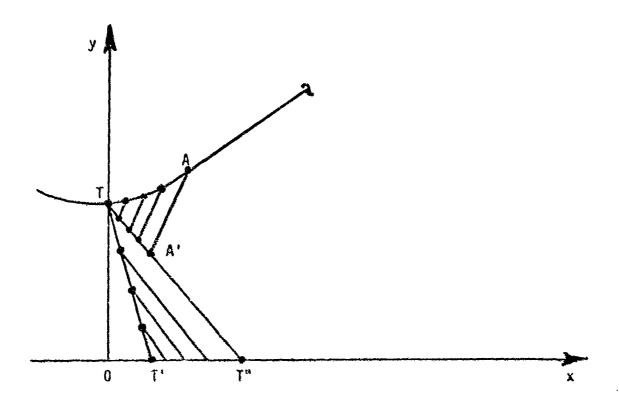


Figure 27. Left-running Mach line initial-expansion flowfield.

and passing through the throat attachment point, point A, defines the downstream extent of the flowfield along the initial expansion contour.

Five different modes of flowfield construction are considered in the program. Those modes of operation are described in the following paragraphs.

Mode 1. The first mode analyzes the flowfield in a nozzle without an extension by constructing a network of right-running Mach lines, as illustrated in Figure 28. The initial expansion flowfield illustrated in Figure 26 is first constructed. Right-running Mach lines are then initiated from points along the nozzle supersonic contour, contour AE. These right-running Mach lines are continued until they intersect the nozzle axis. Right-running Mach line EE', emanating from the nozzle exit lip point, point E, defines the downstream extent of the nozzle flowfield.

Mode 2. This mode analyzes the flowfield in a nozzle without an extension by constructing a network of left-running Mach lines, as illustrated in Figure 29. The initial expansion flowfield illustrated in Figure 26 is first constructed. Left-running Mach lines are then originated from line AA'. These left-running Mach lines are continued until they intersect the nozzle wall, contour AE. Left-running Mach line EE's which passes through the nozzle exit lip point, point E, defines the downstream extent of the nozzle flowfield.

Mode 3. This mode analyzes the flowfield in a sozzle without an extension in which an embedded right-running oblique shock wave is detected and tracked. The flowfield is constructed along a network of left-running Mach lines, as illustrated in Figure 30. The initial

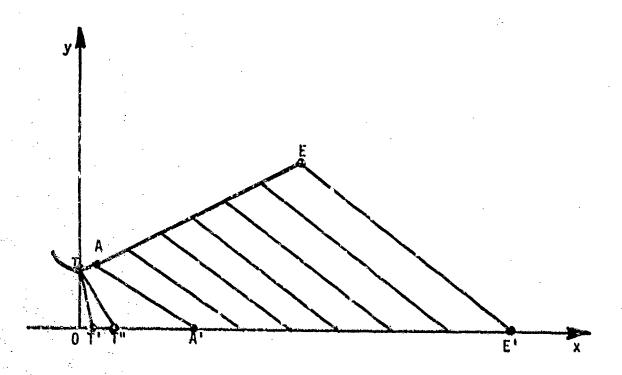


Figure 28. Node 1 Mach line network.

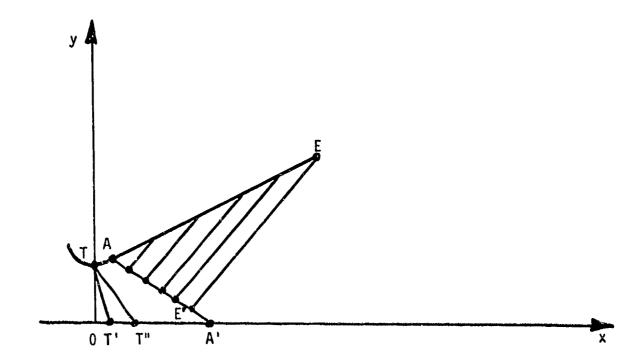


Figure 29. Mode 2 Nach line network.

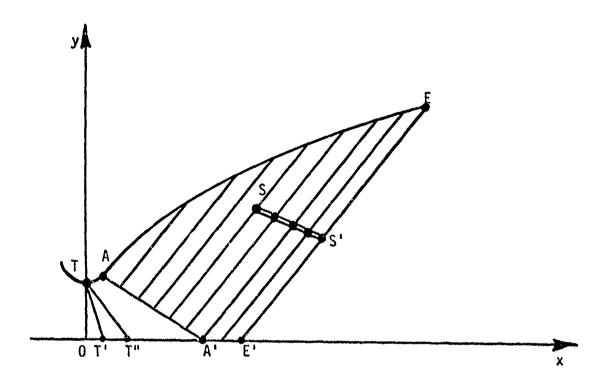


Figure 30. Mode 3 Mach line network.

expansion flowfield illustrated in Figure 26 is first constructed.

Left-running Mach lines are then originated from line AA'. These leftrunning Mach lines are continued until they intersect the nozzle wall,
contour AE. The point of initiation, point S, of the right-running
oblique shock wave SS' must be specified in terms of its (I,J) characteristic coordinates. Subsequent left-running Mach lines pass through
the oblique shock wave, which is illustrated as a double line in Figure
30. The location and properties of the oblique shock wave are determined simultaneously with the construction of the left-running Mach
line network. Left-running Mach line E'E, which passes through the
nozzle exit lip point, point E, defines the downstream extent of the
nozzle flowfield.

Mode 4. This mode analyzes the flowfield in a nozzle with a scarfed nozzle extension in which an attached right-running oblique shock wave emanates from the junction between the basic nozzle and the nozzle extension, point E. The flowfield is constructed along a network of left-running Mach lines, as illustrated in Figure 31. The initial expansion flowfield illustrated in Figure 26 is first constructed. Left-running Mach lines are then originated from line AA'. These left-running Mach lines are continued until they intersect the nozzle wall, contour AE. Left-running Mach line E'E, which passes through the nozzle exit lip point, point E, defines the flow properties at point E. An attached right-running oblique shock wave is originated from point E. Subsequent left-running Mach lines pass through the oblique shock wave. The location and properties of the oblique shock wave are determined simultaneously with the construction of the left-running Mach line

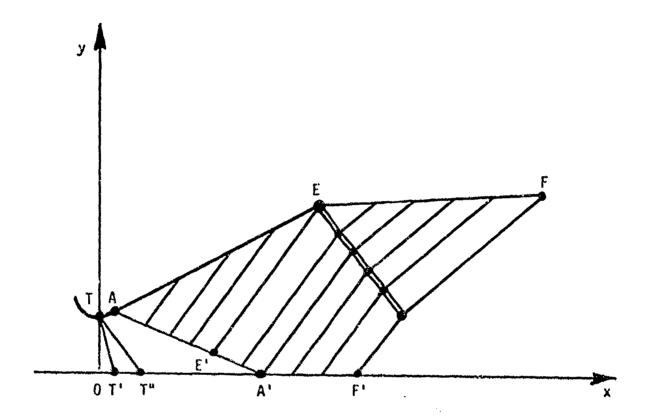


Figure 31. Node 4 Nach line network.

network. Left-running Mach line F'F, which passes through point F, the exit lip point of the nozzle extension, defines the downstream extent of the flowfield in the scarfed nozzle extension.

Mode 5. This mode is identical to the fourth mode described in the previous paragraph, except that the initial expansion flowfield illustrated in Figure 27 is employed instead of the initial expansion flowfield illustrated in Figure 26. The remainder of the flowfield is constructed exactly as described in the previous paragraph. The resulting flowfield is illustrated in Figure 32.

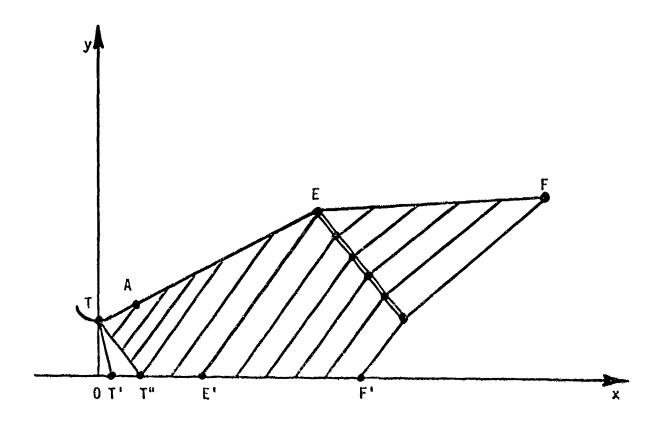


Figure 32. Mode 5 Mach line network.

2. PROBLEM SPECIFICATIONS

The first step in the procedure is the specification of the problem to be analyzed. The nozzle geometric model must be specified as described in Section II and illustrated in Figure 25. The nozzle throat is specified by the throat radius y_t , the throat upstream radius of curvature ρ_{tu} which is employed in the transonic flow analysis, and the throat downstream radius of curvature ρ_{td} which controls the initial supersonic expansion. The x location of the nozzle throat is chosen to be 0.0.

The basic nozzle supersonic contour, y(x), attaches smoothly to the throat contour at the attachment angle θ_a . The length of the basic nozzle is x_e . If the supersonic expansion contour is conical, the cone angle must be equal to θ_a . A conical nozzle of length x_e is completely specified by any one of the variables θ_a , y_e , or ϵ . If the supersonic contour is quadratic, θ_a and x_e must be specified. Then, one of the variables θ_e , y_e , or ϵ completes the specification of the quadratic nozzle. For the tabular nozzle, a set of (x,y) pairs must span the region from point A to point E. The first point in the table must be the first point downstream of point A. The last point in the table becomes point E. The basic nozzle geometry is described in the computer program in subroutine B&UNDYE.

The conical extension is specified completely by the values of $\mathbf{x_f}$ and $\mathbf{\theta_f}$. The geometry of the conical extension is described in the computer program in subroutine BOUNDYW.

The gas thermodynamic model is specified by the gas specific heat ratio γ and the gas constant R. Subroutines THERMOI and THERMOR

evaluate the equations of state for a thermally and calorically perfect gas flowing in an irrotational and rotational flowfield, respectively.

The nozzle operating conditions are specified by the stagnation pressure $P_{\bf t}$, the stagnation temperature $T_{\bf t}$, and the ambient pressure $P_{\bf a}$.

The parameters described in the above paragraphs completely specify a specific nozzle and its operating conditions. Numerous other parameters must be specified to enable a numerical solution of the flowfield. Those parameters are described in the following paragraphs as they are encountered.

The input data are read into the computer program by subroutine INPUT in overlay LINK10.

3. Initial-value line

The first step in the numerical solution is the determination of a supersonic initial-value line spanning the nezzle throat region from which the method of characteristics solution for the supersonic flow-field can be initiated. Two options for obtaining a supersonic initial-value line are contained in the computer program. The first option is an internally generated initial-value line. The second option is to input a tabular initial-value line obtained from any other source.

A. Internally Generated Initial-Value Line

The internally generated initial-value line is based on Kliegel's analysis [Reference (2)]. That transonic flowfield model is described in Section III.2. The only additional information required to specify the location of the initial-value line are the number of points desired and the spacing of those points. The first point is the point on the nozzle axis $(\varepsilon, 0.0)$, and the last point is the nozzle throat point $(0.0, y_*)$.

The total number of points on the initial-value line is HI. These NI points may be spaced uniformly along the initial-value line, or they may be spaced according to a geometric progression in λy so that the ratio of the final Δy adjacent to the nozzle wall to the initial Δy adjacent to the nozzle axis is given by the parameter DYRATIO. This latter option, spacing according to a geometric progression, is useful when either throat radius of curvature, ρ_{tu} or $\rho_{tt'}$ is small compared to the throat radius y_t , so that large gradients in flow properties occur adjacent to the nozzle wall in the throat region.

Figure 33 illustrates a typical initial-value line determined by the present algorithm.

B. Tabular Initial-Value Line

The supersonic initial-value line may be determined from any external source and read into the computer program in tabular form. The order of the tabular data must be consistent with the order in which the data are stored and used by the remainder of the program. The first tabular point must be the nozzle axis point, point T. and the last tabular point must be the nozzle wall point, point T. The number of points and their spacing is arbitrary.

The tabular initial-value line is specified by the number of points on the line, NI. the location of the points (i.e., x and y), and the Mach number A and flow argle 6 at each point. From this information, the remainder of the flow properties at each point are calculated internally (i.e., P, p. T, Y, u, and v).

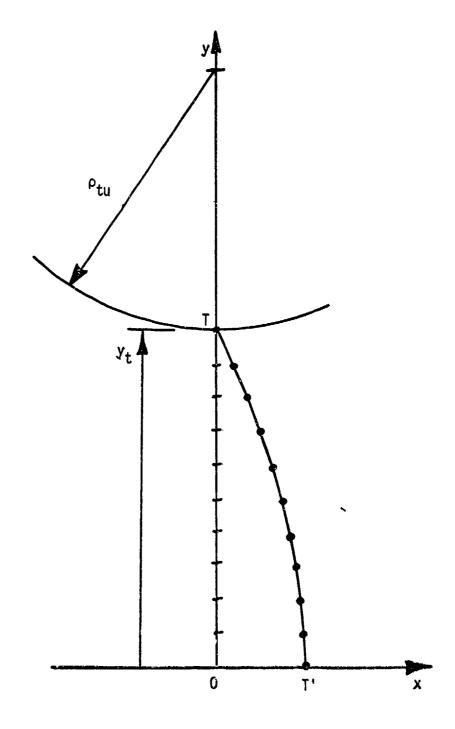


Figure 33. Supersonic initial-value line.

4. FLOWFIELD FROM THE INITIAL-VALUE LINE

The flowfield model in the basic nozzle is described in Section III.3. The numerical solution for this flowfield is obtained by applying the method of characteristics for steady two-dimensional irrotational supersonic flow. That procedure is implemented by constructing left-running and right-running Mach lines, starting from the initial-value line, until they crisscross the entire supersonic flowfield. The solution of the flowfield from the initial-value line is described in this section.

Figure 34 illustrates schematically the application of the unit process for an interior point at the first two points located on the form portion of the initial-value line, for determining the location of and the flow properties at the downstream point of intersection of the two Mach lines emanating from the two initial-value line points. Figure 35 illustrates schematically the application of the unit process for an axis point, for determining the location of and the flow properties at the downstream point of intersection of the right-running Mach line from the interior point and the nozzle axis.

The foregoing procedure is repeated from the next two points on the initial-value line, thus extending the corresponding right-running Mach line to the nozzle axis. The procedure is repeated until the complete region determined by the initial-value line has been determined. Figure 36 illustrates the resulting network of lach lines emanating from the initial-value line. The flowfield from the initial-value line is now complete. To continue the solution further, the wall boundary conditions must be employed.

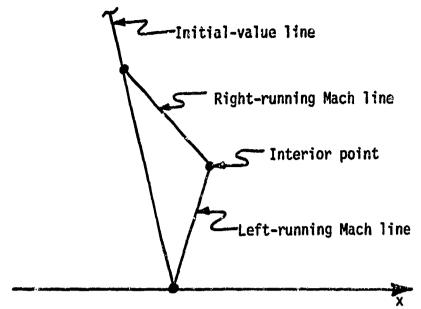


Figure 34. Application of the unit process for an interior point from the first two points on the initial-value line.

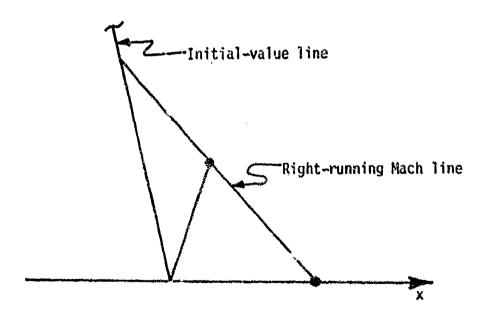


Figure 35. Application of the unit process for an axis point.

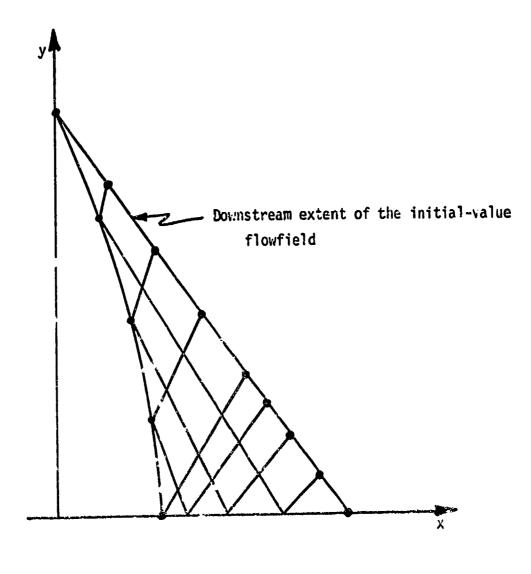


Figure 36. Extent of the initial-value problem.

In the computer program, the logic described in this section for determining the flowfield from the initial-value line is implemented by subroutine MOCIVL in overlay LINK2O.

5. FLOWFIELD FROM THE CIRCULAR ARC THROAT

The flowfield from the circular arc throat is determined by prespecifying the number and location of points along the circular arc, and then applying the method of characteristics to determine the flow properties at those points. Two different Mach line network construction procedures are available for determining the flow properties along the initial expansion contour.

The first Mach line network construction procedure constructs rightrunning Mach lines from the prespecified points along the circular arc
throat contour. These right-running Mach lines are propagated across the
flowfield to the nozzle axis just as the right-running Mach lines from the
initial-value line points were propagated to the nozzle axis (see Section
V.4). This procedure is continued until the right-running Mach line from
point A, the attachment point between the circular arc throat and the
supersonic expansion contour, has been constructed. This Mach line network, illustrated in Figure 26, is used for MODE = 1 to 4, as discussed
in Section V.1 and illustrated in Figures 28 to 31.

The second Mach line network construction procedure constructs leftrunning Mach lines from the right-running Mach line at the downstream
extent of the flowfield from the initial-value line, illustrated in
Figure 36. These left-running Mach lines are propagated across the flowfield to the circular arc throat contour. This procedure is continued
until a left-running Mach passes downstream of the last prespecified wall
point on the circular arc throat contour, point A. This Mach line network,
illustrated in Figure 27, is used for NODE = 5, as discussed in Section

V.1 and illustrated in Figure 32.

The total turning angle along the circular arc throat is θ_a . A number of points are prespecified along the circular arc contour. The first point is a point just downstream of the throat point (i.e., point T), and the last point is the attachment point between the circular arc throat and the supersonic contour (i.e., point A). The only additional information required to specify the location of these points are the number of points desired and the spacing of those points.

The total number of points on the circular arc throat (not including the throat point, point T) is NT. These NT points may be spaced along the circular arc in equal angular increments, or they may be spaced according to a geometric progression in $\Delta\theta$ so that the ratio of the final $\Delta\theta$ adjacent to point A to the initial $\Delta\theta$ adjacent to point T is given by the parameters DARATIO. This latter option, spacing according to a geometric progression, is useful when the throat downstream radius of curvature ρ_{td} is small compared to y_t , so that large gradients in flow properties occur adjacent to the nozzle wall along the circular arc contour. Figure 37 illustrates a typical distribution of points along the throat downstream circular arc contour.

Figure 38 illustrates schematically the application of the inverse wall point unit process to determine the first prespecified point on the threat downstream circular arc contour. A left-running Mach line is projected rearward from the prespecified wall point to intersect the previous right-running Mach line, as illustrated in Figure 38. If the point of intersection falls below the second point, then the point of intersection is assumed to fall between the second and third points, as

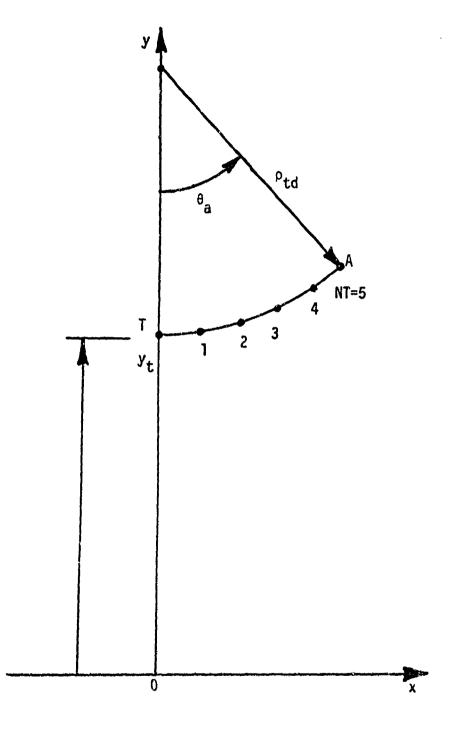
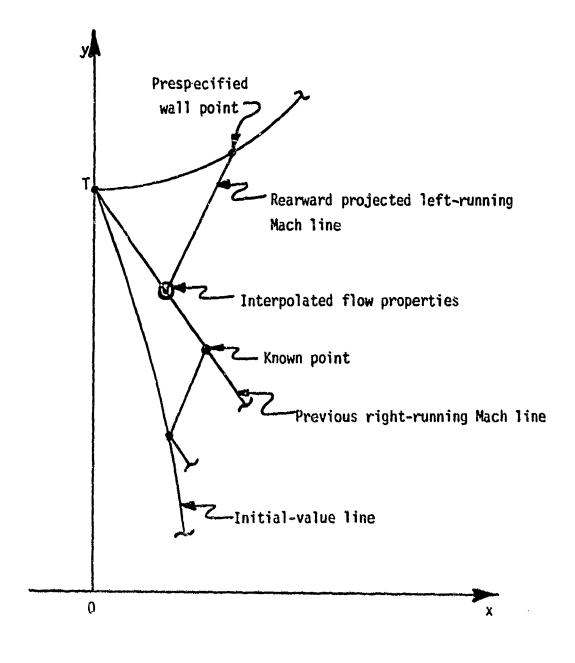


Figure 37. Distribution of points along the throat downstream circular arc contour.



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Figure 38. Application of the inverse wall point unit process.

illustrated in Figure 39. This procedure is continued until the two points bracketing the rearward projected left-running Mach line are located. In most cases, these two points are the first two points on the previous right-running Mach line. The flow properties at the point of intersection are then calculated by linear interpolation. The flow properties at the prespecified wall point are then found by applying the compatibility relation along the left-running Mach line and the solid wall boundary conditions.

The above procedure for determining the first inverse wall point is the same for both of the Mach line network construction procedures described at the beginning of this section. From here on, however, the two procedures are different. They are described in the next two sections.

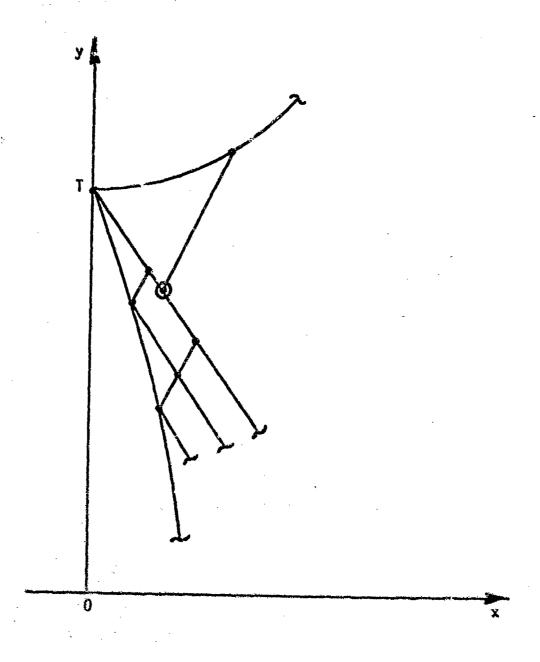


Figure 39. Modification of the inverse wall point unit process.

5. RIGHT-RUNNING MACH LINE NETWORK FROM THE CIRCULAR ARC THROAT

Once the solution for the inverse wall point has been obtained, as illustrated in Figure 38, a right-running Mach line is originated from the inverse wal! point and continued until it intersects the nozzle axis, as described in Section V.4 for the right-running Mach lines emanating from points on the initial-value line. The above procedure is repeated at each successive prespecified wall point on the nozzle circular arc throat contour until the region of the flowfield determined by that portion of the nozzle contour has been determined. Figure 40 illustrates the resulting network of Mach lines emanating from the nozzle circular arc throat contour. In the computer program, the logic described in this section is implemented in subroutine MOCARCR in overlay LINK20.

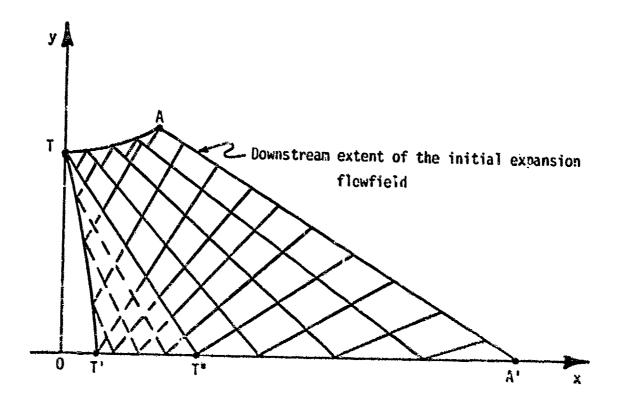


Figure 40. Extent of the flowfield determined by the nozzle circular arc throat contour.

7. CHARACTERISTIC COORDINATE SYSTEM

By this point, it should be obvious that the numerical solution is being obtained along the network of left-running and right-running Mach lines that crisscross the flowfield. One of the major advantages of the numerical method of characteristics is that numerical information propagates along the Mach lines, which are the actual paths of propagation of physical information. Consequently, the Mach line network provides an instant view of the domain of dependence and range of influence of each and every point in the flowfield. Hence, it is very useful to be able to identify specific individual Mach lines in a flowfield.

The identification of individual Mach lines is accomplished in the numerical solution by employing a characteristic coordinate system, as illustrated in Figure 41. The characteristic coordinate system consists of the right-running Mach lines, called I characteristics, and the left-running Mach lines, called J characteristics. The characteristic coordinate system is obviously not an orthogonal coordinate system, but a curvilinear coordinate system composed of the right-running and left-running Mach lines.

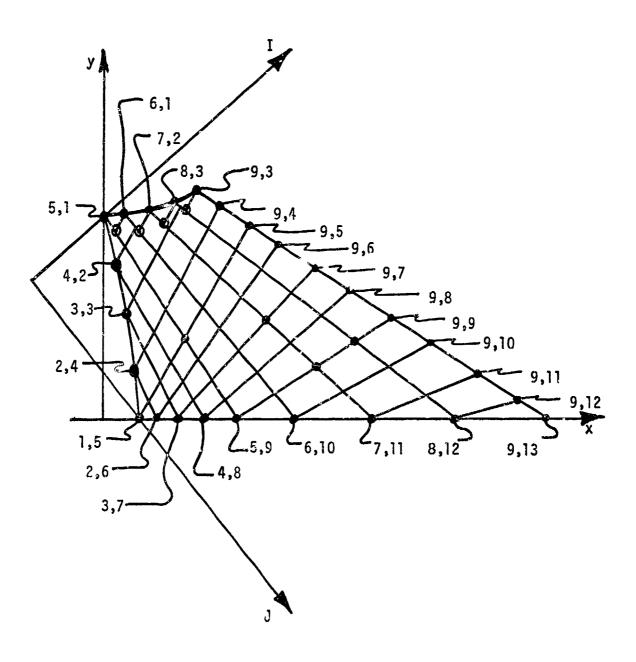


Figure 41. Characteristic coordinate system.

In the characteristic coordinate system, the right-running Mach lines are numbered starting at one for the point on the initial-value line at the nozzle axis, and increasing by one for each succeeding right-running Mach line. The left-running Mach lines are numbered starting at one for the point on the initial-value line at the nozzle throat, and increasing by one for each succeeding left-running Mach line. By following a constant value of the I or J characteristic coordinate through the flowfield, a unique right-running or left-running Mach line can be traced from its point of origin to its final point in the flowfield. The influence of initial-value line points and wall points on the flowfield can be easily tracked in this manner.

To illustrate the characteristic coordinate system further, several special points are considered. The coordinates of the point on the initial-value line at the nozzle axis are (1, NI), and the coordinates of the point at the nozzle throat are (NI, 1). The coordinates of all of the initial-value line points are specified on Figure 41, where NI = 5. Note that points (1,1) to (1.4), (2,1) to (2,3), (3,1), (3,2), and (4,1) are not defined.

The first prespecified wall point has an I value of NI + 1. The corresponding J value is that of the left-running Mach line just above the rearward projected left-running Mach line through the wall point. Consequently, the J value of a prespecified wall point is the same as the J value of the previous wall point when the rearward projected left-running Mach line intersects the previous right-running Mach line between the first two points on that Mach line. When that intersection point falls between the second and third point, the J value of the

prespecified wall point is one greater than the J value of the previous prespecified wall point. The coordinates of the prespecified wall points are presented on Figure 41. Clearly, as left-running Mach lines intersect the wall and dissappear from the Mach line network, the characteristic coordinate system reflects this situation by the increasing J values of the wall points.

On the nozzle axis, right-running Mach lines from the initialvalue line or the nozzle wall intersect the axis and left-running Mach
lines are initiated from the intersection point and propagated into the
flowfield. Thus, new left-running characteristics having increasing J
values are generated from the nozzle axis. The coordinates of the axis
points are specified on Figure 41.

The characteristic coordinates of all of the interior points are simply the I and J values of the right-running and left-running Mach lines that intersect at each point. The coordinates of the interior points along the right-running Mach line from the last prespecified wall point illustrated in Figure 41 are specified to illustrate the coordinates for interior points. The coordinates of several other interior points are specified on Figure 41 for illustration.

8. LEFT-RUNNING MACH LINE NETWORK ADJACENT TO THE CIRCULAR ARC THROAT

The first step in determining a left-running Mach line network along the initial expansion contour is the same as the procedure illustrated in Figure 38. After determining the first inverse wall point, the right-running Mach line emanating from that point is extended into the flowfield to find its point of intersection with the left-running Mach line just below the interpolated point on the previous right-running Mach line. That left-running Mach line becomes the control left-running Mach line in the following procedure.

So far, this procedure is identical to the procedure presented in Section V.6. However, instead of continuing the right-running Mach line to the nozzle axis as desribed in Section V.6, the next inverse wall point is determined, as illustrated in Figure 42. The right-running Mach line emanating from that inverse wall point is then extended into the flowfield to intersect the control left-running Mach.

This procedure is continued from succeeding prespecified inverse wall points until the point of intersection of the rearward projected left-running Mach line from the next inverse wall point and the previous right-running Mach line lies below the control left-running Mach line, as illustrated in Figure 42. At that point, that particular control left-running Mach line is complete, and the entire procedure is repeated, starting from the next point on the last right-running Mach line at the downstream extent of the initial-value flowfield.

The above procedure is repeated from succeeding points on the rightrunning Mach line at the downstream extent of the initial-value flowfield

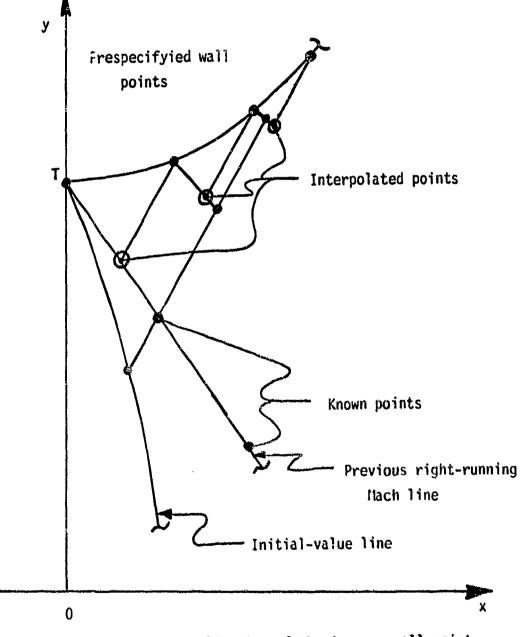


Figure 42. Modified application of the inverse wall point unit process.

until the last prespecified inverse wall point has been determined. The final control left-running Mach line is then extended forward, by the direct wall point unit process, to intersect the supersonic turning contour. Figure 43 illustrates the resulting network of Mach lines. This Mach line network is employed when analyzing a scarfed nozzle by MODE = 5. In the computer program, the logic described in this section is implemented in subroutine MOCARCL in overlay LINK2O.

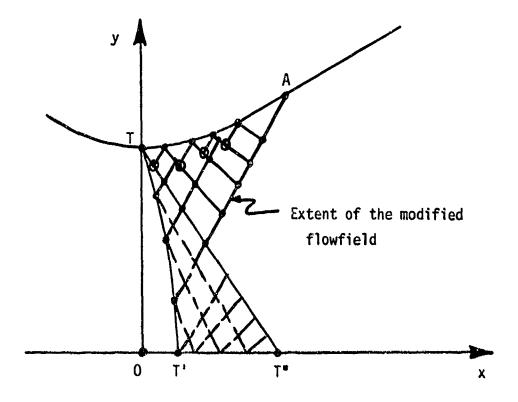


Figure 43. Extent of the modified flowfield determined by the nezzle circular arc throat contour.

9. RIGHT-RUNNING MACH LINE NETWORK IN THE BASIC NOZZLE

The flowfield from the supersonic contour (i.e., conical, quadratic, or tabular contour) is determined by extending a left-running Mach line from the second point on the right-running Mach line emanating from the previous wall point until it intersects the supersonic contour. The flow properties at the point of intersection are determined by the method of characteristics. A right-running Mach line is then emanated from the wall points. The right-running Mach line is propagated across the flowfield to the nozzle axis just as the right-running Mach lines from the initial-value line (see Section V.4) and right-running Mach lines from the throat circular arc contour (see Section V.6) were propagated to the nozzle axis. This procedure is continued until the intersection of the left-running Mach line from the previous rightrunning Mach line and the supersonic contour falls on the projection of the supersonic contour beyond the end point of the basic nozzle, point The flow properties at point E are determined by the inverse wall point method described in Section V.4. The final right-running Mach line from point E is propagated to the nozzle axis, thus completing the flowfield determined by the basic nozzle supersonic contour.

Figure 44 illustrates schematically the application of the direct wall point unit process to determine the first point on the supersonic contour just downstream of the attachment point, point A, where the throat circular arc contour joins the supersonic contour. The flow properties at the direct wall point are then found by applying the compatibility relation along the left-running Mach line and the solid wall boundary conditions. This procedure is implemented in the computer

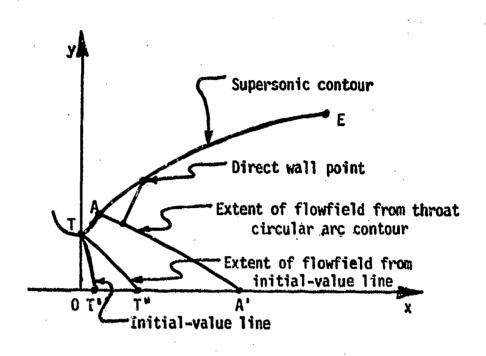


Figure 44. Application of the direct wall point unit process.

program by subroutine DRWALLI and BOUNDYE in overlay LINK20.

The complete flowfield along right-running Mach lines emanating from the supersonic contour of the basic nozzle is illustrated schematically in Figure 45. This Mach line network is employed when analyzing a basic nozzle (MODE = 1) or a basic nozzle with an embedded right-running oblique shock wave (MODE = 3). In the computer program, the logic described in this section is implemented in subroutine MOCRRC in overlay LINK2O.

The procedure described above is very straightforward in a contine uous flowfield. However, when compression waves are present they tend to coalesce, and if enough coalescence occurs, oblique shock waves are formed. The coalescence of Mach lines and the appearance of oblique shock waves complicates both the flowfield itself and the logic procedure described above for calculating the flowfield. In the present investigation, the assumption is made that no strong oblique shock waves occur in the basic nozzle.

When right-running Mach lines coalesce (i.e., cross), the original Mach line is retained and the new Mach line is terminated downstream of the point of crossing. This procedure is illustrated in Figure 46. Since the computational logic is based on following right-running Mach lines across the flowfield from the nozzle wall to the nozzle axis, each Mach line must reach the nozzle axis. The above manner of handling crossing Mach lines achieves that goal. The new Mach line thus consists of the new points calculated before the point of crossing and the old points downstream of the point of crossing. Thus, the new Mach line completely spans the nozzle flowfield from the nozzle wall to the nozzle axis. In

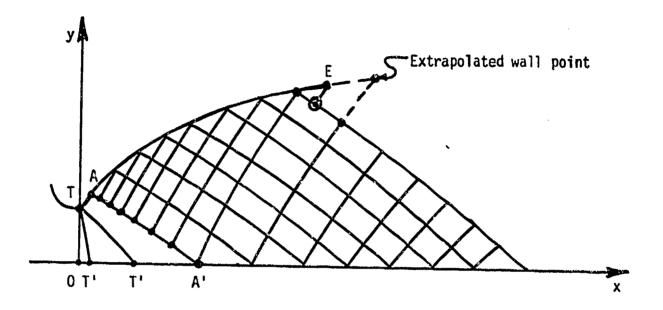
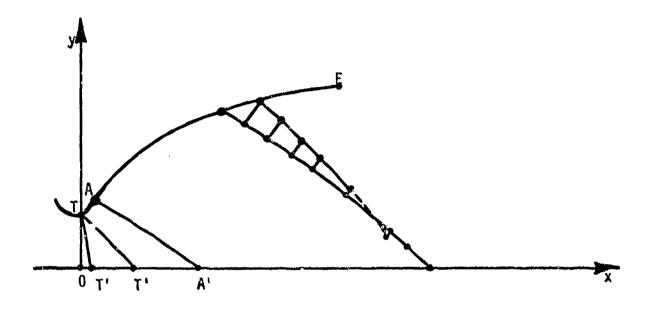


Figure 45. Extent of the right-running Mach line network along the basic nozzle contour.



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Figure 46. Crossing right-running Mach lines.

the computer program output, the points downstream of the point of crossing are written out as the solution for both right-running Mach lines.

When left-running Mach lines cross, the new Mach line is terminated. Since the computational logic is based on following right-running Mach lines across the flowfield instead of following left-running Mach lines, no further action is necessary when left-running Mach lines cross. This procedure is illustrated in Figure 47. Essentially, the left-running Mach line that crossed its upstream neighbor has been assumed to cease to exist. This would create errors in the (I,J) characteristic coordinate system if it were not accounted for. In the present computer program, when a left-running Mach line crosses its upstream neighbor, the solution is aborted and the x location of the solution point is specified by -1.0. The aborted point is given a location in the solution storage arrays and the (I,J) characteristic coordinate system. That left-running Nach line then acts like a phantom Mach line until it passes out of the flowfield at the nozzle wall, as illustrated in Figure 48. In any particular flowfield, several phantom left-running Mach lines may exist simultaneously, adjacent to each other or separated by real leftrunning Mach lines. In the computer program output, these phantom points are written out with their I.J characteristic coordinates and a blank line of data.

The occurence of crossed right-running Mach lines and phantom
left-running Mach lines considerably complicates the program logic.
Kowever, the procedure described above works very well and gives very accurate results as long as the Mach line crossings are not too numerous.

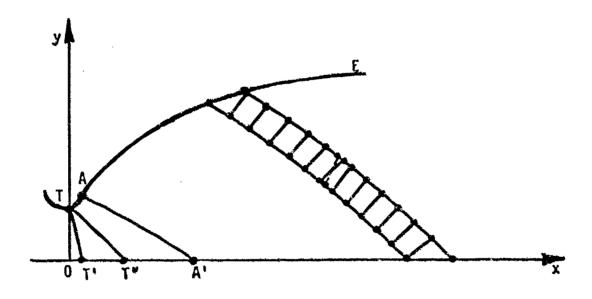


Figure 47. Crossing left-running Mach lines.

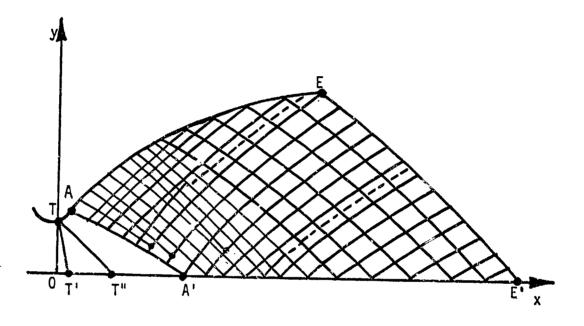


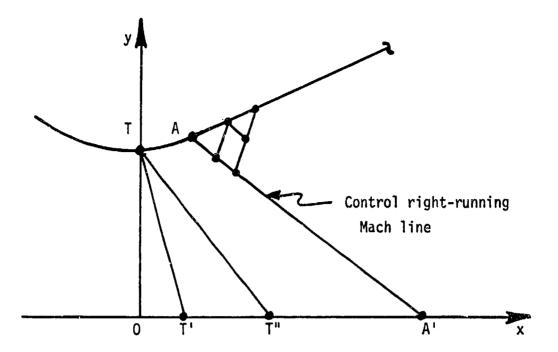
Figure 48. Propagation of phantom left-running Mach lines.

10. LEFT-RUNNING MACH LINE NETWORK IN THE BASIC NOZZLE

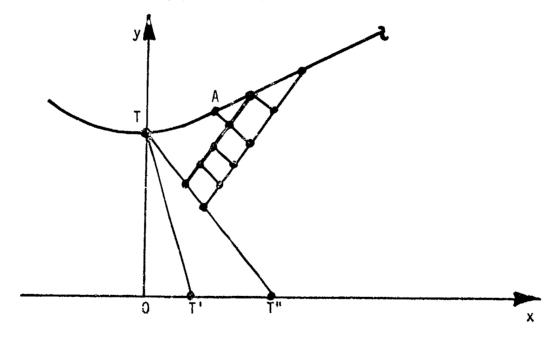
A left-running Mach line network can be constructed in the basic nozzle starting from either line AA', illustrated in Figure 40, or line TT", illustrated in Figure 43. Modes 2, 3, and 4, described in Section V.2, construct left-running Mach line networks starting from line AA'. Mode 5 constructs a left-running Mach line network starting from line TT". Those two network construction procedures are described in this section.

When the left-running Mach line network construction procedure illustrated in Figure 40 is employed, the flowfield is constructed along left-running Mach lines emanating from the last right-running Mach line at the downstream extent of the initial expansion flowfield. That right-running Mach line is called the control right-running Mach line. Starting from the flowfield illustrated in Figure 40, a left-running Mach line is initiated from the second point on the control right-running Mach line. That left-running Mach line is continued upward until it intersects the nozzle wall, as illustrated in Figure 49(a).

This procedure is repeated from successive points along the control right-running Mach line until one of two events occurs: (1) a left-running Mach line from the last point (i.e., the point on the axis) on the control right-running Mach line has been generated, or (2) the intersection of the left-running Mach line and the supersonic contour falls on the projection of the supersonic contour beyond the end of the basic nozzle, point E. In the first case, an axis point is determined as illustrated in Figure 35, and the next left-running Mach line emanates from that axis point. In the second case, the flow properties at point E.



(a) Starting from line AA'.



(b) Starting from line TT".

Figure 49. Propagation of a left-running Mach line to the basic nozzle wall.

are determined by the inverse wall point method described in Section V.4. This completes the flowfield determined by the basic nozzle contour.

When the left-running Mach line network construction procedure illustrated in Figure 43 is employed, the flowfield is constructed along left-running Mach lines emanating from the last right-running Mach line at the downstream extent of the flowfield from the initial-value line, line TT'. That right-running Mach line is called the control right-running Mach line. Enough left-running Mach lines have already been initiated from the control right-running Mach line to determine the flow properties along the initial expansion contour, contour TA. That procedure is discussed in Section V.8 and illustrated in Figure 42.

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The procedure described in Section V.8 is continued from successive points along the control right-running Mach line. Each successive left-running Mach line is continued upward until it intersects the nozzle wall, as illustrated in Figure 49(b). The procedure from here on is identical to that described above where line AA' is the control right-running Mach line.

The flowfield determined by either of the two aforementioned Mach line network construction procedures is illustrated schematically in Figure 50. This Mach line network construction procedure is employed when analyzing a basic nozzle (MODE = 2), a basic nozzle with an embedded right-running oblique shock wave (MODE = 3), or a scarfed nozzle with an attached right-running oblique shock wave (MODE = 4 or 5). In the computer program, the logic described in this section is implemented in subroutine MOCLRCI in overlay LINK20.

The procedure described above is very straightforward in a

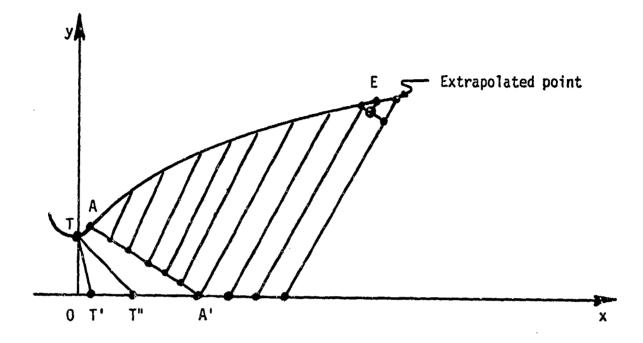


Figure 50. Extent of the left-running Mach line network along the basic nozzle contour.

continuous flowfield. However, when compression waves are present they tend to coalesce, and if enough coalescence occurs, oblique shock waves are formed. The coalescence of Mach lines and the appearance of oblique shock waves complicates both the flowfield itself and the logic procedure described above for calculating the flowfield. In the present investigation, the assumption is made that no strong oblique shock waves occur in the basic nozzle.

When left-running Mach lines coalesce (i.e., cross), the original Mach line is retained and the new Mach line is terminated downstream of the point of crossing. This procedure is illustrated in Figure 51.

Since the computational logic is based on following left-running Mach lines across the flowfield from the downstream extent of the initial value flowfield or the nozzle axis to the nozzle wall, each left-running Mach line must reach the nozzle wall. The above manner of handling crossing Mach lines achieves that goal. The new Mach line thus consists of the new points calculated before the point of crossing and the old points after the point of crossing. Thus, the new Mach line completely spans the nozzle flowfield from the downstream extent of the initial value flowfield or the nozzle axis to the nozzle wall. In the computer program output, the points after the point of crossing are written out as the solution for both left-running Mach lines.

When right-running Mach lines cross, the new Mach line is terminated. Since the computational logic is based on following left-running Mach lines across the flowfield instead of following right-running Mach lines, no further action is necessary when right-running Mach lines cross. This procedure is illustrated in Figure 52. Essentially, the right-running

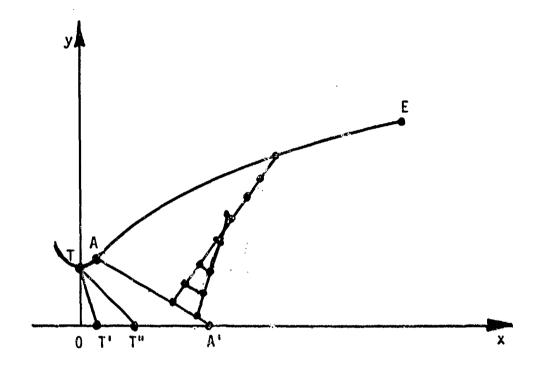


Figure 51. Crossing left-running Mach lines.

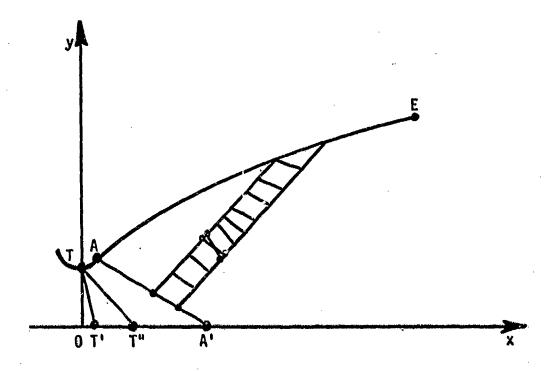


Figure 52. Crossing right-running Mach lines.

Mach line that crossed its upstream neighbor has been assumed to cease to exist. This would create errors in the (I,J) characteristic coordinate system if it were not accounted for. In the present computer program, when a right-running Mach line crosses its upstream neighbor, the solution is aborted and the x location of the solution point is specified by -1.0. The aborted point is given a location in the solution storage arrays and the (I,J) characteristic coordinate system. That right-running Mach line then acts like a phantom Mach line until it passes out of the flowfield at the nozzle axis, as illustrated in Figure 53. In any particular flowfield, several phantom left-running Mach lines may exist simultaneously, adjacent to each other or separated by real left-running Mach lines. In the computer program output, these phantom points are written out with their I,J characteristic coordinates and a blank line of data.

The occurence of crossed left-running Mach lines and phantom right-running Mach lines considerably complicates the program logic. However, the procedure described above works very well and gives very accurate results as long as the Mach line crossings are not too numerous.

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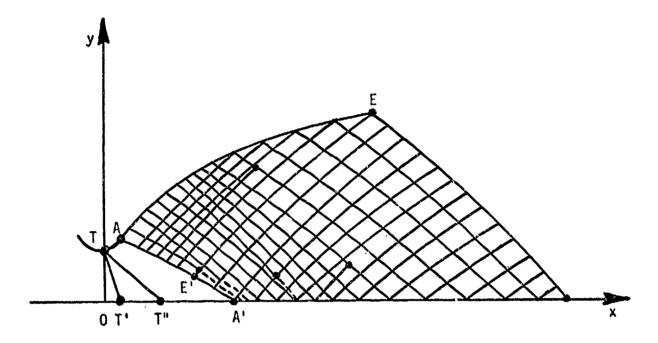


Figure 53. Propagation of phantom right-running Mach lines.

11. OBLIQUE SHOCK WAVE

An oblique shock emanates from the junction point of the basic nozzle and the nozzle extension, point E in Figure 25. This oblique shock wave propagates upstream into the flowfield emanating from the supersonic contour of the basic nozzle. The flowfield model for this oblique shock wave is presented in Section III.3. The implementation of that flowfield model into the overall numerical algorithm is discussed in the present section.

The first step in the determination of the oblique shock wave is the determination of the wave angle ε and the flow properties on the downstream side of the oblique shock wave at point E itself. This is accomplished by assuming a value for the wave angle ε , solving for the flow turning angle δ from equation (77), and comparing the resulting flow angle (i.e., $\theta_{\rm e}$ - δ) with the angle of the conical extension (i.e., $\theta_{\rm f}$). The wave angle ε is varied iteratively until ($\theta_{\rm e}$ - δ) converges to $\theta_{\rm f}$ within a prespecified convergence tolerance. The downstream Mach number M_2 is then determined from equations (74) to (76), and the remaining flow properties are determined from equations (78) to (82). This completes the solution for the oblique shock wave at point ε .

One small modification in the characteristic network is made before the oblique shock wave solution is initiated. The interpolated point at the intersection of the reward projection of the left-running Mach line through point E is substituted into the solution arrays in place of the point just below it, as illustrated in Figure 54.

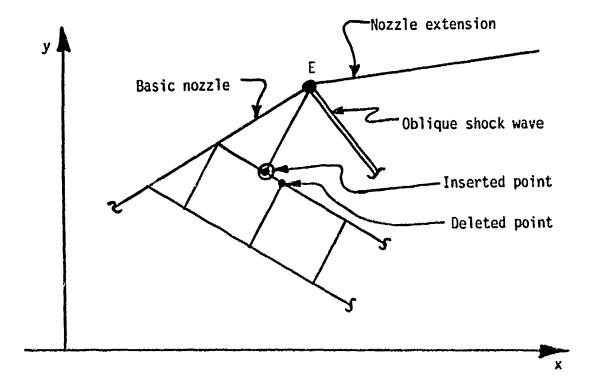


Figure 54. Insertion of interpolated point in the solution arrays.

The next step in determining the flowfield downstream of the oblique shock wave is to initiate a left-running Mach line from the next point on the right-running Mach line at the downstream extent of the upstream flowfield or from the nozzle axis if the upstream flowfield has reached the axis. That left-running Mach line is then propagated toward the scarfed nozzle extension until it reaches the oblique shock wave. It penetrates the oblique shock wave and continues until it intersects the scarfed nozzle extension, as illustrated in Figure 55.

The numerical procedure for obtaining the solution at the intersection of the left-running Mach line and the oblique shock wave is presented in Section III.4.C. The application of that model to the first point on the oblique shock wave downstream of point E is illustrated in Figure 56. The circled points are temporary points inserted on rearward projected right-running Mach lines. The shock wave point is located at the downstream intersection of the oblique shock wave from point E and the left-running Mach line from the upstream flowfield. The oblique shock wave is illustrated schematically as a double line to indicate that there are two sets of flow properties at each point on the shock wave; the upstream properties determined by application of the upstream interio point unit process and the downstream properties determined by the oblique shock wave equations. The first wall point downstream of point E is also determined during the iteriative procedure for the shock wave angle at the shock wave point.

After the first point on the oblique shock wave downstream of point E has been determined, the next point on the shock wave is determined by the procedure discussed in Section III.4.C. That procedure is illustrated

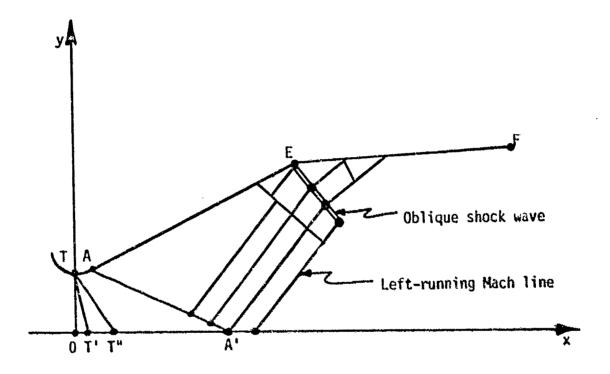
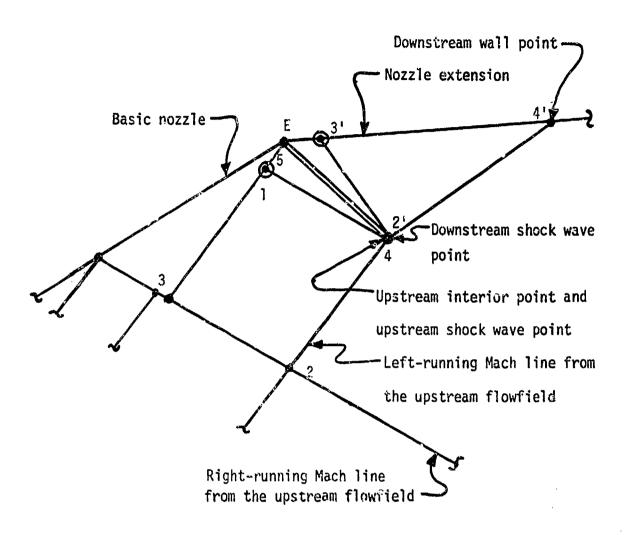


Figure 55. Propagation of a left-running Mach line from the upstream flowfield to the scarfed nozzle extension.



Known points:

2, 3, and 5

Solution points:

4, 2', and 4'

Interpolated points: 1 and 3'

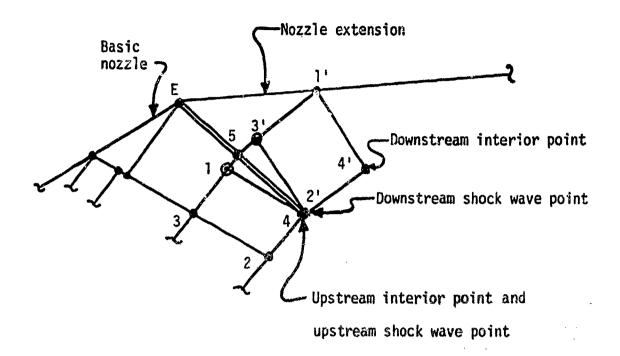
Figure 56. First point on the oblique shock wave downstream of point E.

in Figure 57. This procedure is repeated from subsequent points at the downstream extent of the upstream flowfield until the entire flowfield in the scarfed nozzle extension has been determined.

Since the strength of the oblique shock wave is greater than the strength of a Mach line, the oblique shock wave propagates at a steeper angle than the upstream right-running Mach line. Eventually the shock wave will overtake the upstream Mach line.

When the oblique shock wave overtakes the upstream right-running Mach line, that crossed right-running Mach line is simply terminated. The solution for the oblique shock wave is then obtained from the right-running Mach line preceding the crossed right-running Mach line. The crossed right-running Mach line is then treated as a phantom right-running Mach line in the characteristic coordinate system, as discussed in Section V.10 and illustrated in Figure 53. This procedure of terminating the crossed right-running Mach line and dropping back to the previous right-running Mach line is repeated until the shock wave point is successfully located. This procedure is illustrated in Figure 58.

The right-running Mach line downstream of the oblique shock wave is steeper than the shock wave because of the compression and decelleration of the flow caused by the shock wave. Eventually the downstream Mach line will overtake the shock wave. When the downstream right-running Mach line overtakes the oblique shock wave, that downstream right-running Mach is terminated. It is then treated as a phantom right-running Mach line in the characteristic coordinate system, as discussed in Section V.10 and illustrated in Figure 53. The next solution point is then determined at the intersection of the control left-running Mach line downstream of



Known points:

1', 2, 3, and 5

Solution points:

4, 2', and 4'

Interpolated points: 1 and 3'

Figure 57. Application of the oblique shock wave point unit process.

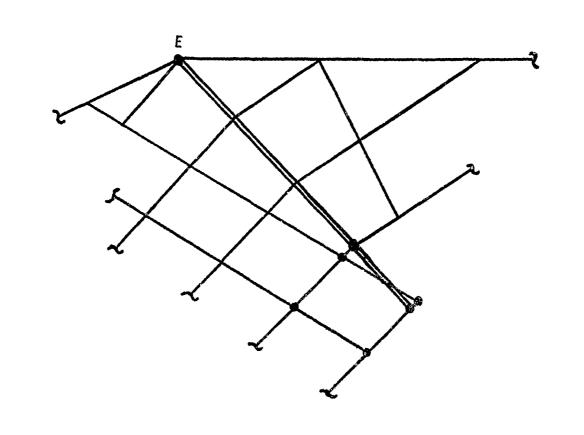


Figure 58. Otlique shock wave overtaking upstream right-running Mach line.

the terminated right-running Mach line. This procedure is illustrated in Figure 59.

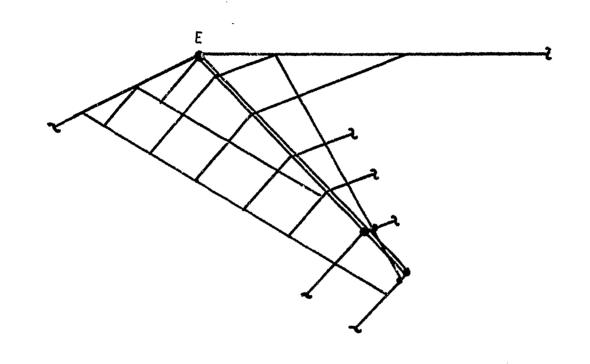


Figure 59. Downstream right-running Mach line overtaking the oblique shock wave.

12. FLOWFIELD IN THE SCARFED NOZZLE EXTENSION

The flowfield in the scarfed nozzle extension consists of the flow-field downstream of the oblique shock wave emanating from point E, the junction point of the basic nozzle and the nozzle extension. Due to the presence of an entropy gradient caused by the curved oblique shock wave, that flowfield is rotational. The flowfield model for a rotational flow is presented in Section III.5. The implementation of that procedure to determine the flowfield in the nozzle extension is described in this section.

The flowfield in the scarfed nozzle extension is determined by initiating left-running Mach lines from the upstream flowfield, as illustrated in Figure 55, and propagating those Mach lines to the scarfed nozzle extension. This procedure is repeated for succeeding left-running Mach lines until one of the three following situations arises:

- The oblique shock wave overtakes the right-running Mach line at the downstream extent of the upstream flowfield before the end of the scarfed nozzle extension has been reached.
- 2. The oblique shock wave intersects the nozzle axis before the end of the scarfed nozzle extension has been reached.
- 3. The left-running Mach line intersects the extension of the scarfed nozzle extension beyond point F. In this case, the indirect wall point unit process is applied at point F, as illustrated in Figure 60, and the scarfed nozzle flowfield is complete. This is the normal mode of operation.

Situations (1) and (2) are illustrated in Figures 61 and 62

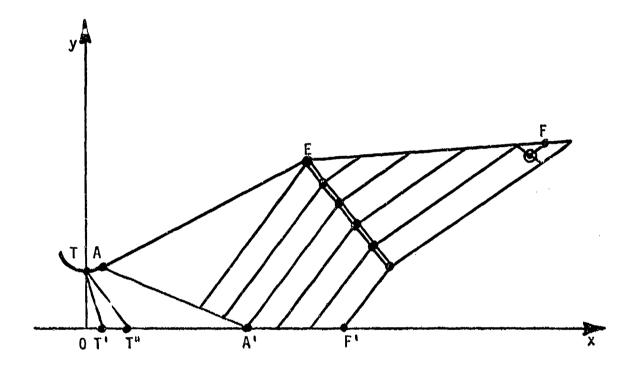


Figure 60. Normal solution for the flowfield in the nozzle extension.

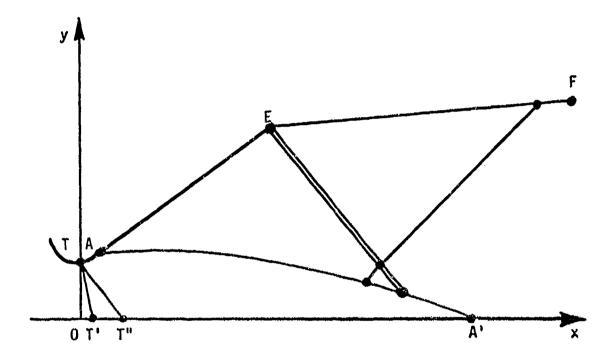
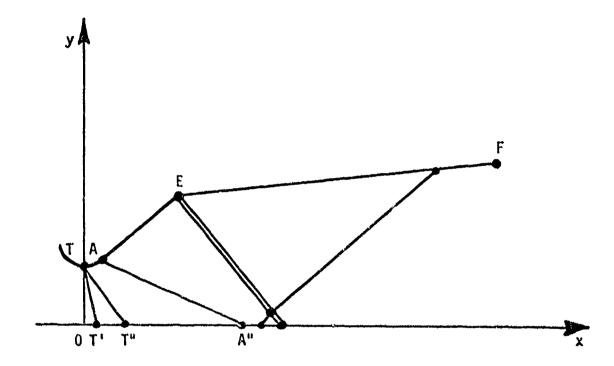


Figure 61. Oblique shock wave reaches the right-running Mach line at the upstream flowfield.

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Figure 62. Oblique shock wave reaches the nozzle axis.

respectively. When the oblique shock wave reaches the nozzle axis, the calculations are terminated. In an actual flowfield, the oblique shock wave generally transitions to a Mach disc and an embedded subsonic zone in the neighborhood of the nozzle axis. That phenomena cannot be predicted by the method of characteristics. In fact, the prediction of the Mach disc and the subsonic flow region is extremely difficult, and considerably beyond the scope of the present investigation. Consequently, in the present investigation, the oblique shock wave is simply terminated at the last interior point above the nozzle axis.

Practically speaking, this situation does not arise frequently.

When it does occur, the last point actually calculated is generally quite close to point F. In this situation, the approximation is made that a uniform flow exists from the last predicted wall point to point F.

Since point F is at the last remnant of the scarfed nozzle extension and the pressure level in this portion of the flowfield is very nearly atmospheric in many cases, the aforementioned assumption is generally quite good. When the last calculated point is far upstream of point F and/or the internal pressure level is greatly different from atmospheric pressure, the aforementioned assumption may introduce a small but significant error into the nozzle side force, and an even smaller and probably insignificant error into the nozzle axial force. In that case, some judgement as to the accuracy of the solution must be made based on experience.

The flowfield described in Situation (1) above can sometimes be continued further by using the MODE = 5 option. That option employs a left-running Mach line network construction procedure for the flowfield along the initial expansion contour. Consequently, the oblique shock

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wave can propagate all the way to the nozzle axis before the calculations must be terminated.

The implementation of the logic described in this section for determining the flowfield in the scarfed nozzle extension is contained in subroutine MOCLRCR in overlay LINK30.

13. SCARFED NOZZLE PERFORMANCE

The scarfed nozzle performance model is presented in Section IV.

The application of that model is discussed in this section. The implementation of this model in the computer program is contained in overlay LINK40.

The entire flowfield in the scarfed nozzle can be calculated by the procedures presented in the preceding sections. The performance of the basic nozzle is specified in terms of the nozzle mass flow rate \hat{m} and the axial thrust F_N .

Two steps are required to evaluate the performance of the scarfed nozzle extension. The first step is the calculation of the angle ψ , given by equation (121). That calculation is performed in subroutine GEON in overlay LINK30. The second step is the calculation of the scarfed nozzle thrust components F_{x} , SCE and F_{y} , SCE, given by equations (139) and (140). That calculation is performed in subroutine FORNON in overlay LINK30.

The final step in the overall numerical algorithm is the calculation of the performance of the entire nozzle. That performance is specified by F_X , F_Y , $(I_{SP})_{X^2}$, $(I_{SP})_Y$, and the effective scarfing angle β_{eff} . Those parameters are defined in Section IV.4. Figure 63 illustrates the relationship between the thrust components of the scarfed nozzle and the thrust components transmitted to the missile. The calculation of the overall performance parameters is accomplished in subroutine FORMON in overlay LINK3O.

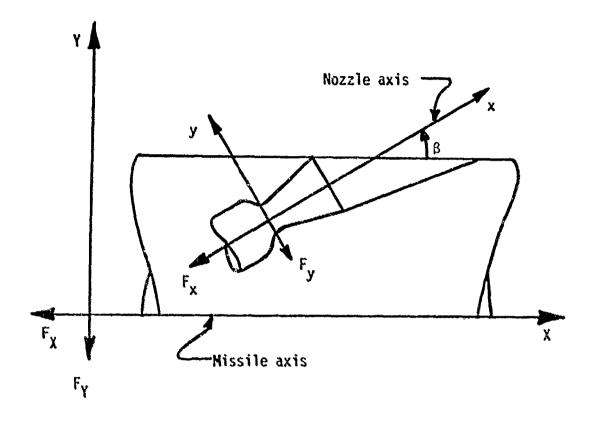


Figure 63. Overall nozzle performance model.

SECTION VI

COMPUTER PROGRAM

1. INTRODUCTION

A computer program has been developed to calculate the performance of a scarfed propulsive nozzle. A brief description of that program is presented in this section. The overlay structure of the program is described, and a brief description of the function of each program and subroutine is presented.

2. PROGRAM OVERLAY STRUCTURE

The overall program consists of 52 routines (5 program routines and 47 subroutines). The program is organized into an overlay scheme to minimize the amount of computer memory required. The overlay structure is presented in Figure 64. Two overlay levels are used: the resident overlay level and the primary overlay level.

OVERLAY (0,0) is the resident overlay which controls the overall execution of the program. Several commonly used subroutines are included in this overlay. OVERLAY (1,0) perofrms data input, parameter initialization, and generation of the initial-value line across the nozzle throat. OVERLAY (2,0) performs the evaluation of the flowfield in the basic nozzle. OVERLAY (3,0) constructs the oblique shock wave emanating from the junction between the basic nozzle and the nozzle extension and evaluates the flowfield in the nozzle extension. OVERLAY (4,0) calculates the thrust components in the scarfed extension, and determines the

LINKOO BOUNDYN MAINOO THERMOI BOUNDYE THERMOR INTERI AXISI OUTPUT **PDATA** THRUST PRINT FINAL PLOTWAL DWDOTR DWDOTL **PLOTRRC** LINK30 LINK40 LINK20 LINKTO HAIN40 MAIN30 MAINZO MAINTO MOCLRCR GEOM INPUT MUCIVL SHIFT? FORMON MOCARCR IVLINE HOVES KLIES MOCARCL POINTE IVLTAB MOCRRC SHOCK MOCLRCI IVLPERF . SHOCKU WESET COMPRES SHIFT SHOCKD ATTACH MOVET RRCHAR DRUALLR DRWALLI INWALLR INWALLI INTERR POINTEE AXISR

Figure 64. Frogram overlay structure.

performance of the overall scarfed propulsive nozzle.

In the following sections, a brief description is given of the function of each subroutine in the computer program. This information supplements the information available within the program in the form of comment statements.

3. OVERLAY (0,0)

MAINOO. This program routine is the main control routine in OVERLAY (0,0), the resident overlay. MAINOO first calls OVERLAY (1,0) for data input, parameter initialization, and generation of the initial-value line. MAINOO then calls OVERLAY (2,0) to calculate the flowfield in the basic nozzle, OVERLAY (3,0) to calculate the oblique shock wave and the flowfield in the nozzle extension, and OVERLAY (4,0) to calculate the nozzle performance.

THERMOI. This subroutine calculates the temperature T, pressure P, density ρ , speed of sound a, and Mach number M, corresponding to a specified value of the velocity magnitude V, for the irrotational flow of a thermally and calorically perfect gas.

THERMOR. This subroutine calculates the temperature T, speed of sound a, and Mach number M, corresponding to specified values of the pressure P, density ρ , and velocity magnitude V, for the rotational flow of a thermally and calorically perfect gas.

OUTPUT. This subroutine writes out the flow properties I, J, x, y, u, v, H, V, θ , P, ρ , T, M, m, and F at a point in the flowfield.

THRUST. This subroutine calculates the increment in thrust along the wall due to the pressure forces acting on the wall, the total nozzle thrust, and the nozzle specific impulse, as described in Section IV.2. The isentropic one-dimensional values are calculated for comparison, and the nozzle performance parameters F, F_{1-0} , η_F , I_{sp} , $I_{sp,1-0}$, and η_I are written out.

<u>DWDOTR</u>. This subroutine calculates the increments in mass flow rate and thrust between two adjacent points on a right-running Mach line

using the trapezoidal rule.

<u>DWDOTL</u>. This subroutine calculates the increment in mass flow rate and thrust between two adjacent points on a left-running Mach line using the trapezoidal rule.

<u>BOUNDYW</u>. This subroutine contains the description of the supersonic contour of the basic nozzle (see Section II.2). The point of intersection of a left-running Mach line and the supersonic contour is located in this subroutine.

<u>BOUNDYE</u>. This subroutine contains the description of the nozzle extension contour (see Section II.3). The point of intersection of a left-running Mach line and the nozzle extension is located in this subroutine.

<u>INTERI</u>. This subroutine implements the method of characteristics solution for an interior point in an irrotational flowfield, as illustrated in Figure 27.

AXISI. This subroutine implements the method of characteristics solution for an axis point in an irrotational flowfield, as illustrated in Figure 28.

<u>PDATA</u>. This subroutine saves the coordinates of the points along right-running Nach lines for subsequent plotting of the right-running Mach line network.

PRINT. This subroutine prints out the data saved by subroutine PDATA.

<u>PLOTWAL</u>. This subroutine initializes the plotting subroutines and plots the wall contour.

<u>PLOTRRC.</u> This subroutine plots the right-running Mach line network, using the data saved by subroutine PDATA.

4. OVERLAY (1,0)

MAIN10. This program is the main control routine in OVERLAY (1,0). Subroutine INPUT is called to read in and write out the input data. Subroutine IVLINE is called to generate the supersonic initial-value line.

INPUT. This subroutine reads in the input data and problem specifications through namelists DATA and WALL. All of the input variables are defined by comment cards. Default values are assigned to most of the input variables before the namelists are read in. Several parameters are initialized, and the units conversion factors are specified. The program description, the job title, and the problem specifications are written out. If requested, subroutine COMPRES is called to compress the tabular wall contour.

IVLINE. This subroutine selects the type of supersonic initial-value line to be employed.

KLIEG. This subroutine internally generates a supersonic initial-value line by the perturbation analysis developed by Kliegel and Levine (2). The location (i.e., x and y coordinates) of the points on the initial-value line are specified, and the velocity components (i.e., u and v) are calculated.

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IVLTAB. This subroutine reads in the coordinates (i.e., x and y), the Mach number M, and the flow angle θ along a tabular initial-value line. The velocity components (i.e., u and v) are internally calculated at each point.

 x and y of the points on the initial-value line. The calculated parameters include the Mach number M, the velocity magnitude V, the flow angle 0, the static pressure P, density ρ , and temperature T, the stagnation pressure P_t and temperature T_t , the mass flow rate \dot{m} flowing between each point and the nozzle axis, and the thrust associated with that mass flow rate. The nozzle overall performance parameters are also written out. These include the mass flow rate \dot{m} , the discharge coefficient C_D , the thrust F, the thrust efficiency η_T , the specific impulse I_{sp} , and the specific impulse efficiency η_T .

<u>COMPRES</u>. This subroutine compresses the axial coordinates of the tabular wall data.

ATTACH. This subroutine attaches the compressed tabular contour smoothly to the circular arc initial-expansion contour.

5. OVERLAY (2,0)

MAIN20. This program is the main control routine in OVERLAY (2,0). Subroutines MOCIVL, MOCARCR, MOCARCL, MOCRRC, and MOCLRCI are called to calculate the flowfield emanating from the initial-value line, the nozzle throat downstream circular arc contour, and the supersonic expansion contour.

MOCIVL. This subroutine is the control subroutine for calculating the flowfield emanating from the supersonic initial-value line. Subroutines INTERI and AXISI are called, as described in Section V.4, to calculate the flowfield. Subroutine MOVEl is called to transfer data between computational points and storage arrays, as required. Subroutine DWDOTR is called to integrate the mass flow rate and thrust along the right-running Mach lines as they are constructed. Subroutine OUTPUT is called to write out the flow properties at each solution point.

MOCARCR. This subroutine is the control subroutine for calculating the flowfield along right-running Mach lines emanating from the throat downstream circular arc contour of the basic nozzle as described in Section V.6. Subroutines INWALLI, INTERI, and AXISI are called to calculate the flowfield. Subroutine MOVEl is called to transfer data between computational points and solution arrays, as required. Subroutine THRUST is called to integrate the thrust generated along the nozzle wall, as described in Section IV.2. Subroutine DWDOTR is called to integrate the mass flow rate and thrust along the right-running Mach lines as they are constructed. Subroutine GUTPUT is called to write out the flow properties at each solution point.

MOCARCL. This subroutine is the control subroutine for calculating the flowfield along left-running Mach lines along the initial expansion contour of the basic nozzle, described in Section V.8. Subroutines INWALLI, INTERI, and AXISI are called to calculate the flowfield. Subroutine MOVEl is called to transfer data between computational points and solution arrays, as required. Subroutine THRUST is called to integrate the thrust generated along the nozzle wall, as described in Section IV.2. Subroutine DWDOTL is called to integrate the mass flow rate and thrust along the left-running Mach lines as they are constructed. Subroutine OUTPUT is called to write out the flow properties at each solution point.

MOCRRC. This subroutine is the control subroutine for calculating the flowfield along right-running Mach lines emanating from the supersonic contour of the basic nozzle as described in Section V.9.

Subroutines DRWALLI, INTERI, and AXISI are called to calculate the flowfield. Subroutine MOVEL is called to transfer data between computational points and solution arrays, as required. Subroutine THRUST is called to integrate the thrust generated along the nozzle wall, as described in Section IV.2. Subroutine DWBOTR is called to integrate the mass flow rate and thrust along the right-running Mach lines as they are constructed. Subroutine OUTPUT is called to write out the flow properties at each solution point.

MOCLRCI. This subroutine is the control subroutine for calculating the flowfield along left-running Nach lines that intersect the supersonic contour of the basic nozzle, as described in Section V.10. Subroutines DRWALLI, INTERI, and AXISI are called to calculate the flowfield.

Subroutine MOVEI is called to transfer data between computational points and solution arrays, as required. Subroutine THRUST is called to integrate the thrust generated along the nozzle wall, as described in Section IV.2. Subroutine DWDOTL is called to integrate the mass flow rate and thrust along the left-running Mach lines as they are constructed. Subroutine OUTPUT is called to write out the flow properties at each solution point.

RESET. This subroutine resets points along a right-running Mach line after it crosses the previous right-running Mach line, so that the remaining points to be calculated are defined to be the remaining points on the previous right-running Mach line, as illustrated in Figure 39.

SHIFT. This subroutine resets points along a left-running Mach line after it crosses the previous left-running Mach line, so that the remaining points to be calculated are the remaining points on the previous left-running Mach line, as illustrated in Figure 44.

MOVEL. This subroutine transfers data between grid points (points 1 to 4) and the storage arrays [X(150), Y(150), etc.] as requested by the logic control subroutines (MOCIVL, MOCARCR, MOCARCL, MOCRRC, and MOCLRCI).

<u>DRWALLI</u>. This subroutine implements the method of characteristics solution for a direct wall point in an irrotational flowfield, as illustrated in Figure 37.

<u>INVALLI</u>. This subroutine implements the method of characteristics solution for an inverse wall point in an irrotational flewfield, as illustated in Figures 31 and 32.

<u>POINTEE</u>. This subroutine defines the properties at the beginning of an embedded shock wave at point (IE, JE) in the characteristic coordinate system. This feature is used in MODE = 3 and only in MODE = 3.

6. OVERLAY (3,0)

MAIN30. This program is the main control program for OVERLAY (3,0), which calculates the rotational flowfield downstream of the oblique shock wave. Subroutine MOCLRCL is called to control the logic for calculating the flowfield. If an attached oblique shock wave exists at point E (MODE = 4 or 5), subroutine POINTE is called first to determine the properties at the shock wave at point E.

MOCLRCR. This subroutine contains the logic for constructing the oblique shock wave emanating from the junction of the basic nozzle and the nozzle extension (see Section V.11). Subroutine POINTE is called to determine the properties of the oblique shock wave at point E itself. Then subroutines INTERI, SHOCK, INTERR, and DRWALLR are called in sequence to determine the solution for an oblique shock wave point, as illustrated in Figure 48. Subroutine THRUST is called to calculate the thrust at the wall points. Subroutine MOVE2 is called to transfer data between computational points and storage arrays, as required. Subroutine DWDOTL is called to integrate the mass flow rate and thrust along the left-running Mach lines as they are constructed. Subroutine OUTPUT is called to write out the flow properties at each solution point. When the shock wave reaches the downstream extent of the upstream flowfield, as illustrated in Figure 53 or 54, the solution is terminated and the properties at the exit of the scarfed nozzle extension are extrapolated.

SHIFT2. This subroutine resets points along a left-running Mach line after it crosses the previous left-running Mach line, so that the remaining points to be calculated are the remaining points on the previous

left-running Mach line, as illustrated in Figure 44.

MOVE2. This subroutine transfers data between grid points (points 1 to 5) and the storage arrays [X(150), Y(150), etc.] as required by the logic control subroutine MOCLRCR and the other computational subroutines.

<u>POINTE</u>. This subroutine calculates the properties of the downstream side of the oblique shock wave at point E, as described in Section V.11.

SHOCK. This subroutine controls the logic for calculating the solution at a shock wave point, as discussed in Section V.11.

SHOCKU. This subroutine calculates the properties on the upstream side of the oblique shock wave as discussed in Section V.11 and illustrated in Figures 49 and 50.

SHOCKD. This subroutine calculates the properties on the downstream side of the oblique shock wave as discussed in Section V.11 and illustrated in Figures 49 and 50.

RRCHAR. This subroutine applies the compatibility relation along a rearward projected right-running Mach line from the downstream side of the oblique shock wave to check for overall convergence of the oblique shock wave calculation procedure. This procedure is described in Section V.11 and illustrated schematically in Figures 49 and 50.

<u>DRWALLR</u>. This subroutine implements the method of characteristics solution for a direct wall point in a rotational flowfield, as illustrated schematically in Figure 48.

INVALLE. This subroutine implements the method of characteristics solution for an inverse well point in a rotational flowfield, as discussed in Section V.12 and illustrated schematically in Figure

INTERR. This subroutine implements the method characteristics solution for an interior point in a rotational transferd, as discussed in Section V.12 and illustrated schematically in Figure 57.

AXISR. This subroutine implements the method of characteristics solution for an axis point in a rotational flowfield. In the present program, right-running Mach lines from the nozzle extension never reach the nozzle axis, as discussed in Section V.11 and illustrated schematically in Figures 53 and 54. Consequently, subroutine AXISR is not called in the present program.

7. OVERLAY (4,0)

<u>MAIN40</u>. This program is the main control program in OVERLAY (4,0). The purpose of OVERLAY (4,0) is to implement the calculation of the scarfed nozzle performance, as discussed in Section V.13. MAIN40 calls subroutine GEOM to calculate the angle ψ , given by equation (121). Subroutine FORMOM is called to calculate the performance of the scarfed nozzle and the missile performance.

<u>GEOM</u>. This subroutine calculates the angle ψ , given by equation (121).

FORMOM. This subroutine calculates the performance of the scarfed nozzle and the overall missile performance, as discussed in Section V.13.

SECTION VII

INPUT PARAMETERS

1. INTRODUCTION

The problem specifications are defined by input parameters which are read in by namelists DATA and WALL in subroutine SNPUT and by namelist IVSL in subroutine IVLTAB, all in OVERLAY (1,0). Only those parameters that are pertinent for a particular problem need be read in. Many parameters have default values, and need not be specified unless values different from the default values are to be considered. All of the input parameters are discussed in this section. Where appropriate, default values are specified. The default values specify Sample Case No. 1.

2. TITLE CARD

The first card of each data deck is a title card consisting of 80 alphanumeric characters of identifying information. This card may be blank, or contain any combination of allowable FORTRAN characters. This card must be the first card of every data deck, even if the card is blank. The format of the card is (8A10).

3. NAMELIST DATA

The parameters specified by namelist DATA are described in this section. All of these parameters have default values, which are reset to their original values before each data case. The first 14 parameters (IUNITS to NWRITE) are logic control parameters. The next two parameters (G and RG) specify the gas thermodynamic model. The following three parameters (PS, TS, and PA) define the nozzle operating conditions. The next ten parameters (DELTA to XMAX) specify the geometry of the basic nozzle. The following three parameters (AF, XF, and BETA) specify the geometry of the scarfed conical extension. The next two parameters (DYRATIO and DARATIO) are used to refine the spacing of the charactermistic network. The final five parameters (ICORI to IDUMP) are convergence control parameters and dump flags.

There are actually eight more parameters included in namelist DATA. These are ICMP, IPLOT, ICPLOT, YMAX, XDES, YDES, IE, and JE. These parameters are not used in the analysis of a conventional nozzle (MODE = 1 or 2) or in the analysis of a scarfed nozzle (MODE = 4 or 5). They are used in the analysis of a compressed nozzle with an embedded right-running oblique shock wave. Those eight parameters are not described in this report.

IUNITS A positive integer variable denoting the unit system employed in the analysis.

IUNITS	Unit system
1	English engineering (EE) units
2	System International (SI) units

The default value of IUNITS is 1.

MODE A positive integer variable denoting the type of analysis to be performed.

MODE	Type of analysis to be performed
0	Transcric initial-value line only
1	Conventional nozzle following right-running Mach
	lines
2	Conventional nozzle following left-running Mach
	lines
3	Conventional nozzle with an embedded RR shock wave
4	Scarfed nozzle following RRCS and LRCS
5	Scarfed nozzle following LRCS

The default value of MoDE is 4.

IWALL A positive integer variable denoting the type of basic nozzle contour to be analyzed.

IWALL	Type of basic nozzle contour
1	Conical contour
2	Quadratic contour
5	Tabular contour read in from TAPE5
8	Tabular contour read in from TAPE8

The default value of IWALL is 1.

JWALL A positive integer variable which controls punching out the basic nozzle contour.

JWALL.	Basic nozzle contour punch option
0	Don't punch
1	Do punch

The default value of JWALL is 0.

A positive integer variable denoting the number of points on the initial-value line across the nozzle throat. The default value of NI is 11.

A positive integer variable denoting the number of indirect wall points on the nozzle throat downstram circular arc contour.

The default value of NY is 15.

IVS A positive integer variable denoting the type of initial-value line to be employed.

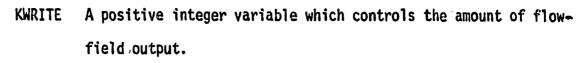
1 Kliegel's analysis (internally generated)
5 Tabular data read in from TAPE5
7 Tabular data read in from TAPE7
The default value of IVS is 1.

IOUT A positive integer variable which controls the width of the output.

TOUT	<u>Width of output</u>
Ü	80 column output
1	132 column output

IWRITE A positive integer variable specifying the right-running characteristic on which KWRITE is set to 2. The default value of IWRITE is 0.

JWRITE A positive integer variable specifying the left-running characteristic on which KWRITE is set equal to 2. The default value of JWRITE is 0.



KWRIT	<u>E</u> Outpu	option		
1	Write	out only first a	and last poir	its on each Mach
	line			
2	Write	out all points o	on each Mach	line, including
	delet	d points.		
3	Write	out all points o	on each Mach	line, excluding
	delet	d points		
4	Write	out only the fir	rst, second,	and last pages
The d	efault value	of KWRITE is 1.		

- LWRITE A positive integer variable specifying the right-running characteristic at which all of the 10 dump flags IDUMP(I) are set equal to 1. The default value of LWRITE is 0.
- MWRITE A positive integer variable specifying the left-running characteristic at which all 10 dump flags IDUMP(I) are set equal to 1. The default value of MWRITE is 0.
- The gas specific heat ratio γ (dimensionless). The default value of C is 1.2.
- RG The gas constant R (ft-1bf/1bm-R or k_0/k_0 -K). The default value of RG is 65.0.
- PS The nozzle inlet stagnation pressure P_t (lbf/in.² or N/m²). The default value of PS is 1000.0.
- TS The nozzle inlet stagnation temperature T_t (R or K). The default value of TS is 5000.0.

PA The ambient (i.e., atmospheric) pressure P_a (lbf/in.² or N/m²). The default value of PA is 14.696.

The nozzle geometric model is illustrated in Figure 65. The nozzle throat consists of a double circular arc contour. The supersonic contour of the basic nozzle attaches smoothly to the throat downstream circular arc contour (i.e., at point A, y_a and θ_a for the circular arc and the supersonic contour are the same). The length of the basic nozzle is x_e .

The basic nozzle contour may be conical (IWALL = 1). The conical contour may be specified by any of the following five sets of parameters.

- 1. The throat attachment angle θ_a and the nozzle length x_e .
- 2. The throat attachment angle θ_a and the exit lip radius y_e .
- 3. The throat attachment angle θ_{a} and the nozzle area ratio ϵ_{\star}
- 4. The nozzle legth x_e and exit lip radius y_e .
- 5. The nozzle length $\mathbf{x}_{\mathbf{e}}$ and nozzle area ratio $\boldsymbol{\epsilon}.$

The basic nozzle contour may be a quadratic (IWALL = 2). The quadratic contour may be specified by any of the following three sets of parameters.

- 1. The throat attachment angle θ_a , the exit lip angle θ_e , and the nozzle length x_e .
- 2. The throat attachment angle θ_a , the nozzle length x_e , and the exit lip radius y_e .
- 3. The throat attachment angle θ_a , the nozzle length x_e , and the nozzle area ratio ϵ .

The default values of all these parameters are 0.0. The desired option is selected simply by specifying nonzero values for the appropriate variables. The remaining variables are internally calculated

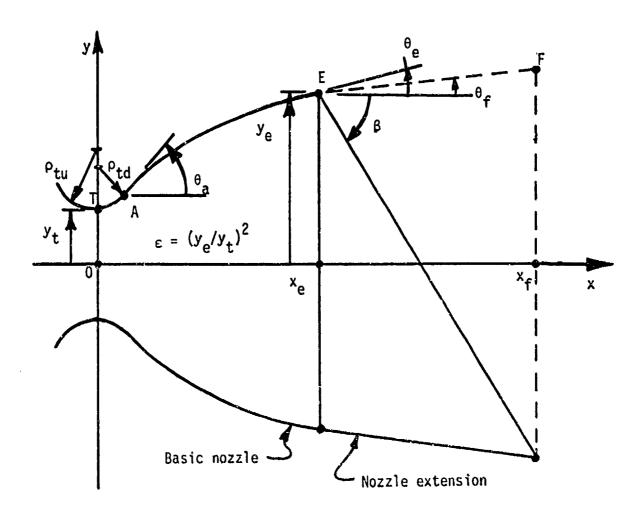


Figure 65. Nozzle geometric model.

from those specified.

The remaining option for specifying the supersonic contour of the basic nozzle, the tabular wall option, is discussed in Section VII.4.

The conical extension is completely specified by the cone angle $\theta_{\mathbf{f}}$ and the x coordinate of the end of the extension, $\mathbf{x}_{\mathbf{f}}$. The conical extension is attached to the basic nozzle at point E.

DELTA A real variable denoting the type of coordinate system to be considered.

DELTA Coordinate system

0.0 Two-dimensional planar

1.0 Two-dimensional axisymmetric

The default value of DELTA is 1.0.

YT The nozzle throat radius y_t (in. or m). The default value of YT is 1.0.

RTU The nozzle throat upstream circular arc radius of curvature ρ_{tu} (in. or m). The default value of RTU is 1.0.

RTD The nozzle throat downstream circular arc radius of curvature ρ_{td} (in. or m). The default value of RTD is 0.01.

AA The nozzle throat attachment angle θ_a (deg). The default value of AA is 0.0.

AE The basic nozzle exit lip angle at point E, θ_e (deg). The default value of AE is 0.0.

EPS The basic nozzle area ratio $\varepsilon = (y_e/y_t)^2$ (dimensionless). The default value of EPS is 0.0.

YE The basic nozzle exit lip radius y_e (in. or m). The default

value of YE is 0.0.

XE The x coordinate of the length of the basic nozzle x_e (in. or m). The default value of XE is 0.0.

Mach lines may be terminated to save computational effort when the flowfield beyond the end of the nozzle is not required, as illustrated in Figure 66. XMAX is only used when analyzing a basic nozzle following right-running Mach lines (MODE = 1). In that case, XMAX must be at least as large as XE, the length of the basic nozzle. For the analysis of scarfed nozzles (MODE = 4 or 5), XMAX must be a large number so that the entire flowfield from the basic nozzle is obtained. The default value of XMAX is 1000.0.

The conical nozzle extension is also illustrated in Figure 65. The conical extension may be specified by either of the following options.

- 1. The cone angle $\theta_{\mbox{\bf f}}$ and the x coordinate of the end of the conical extension $x_{\mbox{\bf f}}$.
- 2. The cone angle $\theta_{\mbox{\scriptsize f}}$ and the nozzle scarfing angle $\beta_{\mbox{\scriptsize t}}$

The default values of x_f and β are both 0.0. The desired option is selected by simply specifying a nonzero value for either x_f or β . The other variable is internally computed from the one specified.

- AF The conical extension wall angle θ_f (deg). The default value of AF is 0.0.
- XF The x coordinate of the end of the conical extension x_f (in ϕ) m). The default value of XF is 0.0.

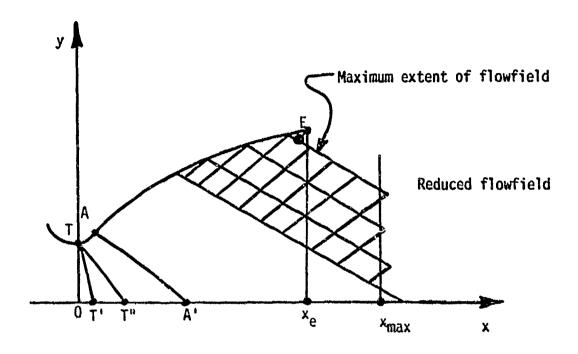


Figure 66. Definition of the variable x_{max} .

BETA The nozzle scarfing angle. The default value of BETA is 0.0.

El

The next two variables, DYRATIO and DARATIO, are used to refine the spacing of the characteristic network. DYRATIO is discussed in Section V.3, and DARATIO is discussed in Section V.5.

DYRATIO A real variable specifying the ratio of the final Δy on the initial-value line adjacent to the nozzle wall to the initial Δy adjacent to the nozzle axis (dimensionless). The default value of DYRATIO is 1.0.

DARATIO A real variable specifying the ratio of the final $\Delta\theta$ on the nozzle throat downstream circular arc adjacent to the throat attachment point, point A in Figure 37, to the initial $\Delta\theta$ adjacent to the nozzle throat, point T in Figure 37 (dimensionless). The default value of DARATIO is 1.0.

The remaining five parameters are convergence control parameters and dump flags.

ICOR A positive integer variable specifying the number of applications of the corrector in the modified-Euler predictor-corrector numerical solution of the characteristic equations and compatibility relations. ICOR = 0 yields the Euler predictor method,

ICOR = 1 yields the modified-Euler predictor-corrector method, and ICOR ≥ 2 yields the modified-Euler predictor-corrector method with (ICOR - 1) iterations. The default value of ICOR is 2.

A positive real variable denoting the tolerance (in. or m) for location convergence in the numerical method of characteristics calculations. The default value of El is 0.0, causing the unit process calculations to terminate after ICOR applications of the

integration algorithm.

E2 A positive real variable denoting the fractional convergence tolerance (dimensionless) for flow property convergence in the numerical method of characteristics. The default value of E2 is 0.0, causing the unit process calculations to terminate after ICOR applications of the integration algorithm.

IDUMP A one-dimensional integer variable array dimensioned at 10. Each element of IDUMP activates output dumps of selected parameters during the flowfield calculations for use in debugging problems.

A value of 0 suppresses the dumps, and a value of 1 activates the dumps. The default values of all 10 elements of IDUMP are 0.

IDUMP (I)	Subroutines with dumps activated
1	MOCLRCI, MAIN3O, MOCLRCR
2	RESET, SHIFT, SHIFT2
3	POINTEE, POINTE, SHOCK, SHOCKD, SHOCKD
4	COMPRES, ATTACH
5	RRCHAR
6	BOUNDYW, BOUNDYE, DRWALLI, DRWALLR
7	INTERI, INTERR
8	INWALLI, INWALLR
9	AXISI, AXISR
10	not used

This completes the specification of all of the input data read in by namelist DATA.

4. NAMELIST WALL

The parameters specified by namelist WALL are described in this section. These parameters specify the tabular wall contour option (IWALL = 5 or 8 in namelist DATA) for the basic nozzle. For the analytical wall options (IWALL = 1 or 2), this namelist should be omitted from the data deck. The tabular nozzle option is discussed in Section II.3.d.

The tabular nozzle geometry is illustrated in Figure 67. The wall table, when complete, will have NWALL pairs of x and y coordinates in the arrays XW and YW, respectively. The first point in the table will be the throat attachment point, point A. The coordinates of point A (i.e., x_a and y_a) are calculated internally and stored in XW(1) and YW(1). Consequently, the values of ho_{td} and $heta_a$ (i.e., RTD and AA) in namelist DATA must be specified. The remaining wall points are read in through namelist WALL. The first input point is XW(2) and YW(2). The succeeding points follow in order. The last point specified in the tabular data is defined internally to be point E (i.e., x_e and y_e). The number of points actually read in through namelist WALL is denoted by NWALL. The total number of points in the wall table is NWALL + 1, due to the insertion of point A as the first point in the table. The input value of NWALL is increased by one internal to the program to reflect the total number of points in . . . wall table. The computer program locally fits a quadratic polynomial to the tabular data when calculating the intersection point of a left-running Mach line and the tabular nozzle wall.

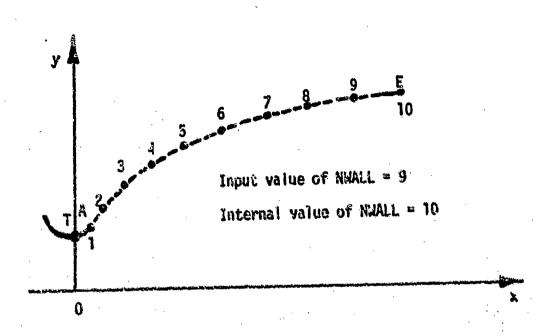


Figure 67. Tabular nozzle geometry.

NWALL An integer variable specifying the number of points in the tabular wall table (must be 99 or less). No default value is specified for NWALL.

A one-dimensional real variable array dimensioned at 99. XW(I)

(I = 1, NWALL) denotes the axial location (in. or m) of the 1th
tabular wall point. No default values are specified for XW(I).

YW A one-dimensional real variable array dimensioned at 99. YW(I)

(I = 1, NWALL) denotes the radial location (in. or m) of the 1th
tabular wall point. No default values are specified for YW(I).

5. NAMELIST IVSL

The parameters specified by namelist IVSL are described in this section. These parameters specify the tabular initial-value line option (IVS = 5 or 7 in namelist DATA). For the internally generated initial-value line option (IVS = 1), this namelist should be omitted from the data deck. The tabular initial-value line option is discussed in Section V.3.B.

The tabular initial-value line is illustrated in Figure 68. The first point on the tabular initial-value line must be the nozzle axis point, and the last point must be the nozzle throat point, point T. The number of points NI (specified in namelist DATA) and their spacing is arbitrary. The maximum number of tabular initial-value line points is 30.

- XIV A one-dimensional real variable array dimensioned at 30. XIV(I) (I = 1, NI) denotes the axial location (in. or m) of the $I\underline{th}$ tabular initial-value line point. No default values are specified for XIV(I).
- RIV A one-dimensional real variable array dimensioned at 30. RIV(I) (I=1, NI) denotes the radial location (in. or m) of the $I\underline{th}$ tabular initial-value line point. No default values are specified for RIV(I).
- MIV A one-dimensional real variable array dimensioned at 30. MIV(I) (I = 1, NI) denotes the Mach number (dimensionless) at the Ith point on the tabular initial-value line. No default values are specified for MIV(I).

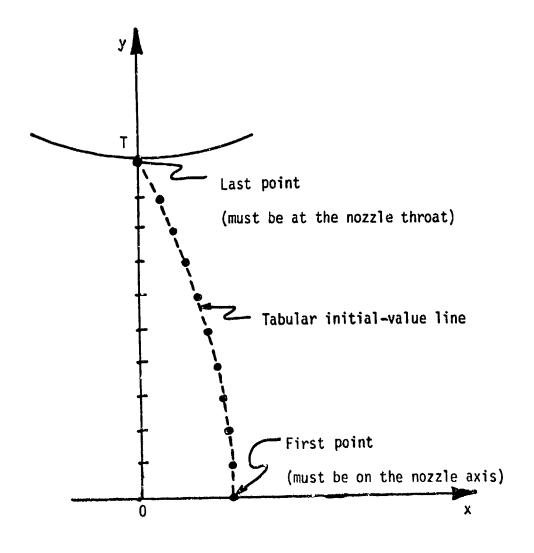


Figure 68. Tabular initial-value line.

A one-dimensional real variable array dimensioned at 30. TIV(I) (I = 1, NI) denotes the flow angle (deg) at the Ith point on the tabular initial-value line. No default values are specified for TIV(I).

SECTION VIII SAMPLE CASES

1. INTRODUCTION

Fifteen sample cases are presented in this section to illustrate the application of the computer program for calculating the flowfield in scarfed propulsive nozzles. For each sample case, a discussion of the problem is presented, the input data are discussed, and the data deck required to execute the program is presented. Selected portions of the computer output are presented for some of the sample cases.

The first eight sample cases are all concerned with the same scarfed nozzle configuration. That nozzle consists of a conical nozzle and a cylindrical extension. The nominal case is specified by a nozzle cone angle θ_a = 15 deg, a nozzle area ratio ϵ = 10, a cylindrical extension with θ_f = 0 deg, and a scarf angle β = 30 deg.

Sample Case No. 1 includes a complete discussion of the problem and the input data for the case in which the conical contour is specified by the cone angle θ_a and the area ratio ϵ . Sample Cases No. 2 to 4 are identical to Sample Case No. 1, except that the conical contour is specified by the exit lip radius y_{ϵ} , the note another x_{ϵ} , and tabular data, respectively. Sample Case No. 5 illustrates the tabular initial-value line option. Sample Case No. 6 illustrates an optional Nach line network construction feature. Sample Case No. 7 illustrates the variable initial-value line spacing option and the variable inverse wall

point spacing option on the nozzle throat downstream circular arc.

Sample Case No. 8 is the same as Sample Case No. 1, except that the length x_f of the conical extension is specified instead of the scarf angle β . Sample Case No. 9 has the same nozzle as Sample Case No. 1, but the scarfed nozzle extension is a cone with $\theta_f = 5$ deg.

Sample Cases No. 10 to 12 illustrate various options for analyzing a simple conical nozzle without an extension. Sample Case No. 10 considers the nominal case modified by specifying the angle $\theta_{\rm f}$ of the conical extension to be the same as the nozzle cone angle $\theta_{\rm a}$. Sample Case No. 11 analyzes the combined nozzle and extension considered in Sample Case No. 10 as a single conical nozzle without an extension. Sample Case No. 12 illustrates an optional Mach line network construction feature for a single nozzle. This same feature is illustrated in Sample Case No. 6 for the nominal case with a cylindrical extension.

Sample Cases No. 13 to 15 are all concerned with the same scarfed nozzle configuration. That nozzle consists of a quadratic nozzle contour and a cylindrical nozzle extension. Sample Case No. 13 is the nominal case for the quadratic nozzle contour. The nominal case is specified by a throat attachment angle θ_a = 25 deg, an area ratio ϵ = 10, and a nozzle length x_e = 8.07104661. The scarfed nozzle extension is cylindrical and the scarf angle β = 30 deg. This nozzle has the identical envelope as the conical nozzle nominal case, Sample Case No. 1. Sample Cases No. 14 and 15 are identical to Sample Case No. 13, except that the quadratic nozzle contour is specified by the nozzle exit lip radius y_e and the nozzle exit lip angle θ_e , respectively.

This sample case is the basic problem considered in all of the sample cases. The nozzle to be analyzed is a conical nozzle having a scarfed cylindrical extension as illustrated in Figure 69. The English Engineering (EE) system of units is employed in all of the sample cases (IUNITS = 1). The problem to be considered is the analysis of a scarfed propulsive nozzle (MODE = 4). The basic nozzle is a conical nozzle (IWALL = 1). Punched output of the wall contour is not desired (JWALL = 0). Eleven equally spaced points are desired on the initial-value line (NI = 11 and DYRATIO = 1.0), and 15 equally spaced points are desired on the nozzle throat downstream circular arc (NT = 15 and DARATIO = 1.0). The initial-value line is to be generated internally by Kliegel's method (IVS = 1). Every point on each Mach line is to be written out (KWRITE = 2), and KWRITE is to remain at the value of 2 for the entire calculation (IWRITE = JWRITE = 0).

The gas specific heat ratio γ is 1.2 (GANNA = 1.2) and the gas constant R is 65.0 (ft-lbf)/(lbm-R).

The nozzic inlet stagnation pressure P_t and temperature T_t are 1000.9 lbf/in.² and 5000.0 R. respectively (PS = 1000.0 and TS = 5000.0). The ambient pressure P_a is 14.696 lbf/in.² (PA = 14.696).

The nozzle is axisymmetric (DELTA = 1.0). The throat radius y_t is 1.0 in. (YT = 1.0), the nozzle throat upstream radius of curvature ρ_{tu} is 1.0 in. (RTU = 1.0), and the nozzle throat downstream radius of curvature ρ_{td} is 0.01 in. (RTD = 0.01). The basic nozzle contour is conical (IWALL = 1), and the conical geometry is specified by the nozzle throat

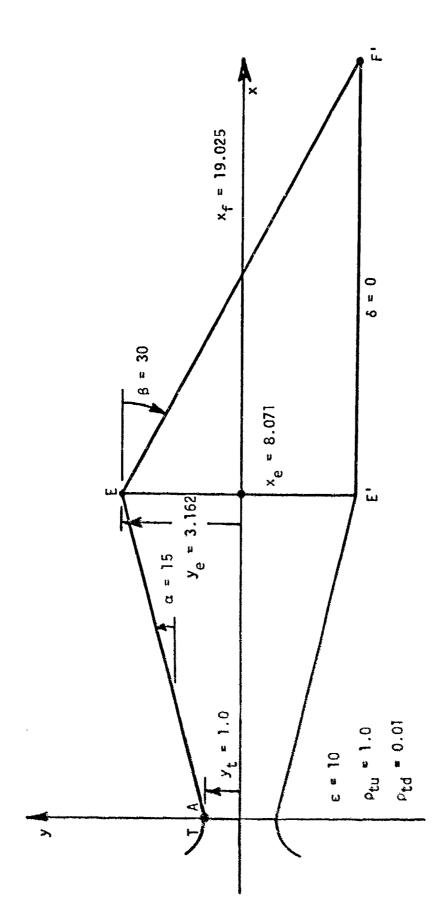


Figure 69. Scarfed nozzle geometry for Sample Case No. 1.

attachment angle θ_a of 15 degrees (AA = 15.0) and an area ratio ε = 10 (EPS = 10.0). The parameters θ_a and ε completely specify the geometry of the conical nozzle. Hence, the exit lip radius y_e and the nozzle length x_e are left at their default values of 0.0 (YE = 0.0 and XE = 0.0). The complete flowfield from the basic nozzle must be calculated since a scarfed extension is present. Thus, XMAX is left at its default value of 1000.0.

The cone angle θ_f of the scarfed conical extension is 0.0 deg (AF = 0.0). The scarf angle β = 30.0 deg (BETA = 30.0). The parameters θ_f and β completely specify the geometry of the scarfed conical extension. Hence, the length x_f of the scarfed conical extension is left at its default value of 0.0 (XF = 0.0).

Two applications of the corrector (ICOR = 2) are to be made in the modified-Euler predictor-corrector numerical integration procedure used in the method of characteristics unit processes. All fluwfield points are to be corrected two times, so the convergence tolerances are set equal to 0.0 (El = 0.0 and E2 = 0.0).

This completes the specification of namelist DATA. Since the initial-value line and the basic nozzle contour are internally generated (IVS = 1 and IWALL = 1), namelists IVSL and WALL are omitted from the data deck. The data deck for Sample Case No. 1 is presented below.

SAMPLE CASE NO. 1. NOMINAL CASE WITH AA=15, EPS=10, AF=0, AND BETA=30 SDATA HODE=4, IWALL=1, NI=11, NT=15, KWRITE=2,

G=1.2, %G=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, EPS=10.0, AF=0.0, BETA=30.0 \$

The computer output for Sample Case No. 1 is presented in Append. A. The first page presents the program abstract, the job title, the problem specifications, and the thermodynamic model. The second page describes the nozzle geometry. The next two pages present a complete listing of namelist DATA. The fifth page presents the initial-value line properties and performance parameters.

The three following pages present the flow properties at each point along the right-running Mach lines emanating from the initial-value line, as illustrated in Figure 36. The first point in each set of data is the initial-value line point. That point is followed by the succeeding points along the right-running Mach line. The last point in each set of data is the point where the right-running Mach line intersects the nozzle axis. Since 11 points were specified on the initial-value line (NI = 11), 11 right-running characteristics are present in this initial portion of the flowfield (I = 1 to 11).

The next nine pages present the flow properties at each point along the right-running Mach lines emanating from the prespecified points on the initial expansion contour, that is, the circular arc throat downstream radius of curvature, as illustrated in Figure 40. Since 15 points were specified along the initial-expansion contour (NT = 15), 15 right-running characteristics are present in this portion of the flowfield (I = 12 to 26). The nozzle performance parameters at each prespecified wall point are listed ahead of the corresponding Mach line data.

The next 10 pages present the flow properties at ea n point along the left-running characteristics emanating from points along the final right-running characteristic from the initial-expansion contour, as

illustrated in Figure 50. The nozzle performance parameters at each direct wall point at the end of each left-running Mach line are listed at the end of each set of Mach line points. These left-running Mach lines are generated until the end of the basic nozzle contour is reached. In this sample case, 30 left-running Mach lines are required (J = 1 to 30). The last point (I = 55, J = 30) is the exit lip point of the basic nozzle, point E.

The next 12 pages present the flow properties at each point along the left-running characteristics that reach the scarfed nozzle extension. All of these Mach lines pass through the right-running oblique shock wave emanating from the junction between the basic nozzle contour and the scarfed nozzle extension, point E. The first point (I = 55, J = 30) is the shock wave point at point E itself. Each shock wave point has two lines of data: the first line contains the flow properties on the upstream side of the shock wave and the second line contains the flow properties on the downstream side of the shock wave. The nozzle performance parameters at each point on the scarfed nozzle extension at the end of each left-running Mach line are listed at the end of each set of Mach line points. These left-running Mach lines are generated until the end of the scarfed nozzle extension is reached. In this sample case, 20 left-running Mach lines are required (J = 31 to 50). The last point (I = 71, J = 50) is the exit lip point of the scarfed nozzle extension, point F.

The next two pages present a summary of the nozzle performance parameters at each point along the nozzle wall contour calculated during the method of characteristics analysis.

The next page summarizes the performance of the scarfed nozzle

extension.

The final page of output presents a summary of the overall scarfed nozzle performance parameters and a summary of the overall missile performance parameters. The final result of the analysis is the axial specific impulse delivered to the missile, denoted by ISPXM. For the present sample case, this value is 206.785 (1bf-sec)/1bm.

Market Har Market Control of the

This computer run required approximately 25 sec of CPU time on a CDC 6500 computer and generated approximately 2100 lines of printed output.

In many cases, the amount of output described above is not needed. A considerable reduction in output is achieved by setting KWRITE = 1.

In that case, every section of output described about is still obtained, but only the first, last, and shock wave points are written out along each characteristic. For Sample Case No. 1 with KWRITE = 1, approximately 21 sec of CPU time were required and approximately 500 lines of printed output were generated. Even less printed output is obtained by setting KWRITE = 4. In that case, only the first two pages describing the problem specifications and the last page summarizing the overall performance parameters are written out. For Sample Case No. 1 with KWRITE = 4, approximately 15 sec of CPU time were required and approximately 70 lines of printed output were generated. The portion of the output for Sample Case No. 1 that is abbreviated by setting KWRITE = 1 is presented in Append. B.

Sample Case No. 2 is identical to Sample Case No. 1, except that the conical portion of the basic nozzle is specified by the exit radius y_e instead of the cone angle θ_a , and the amount of output is reduced. The identical basic nozzle contour is obtained for $y_e = 3.16227766$ in. Thus, the only changes to the input data are to delete AA = 15.0, to add YE = 3.16227766, and to change KWRITE to 1. The data deck for Sample Case No. 2 is presented below. The output is identical to the abbreviated output for Sample Case No. 1.

\$DATA MODERAL IWALL=1, NI=11, NT=15, KWRITE=1,
G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,
YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, YE=3.16227766, AF=0.0, BETA=30.0 \$

Sample Case No. 3 is identical to Sample Case No. 1, except that the conical portion of the basic nozzle is specified by the nozzle length x_e instead of the cone angle θ_a . The identical basic nozzle is obtained for $x_e = 8.07104661$ in. Thus, the only changes to the input data are AA = 0.0 and XE = 8.07104661. The data deck for Sample Case No. 3 is presented below. The output is identical to the abbreviated output for Sample Case No. 1.

\$AMPLE CASE NO. 3. NOMINAL CASE WITH XE=8.07104661 SPECIFIED

\$DATA MODE=4, IWALL=1, NI=11, NT=15, KWRITE=1,

G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, XE=8.07104661, AF=0.0, BETA=30.0 \$

Sample Case No. 4 is identical to Sample Case No. 1, except that the supersonic contour of the basic nozzle is read in as tabular data from TAPES (IWALL=5). For the tabular wall option, the basic nozzle throat contour must still be specified, since the first point in the wall contour table is computed internally as point A. Thus, the throat attachment angle must be explicitly specified as 15.0 deg (AA=15.0). The basic nozzle exit lip point, point E, is internally defined to be the last point in the wall table. Thus, x_e and y_e or ε (i.e., XE, YE, and EPS) do not need to be input. Thus, the only change in the input parameters required in namel at DATA is IWALL=5.

For the tabular wall option, namelist WALL must be specified. Only two points are required to specify a conical wall (NWALL=2). The first point can be any point downstream of the throat attachment point, point A. For a contoured nozzle, the first point should be close to point A to achieve an accurate representation of the nozzle wall. For a conical wall, the location of the first point is irrelevant, as long as it lies between the points A and E. That point is specified as (4.03681741, 2.08130920), which is exactly midway between points A and E. The last point is point E (8.07104661, 3.16227766).

The data deck for Sample Case No. 4 is presented below.

\$AMPLE CASE NO. 4. NOMINAL CASE WITH TABULAR NOZZLE WALL \$DATA MODE=4, IWALL=5, NI=11, NT=15, KWRITE=1, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696, YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, AF=0.0, BETA=30.0, \$WALL NWALL=2, XW(2)=4.03681741, 8.07104661, YW(2)=2.08130920, 3.16227766 \$

The output for Sample Case No. 4 is the same as the abbreviated output for Sample Case No. 1, except that the tabular wall data are written out on the second page of output. That page is presented in Append. C.

Sample Case No. 5 is identical to Sample Case No. 1, except that the initial-value line is read in as tabular data from TAPE5 (IVS=5). The initial-value data are taken from the output of Sample Case No. 1. The only change in the input required in namelist DATA is IVS=5.

For the tabular initial-value line option, namelist IVSL must be specified. The point locations and flow properties are taken directly from the output of Sample Case No. 1.

The data deck for Sample Case No. 5 is presented below.

The output for Sample Case No. 5 is the same as the abbreviated output for Sample Case No. 1.

SAMPLE CASE NO. 5. NOMINAL CASE WITH TABULAR INITIAL-VALUE LINE \$DATA MODE=4, IWALL=1, NI=11, NT=15, IVS=5, KWRITE=1,

G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, AF=0.0, BETA=30.0 \$

\$IVSL XI(1)= 0.262, 0.260, 0.252, 0.239, 0.220, 0.197, 0.168,

0.134, 0.094, 0.050, 0.000,

YI(1)= 0.000, 0.100, 0.200, 0.300, 0.400, 0.500, 0.600, 0.700, 0.800, 0.900, 1.000,

MI(1)= 1.031, 1.032, 1.035, 1.041, 1.049, 1.060, 1.075,

1.095, 1.121, 1.156, 1.203,

TI(1)= 0.000, 0.110, 0.210, 0.300, 0.370, 0.410, 0.400, 0.350, 0.260, 0.130, 0.000 \$

Sample Case No. 6 illustrates the option of starting the left-running characteristic network construction from the last right-running characteristic from the initial-value line instead of from the last right-running characteristic from the initial expansion contour. This option is useful when the shock wave intersects the last right-running characteristic from the initial expansion contour. This option is specified by MODE=5. All other parameters are identical to the data for Sample Case No. 1. The data deck for Sample Case No. 6 is presented below. The output for Sample Case No. 6 is presented in Append. D.

\$DATA MODE=5, IWALL=1, NI=11, NT=15, KWRITE=2, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, EPS=10.0, AF=0.0, BETA=30.0 \$

The first six pages of output for Sample Case No. 6 are identical to the output for Sample Case No. 1. The next two pages present the solution along the first left-running characteristic (J=2) emanating from the downstream extent of the initial-value line. The next ten pages present the solution along the 28 left-running characteristics (J= 3 to 30) that reach the wall of the basic nozzle. Although calculated in a different order from Sample Case No. 1, the same data points are calculated and the same solution is obtained. The remainder of the output is identical to the output for Sample Case No. 1.

Sample Case No. 7 illustrates the use of the variable step size options on the initial-value line and the throat downstream circular arc contour. In all other respects, Sample Case No. 7 is identical to Sample Case No. 1. The value of Δy on the initial-value line at the nozzle wall is to be one half as large as the value of Δy adjacent to the nozzle axis (DYRATIO=0.5). The value of $\Delta \theta$ adjacent to the throat attachment point, point A, is to be twice as large as the value of $\Delta \theta$ adjacent to the nozzle throat (DARATIO=2.0). The data deck for Sample Case No. 7 is presented below.

SAMPLE CASE NO. 7. NOMINAL CASE WITH DYRATIO=0.5 AND DARATIO=2.0 \$DATA MODE=4, IWALL=1, NI=11, NT=15, KWRITE=1, DYRATIO=0.5, DARATIO=2.0, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696, YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, AF=0.0, BETA=30.0 \$

Sample Case No. 8 is the same as Sample Case No. 1, except that the x coordinate of the end of the scarfed nozzle extension, point F, is specified instead of the scarf angle β . The same scarfed nozzle extension is obtained for $x_f=19.02549776$ in. The only change in the input data deck is the deletion of BETA = 30.0 and the addition of XF = 19.02549776. The data deck for Sample Case No. 8 is presented below. The output for Sample Case No. 8 is the same as the abbreviated output for Sample Case No. 1.

\$AMPLE CASE NO. 8. NOMINAL CASE WITH XF=19.02549776 SPECIFIED

\$DATA NODE=4, IWALL=1, NI=11, NT=15, KWRITE=1,

G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, EPS=10.0, AF=0.0, XF=19.02549776 \$

Sample Case No. 9 considers the same basic nozzle and operating conditions as Sample Case No. 1. In Sample Case No. 9, however, the scarfed nozzle extension is a cone with an angle $\theta_{\vec{k}} = 5.0$ deg (AF=5.0). The data deck for Sample Case No. 9 is presented below.

\$DATA MODE=4, IWALL=1, NI=11, NT=15, KWRITE=1,
G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,
YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, EPS=10.0, AF=5.0, BETA=30.0 \$

The output for Sample Case No. 9 is identical to the output of Sample Case No. 1 through the basic nozzle flowfield. The oblique shock wave and the flowfield in the nozzle extension are different. Append. E presents that portion of the output downstream of the basic nozzle.

Sample Case No. 10 illustrates a scarfed conical nozzle where the cone angle of the basic nozzle θ_a and the cone angle of the conical extension θ_f are identical, in this case 15.0 deg (AA=15.0 and AF=15.0). Thus, the flow turning angle at point E, the juction of the basic nozzle and the nozzle extension, is zero. In that case, the oblique shock wave emanating from point E becomes a Mach line, and the flow properties across the junction are continuous. The data deck for Sample Case No. 10 is presented below.

SAMPLE CASE NO. 10. NOMINAL CASE WITH AA=AF=15.0

\$DATA MODE=4, IWALL=1, NI=11, NT=15, KWRITE=1,
G=1.2, RG=65.0, PS=1000.0, TS=50CO.0, PA=14.696,
YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, EPS=10.0, AF=15.0, BETA=30.0 \$

The output for Sample Case No. 10 is identical to the output for Sample Case No. 1 through the basic nozzle flowfield. The flowfield in the nozzle extension is different. Append. F presents that portion of the output downstream of the basic nozzle. Note that the oblique shock wave emancting from point E, the junction point between the basic nozzle and the conical extension, is basically a Nach line. The strength of the oblique shock wave does not remain infinitesimal, as it theoretically should. This discrepancy is due to numerical errors introduced by using the rotational flow method of characteristics downstream of the pseudo shock wave to calculate an irrotational flowfield. However, the errors are relatively minor.

Sample Case No. 11 is identical to Sample Case No. 1, except that the basic nozzle has a length of 28.51233074 in. (XE=28.51233074) and there is no nozzle extension (MDDE=1). Consequently, the problem being analyzed is simply a 15 deg conical nozzle 28.51233074 in. long. The flowfield in this nozzle is identical to the flowfield in the basic nozzle and nozzle extension considered in Sample Case No. 10. The output is slightly different, however, since the present nozzle does not have a special junction point at x = 8.07104661 in. A closs comparison of the output from the two sample cases clearly shows that the flowfields are essentially identical. The data deck for Sample Case No. 11 is presented below.

SAMPLE CASE NO. 11. CONICAL NOZZLE CORRESPONDING TO CASE 10, MODE=1 \$DATA MODE=1, IWALL=1, NI=11, NT=15, KWRITE=1, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, XE=28.51233074 \$

The output for Sample Case No. 11 is identical to the output for Sample Case No. 1 through left-running Mach line 29. The flowfield downstream of that Mach line is the same as the flowfield determined in Sample Case No. 10. The right-running and left-running Mach lines are not identical, however, since the special point at x = 8.07104661 in. considered in Sample Case No. 10 is not considered in Sample Case No. 11. The flowfield downstream of left-running Mach line 29 in Sample Case No.

13. SAMPLE CASE NO. 12

Sample Case No. 12 is identical to Sample Case No. 11 except that the left-running characteristic network is initiated from the last right-running characteristic from the initial-value line instead of from the last right-running characteristic from the initial expansion contour.

Since Sample Case No. 11 involved a basic nozzle only (MODE = 1), this sample case also involves only a basic nozzle. The analysis of a basic only by emanating left-running characteristics from the downstream extent of the initial-value line is selected by specifying MODE=2. Such an analysis in a basic nozzle is analogous to MODE=5 in a scarfed nozzle. The only change required to the data deck for Sample Case No. 11 to obtain Sample Case No. 12 is to change MODE from 1 to 2. The data deck for Sample Case No. 12 is presented below.

SAMPLE CASE NO. 12. SAMPLE CASE NO. 11 WITH MODE=2

\$DATA MODE=2, IWALL=1, NI=11, NT=15, KWRITE=1,

G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696,

YT=1.0, RTU=1.0, RTD=0.01, AA=15.0, XE=28.51233074 \$

The output format for Sample Case No. 12 is similar to the output format for Sample Case No. 6 presented in Section VIII.7.

14. SAMPLE CASE NO. 13

Sample Case No. 13 illustrates the quadratic wall option for the basic nozzle (IWALL=2). The envelope of the quadratic wall is identical to the envelope of the conical nozzle considered in Sample Case No. 1. The throat attachment angle is $\theta_a=25.0$ deg, 25 points are to be specified along the throat downstream circular arc contour, the nozzle area ratio $\varepsilon=10.0$, and the nozzle length is $\mathbf{x}_e=8.07104661$ in. (AA=25.0, NT=25.0, EPS=10.0, and XE=8.07104661). The same cylindrical nozzle extension considered in Sample Case No. 1 is considered here. All of the remaining parameters in namelist DATA are the same as for Sample Case No. 1. The data deck for Sample Case No. 12 is presented below.

SAMPLE CASE NO. 13. QUADRATIC WALL, AA=25.0, XE=8.07104661, EPS=10.0 \$DATA MODE=4, IWALL=2, NI=11, NT=25, KWRITE=1, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696, YT=1.0, RTU=1.0, RTD=0.01, AA=25.0, XE=8.07104661, EPS=10.0, AF=0.0, BETA=30.0 \$

The output for Sample Case No. 13 is presented in Append. H.

15. SAMPLE CASE NO. 14

Sample Case No. 14 is identical to Sample Case No. 13, except that the quadratic portion of the basic nozzle is specified by the exit lip radius y_e instead of the area ratio ε . The identical quadratic contour is obtained for $y_e = 3.16227766$ in. Thus, the only changes to the input data are to delete EPS = 10.0 and to add YE = 3.16227766. The data deck for Sample Case No. 14 is presented below.

SAMPLE CASE NO. 14. QUADRATIC WALL, AA=25.0, XE=8.07104661, YE=3.16227766 \$DATA MODE=4, IWALL=2, NI=11, NT=25, KWRITE=1, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696, YT=1.0, RTU=1.0, RTD=0.01, AA=25.0, XE=8.07104661, YE=3.16227766, AF=0.0, BETA=30.0 \$

The output for Sample Case No. 14 is identical to the output for Sample Case No. 13.

16. SAMPLE CASE NO. 15.

Sample Case No. 15 is identical to Sample Case No. 13, except that the quadratic portion of the basic nozzle is specified by the exit lip angle θ_e instead of the area ratio ϵ . The identical quadratic contour is obtained for θ_e = 3.97861353 deg. Thus, the only changes to the input data are to delete EPS = 10.0 and to add AE = 3.97861353. The data deck for Sample Case No. 15 is presented below.

SAMPLE CASE NO. 15. QUADRATIC WALL, AA=25.0, XE=8.07104662, AE=3.97861353 \$DATA MODE=4, IWALL=2, NI=11, NT=25, KWRITE=1, G=1.2, RG=65.0, PS=1000.0, TS=5000.0, PA=14.696, YT=1.0, RTU=1.0, RTD=0.01, AA=25.0, XE=8.07104661, AE=3.97861353, AF=0.0, BETA=30.0 \$

The output for Sample Case No. 15 is identical to the output for Sample Case No. 13.

SECTION IX

CONCLUSIONS

An analysis is presented for predicting the performance of canted scarfed propulsive nozzles. A computer program based on that analysis was developed. Fifteen sample cases are presented to illustrate the use of the computer program.

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- 8. Reference (4), Sections 7-6 and 7-7, pp. 356-369.
- 9. Reference (1), Chapter 17, pp. 185-266.
- 10. Reference (4), Appendix A, pp. 669-694.

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APPENDIX A

OUTPUT FOR SAMPLE CASE NO. 1

" MOEZIE PERKORRANCE PREDICTION PROGRAM.	
THIS PROCKER WILL AMALTER THE FLOWFIELD AND DERFORMANCE OF PROPULSIVE NOZZEES FOR SCHERAL OFFICHS.	::
MODE 1. INKOTATIONAL FLUG ALGHG AIGHT-RUNNING CHARACTERISTICS.	
HODE 2. IRROTATIONAL FLOW ALGHG LEFT-RUNHING CHARACTERISTICS.	
NOSE S. FLOW WITH AR ENFEDDED RIGHT-RUNKING OBLIGUE SHOCK WAVE.	
RODE 4., FLAW THE WOLLE EVIENSION.	
THE PROGRAM WILL KWALTZE THE PERFORMINCE OF A COMPRESSED PROPUESTYE WOZZLE. ICHP = 1.	
THE PROGRAM CAM DETECT AND TRECK AN EMMEDDED RIGHT-RUMMING OBLIQUE SHOCK WAVE (NODE = 3 ON 52. THE FLOWFISLD AHEAD OF THE CHOCK MAYE IS ASSUMED TO BE ROTATIONAL.	
SCHOOL OF RECHANICAL ENGINEERING.	·*-
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#DD?	20-059	15.867	11.363	9.245	7.483	6.034	4.855	3.900	5. 1.70	2.509	2.012	1.602	1.265	986	•753	547	• 365	.199	+062	003
11	5000	5000	5006-	£000	5000	\$000	5000	5000	5000	5000	5000	5000	5000	5000	5000	5000	5000	5000	5000.	5000
-	1000	1000	1060.	1000	1000	10001	1000	1000.	1000	1000	1000	1000	1000	1000*	1000	0000	.000	1000	1000	1000
1	4568	46.05.	£414.	441.8.	4417.	4413.	4407.	4599.	4589	4579.	4569.	4357.	4546.	4554	4322.	4368.	4293.	4274.	4247.	4211.
RHO	.22534	23519	+23759	.25844	425827	.23724	. 23557	.23343	*23098	.22832	* 2255E	.22271	47.45.50	-21682	.21374	21045	+2067E	.20210	19601	.18776
b	444,242	467.654	473.169	475.412	475.014	472.559	468.564	463.470	457.627	451.504	664.B21	436.047	451-167	424.167	416.955	409.264	400.551	\$89.626	575.804	356.905
11674	0000	1.253	1.620	106.1	2.098	2.226	2,288	2:323	2000	2.242	2-161	2.0%	1.961	1.847	1.699	1.524	2 - 509	1.024	4615	000*6
VARG	\$950.	3864.	58.56.	\$82A.	1826.	5839.	5859.	\$685°	5415.	34400	59AL.	4016.	4052.	A 0 46.	4126.	41.66	4215	4275e	4.546.	** 55.
•	1.203	1-162	2.153	1.11.9	1.145	1.154	1.764	1.169	10 10 10	1.141	1.20.2	1 2 2 1 4	1.27.3	1 + 2 + 1	2 2 7 5 5	1.757	1.283	1.554	14.55.1	\$ 98 * 3
3-		#3.¢	104.	127.	140.	1691	1552	25.74	1 5.A.	156.	4524	147.	240	132.	122.	111	44.	16.	6.7.	*
2	5781.	30.65	634.	4226	5A24.	\$430.	3356	\$542.	4.4.4.4	2.4.2.4.5	3576	4915.	494.2	4286.	4126.	53650	4212.	4269.	# 4.4 K	• 0 2 5 4
*	60			1		4	Ļ			45.14	525						1 4 4	.101	4668	************
	8 - 4 G G	. 119			٠		4111	2. 12 ml	1880		250	149.	* * * *	F 6.2.	e 25.	**************************************	. 517	405	, e , e ,	. 701
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jia.	من جن حد الت	=		1 1		<u></u>	7	-4	<u>.</u>	-4 -44	 		-1 -1	**	*	 		=	<u></u>	÷4

IRRCTATIONAL FLOWFIELD ALONG RIGHT-RUNNING CHARACTERISTICS EMANATING FROM THE INITIAL EXPANSION CONTOUR.

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	L	3832.	3302	2441	2000	1807	1558.	1347.	1167.	1015.	884	774.	67.6	594.	519	451.	4.8.7	324	256	182.	103.	35.	1.	-1-		=1.0162	3632.	1035	2867	2506.	2178	1893.	1655.	I 64 a	1271.	1115	94ª.	875.	.111.	.069	h12.	541.	• ()	*CO*	31.5	25.5	69	26.	•	-1.	=1.0162		2826
1001		17 104	A BY	12,775	10.969	9.426	011.9	6.992	5.042	5.237	4.551	3.972	3.466	3.026	2.637	2.2PS	1.954	1.627	1.281	+06.	.50R	• 169	0.0	007		F. 1 A I	20.069	17.4.1	15.114	13,102	11.366	9.875	8.593	7.445	6.556	5.752	2.052	7/4-6	3.456	100.5	66.	5 173	2000	1,647	1.551	.759	.330	.118	920.	300	ETAI	20.00	40000
=		9000	5000	5000.	5000	5000.	2000	5000.	5000	5000.	5000	5000	2000	5000.	5000	5000	5000	5000	5000	2000	2000	5000.	5000.	5000		206-16	5,000.	5000	5000.	5000	5000	5000	.000S	5900.	5000	5000.	5000	5000	2000				000	5000	5000	5000	5030.	5000.	5990	5000.	187.304	0003	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
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-	4 6 6 4	4290	4299.	4301.	4299.	4293.	4285	4275	263	4250.	4237.	4223.	4209.	4195.	4180.	4165.	4148.	4129.	4105.	4072.	4021.	3943.	3876.	3614.	001	Taci	4243.	58	4265.	4267.	4263.	4256.	4247.	4235.	*222	-504	. 194.	*180	4160.	1100	2117	11014	1017		4021	3968	3888	3818.	3753.	3693.	ISPIO	43.64	
RHO.	2696	20600	-20809	.20872	.20822	.20686	-20485	NI.	.19960	~ •	.19355	.19046	.18731	.18414	.18094	.17766	.17416	-17018	•16526	.15865	.14903	.13516	N.	.11437	100 045	?	.19509	.19853	.20020	.20048	•1996B	-	• 19585	.19322	.19031	18723	10481.	18092	1///1	17130	12700	16447	15068	15557	4899	13946	.125A9	.11504	.10562	.09735	190.946	1001	• • • • •
Ь	486.087	398-893	403.758	405-240	404-075	400.893	396.227	390.509	384.088	311.234	370-147	363.079	355.872	348.662	341.407	333.984	326-105	317.192	306-225	291.571	270.493	240.578	17.06	196.875	169	- 1	373.685	181.59R	385.466	186.108	384.255	380.525	375.432	365.389	362.727	355.701	346.491	241.556	234 112	110 656	115.57	304.407	262.796	284.796	270.375	249.775	220.917	196.272	176-954	162.270	1SP =	160.105	1000
THETA	000	3.806	4.417	4.873	5.208	5.445	5.603	5.697	7.67	96/96	6699	5.637	0.55f	5.444	5.321	5.180	5.015	4.812	4.539	4.1.56	(1) (1) (1) (1) (1) (1) (1) (1) (1) (1)	2.217		00000	1.0060		4.000	4.839	5.461	5.933	6.279	6.524	6.68B	6.785	6.828	6.927	16.1.0	621.9	24000		4.064	6.100	7000	5.615	5.209	4.515	3.276	2.156	1.051	0.000	=1.0060	000	3
VMAG	4570	4221	4196.	4108.	4194	4211	4256	• Z bb •	1000	4000	*27.	4415	•	4495.	4536.	4578.	4623.	4675.	4739.	4827.	1964	5150.	5311.	5407.	11 14 14 14	1	4357.	4314.	4233	4200	4300.	4320	434P.	4341.	4417.	* 00 00 00 00 00 00 00 00 00 00 00 00 00	44700	9000	0.74	46.61	4704	4750	4803	4848	4958.	50.89	5204.	5446.	5593	5728.	ETAF =	4 4 2	• 11
Ŧ	407	1.287		1.275		1.293	1.202	1.503	1.510	1.360	7	1 - 3 - 6	1.5/1	1.386	1.401	1.416	1.433	1.452	1.477	1.510	1.550	1.637	1.703	1.764	1,000		1.335	1.320	1.312	1.311	1.315	1.372	1.3*2	1	1.157	۰,	0000	7	0 1 4 1	1.447	1.46.1		1 2 4 0 1	1.526	55	1.613	0 9	. 75	Ç.	a.	3809.2	12366	
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n.	6266	4212.	4183.	4173.	4177	4192	4615	2010	4217	1010		4 5 9 4	***	2	4517.	260	909	4659.	4724.	- 1	B 4 6 4 1	146	310	3	7 7		4347.	4299.	4274.	4267.	4274.	4292.	431R.	4.350	4386		9001	4000*	45.00	46.32	4676	4723	4777	4 X 4 X 4	4937.	5074.	275	5442.	265	5728.	F10 =	4425	- 1
-	1.660		.861	. 799	.742	069	7 9 9	550	800.	220		40.4	104.	30.	. 393	. 355	.331	400	.272	•252		101	•	0000	36.32.1		1.000	.932		810	757	807	665	229	582 184	100.	100	1 4 7 4	- 1	, t		369	343	•312	273	.221	.152	- 095	•	0000	3832.1	1.000	• 1
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٩	365.206	368.210	368.134	365.702	361.51R	149.773	342,927	335.777	32P.499	321.322	314.076	306.882	299.689	292.370	284.640	275.93R	265.289	251.158	231.104	203.307	181.577	163.174	147.510	133.381	: dSI	2	343.606	6.0	ω'	14	348.216	3	91	7		-!o		0 4	, ,	ع ا ذ	4 0	266.358	25.7.806	247,368	233,569	1-	187.366	ĺω	6	• 03	20.88	ć
THETA	5	6.500	9.6	34	7.592	7.959	(CO. /	7.903	7.866	7.804	7.715	7.607	•	7.337	7,159	966	6.679	6.251	5.556	4.314	3.199	60.	3	0.00	-1.0060		00009	.87	53			æ	8.822	5	•	٠,١	8.426	9	: `	9 4	א ה	319	8.007		7.295	S	333	2	10	0	9	0.00
VKAG	44.04		89	-	4424	000	45.28	89	4610.	4651.	693	4735.	7 78	825	4869.	4923.	4990	5040.	5214.	4 09	5573.	5720.	5855.	5962.	. 145		100	0644	4 78	4 82	498	S	4556.	S.	633	2 2	717	200	* C = C = C = C = C = C = C = C = C = C	E 0	.06.94	r 0	5037	> i	5197	332	5.28	9	840	16	6103.	~
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>	΄σ	497	*	62	585.	3	• • • •	628.	; =	631.	000	627.	622.	616.	608	597.	580.	. 400	5	*01.	310.	203.	90	•0		•	473.		-	56	657.	on l	66		N.	رايح	α •	9	ń.	2		2 5	702	שע) c	: =		17	9	=		•
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>	310	4.5		171	.725	28	•		100	520	464	4.7.1	6.4	426	_	378	349	312		_	38	680	φ.	000		5852.1	000	6.6	882	8.30	783	7.4	0	99	*	3	57	• (52	G .	-	<u>وا:</u>	55.	714	7 d	9.8	233	29	130	0.85	.042	
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MOOT	ETAT	20.059	17.745	15-687	12.323	10.958	9.769	7,830	7.043	6.356	30/00	4.749	4.315	3.911	3.521	2.679	2.164	1.564	• 934	* 556 * 1 +	15.1	0.00	• 005	011	ETAI		20.069	17-832	13.847	12.591	11.262	10.100	4004	7.416	6.733	6.138	2.603	0 + 6 6 8 3 ×	4.273	3.875	54.45.3 1.1.1.1	3.007	1.838	1.15A	•736	• 452
A.A.	1167.911		2	5000		5000	2000	S	2		1	ים י	S.	اری	5000	2	2	80	20			500	18	0	187.914		اء	5000	2000	5000	5300.	5000	0000	5000.	5000	5000	5000	5000	5000.	5000.	5000.	2000	5000	5000	5000.	5000.
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RHO	190.950	.17568	- (*	17811	17	27.77	.16881	.16563	16235	15578	15251	*14927	.14604	• 1 42 75	• 13527 • 13536	.13057	•12420	11516	0 2 0 7 0	08420	.07682	.07027	• 064 39	*0550*	190.952		•16965		17128		.16729	16452	15822	.15490	15156	14850	14179	1.458	.13531	•131A7	12800	11702	.10817	165501	• 08538	.07822
a	= dSI	329.525	334.702	335.014	331.655	120.70	314.122	307.035	299.751	285.263	276.092	271.019	263.987	705.655	241.003	230.802	217.362	176.301	100.211	136.342	122.122	109.729	α.	89-107	ISP =		316.011	321.188	319.650	315.915	310.743	297.784	290.629	283.322	276 -011	261-792	254.803	247.884	240.891	233-557	2235363	202.454	184-135	150	o; .	124.790
THETA	=1.0060	7.000	7.889	-07	9.442	• 1	6.6	0	9.018		9.821	9.709	9.578	9.255	9.040	R.749	8.516	1600	1000	4.103	0	9	6	0.000	=1.0060		8,000	0 • 4 CJ	10.110	10,484	10.747	11.02.2	11.065	11.061	11.020	10.657	10.742	때	10.457	10.280	10-050	9-324	8.590	7.330	6 8 8 9	980 • 6
VMAG	ETAF	4604	4566		4592		4693	4735	872	4865	4909.	4953.	4998	5093	5147.	5216.	9	1 to 1 to 1	ລ ແ	955	ပ	6218	~ ? .	6451	ETAF		589	66.1	4661.	46.92+	4713	47.40	8 33	ac i	9922	5012	10	102	5149.	5197	5422	5416.	5553.	5751.	יות	0 0 0
Z	\$6095	1.426		1.414	1.421	1.444	1.459	5 4 4	1.501	1.525	1.542	1.559	1.577	1.614	1.636	1.664	702	. A. A. A.	91	•	90.	2.110	• 17	22	\$8 19.4	1	4 i 4) P)	•	4	1.467	1.496	1.512	1.523	1 2 2 2	5.02	8	5	Ç,	9290	4.75.0	1.7.7	•	1.992	4	•
>		561.	680 a	721.	753.	798.	414	-	A36.	838.	837°	35.55	222	0	909	793.		621	527.	426-	323.	218	110.	•			ار. د ایر	775.	Ø.		\$ 1.50 0.00 0.00 0.00	916	928.	936.			942.	-57	934.	27.7	51.6	n 7 a.	829.	734.	2 2 2	• 600
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THETA	4.023	2.987	1.974	00000	=1.0960	000°6	9,920	16.621	Ξ.	11.521	11.960	12.060	12.102	12.096	17.032	11.894	11.766	11.630	11-475	11.234	10.769	10.322	9.578	A.311	7.164	5000	4 - 9 - X	2.9.9	1.944	•965	00000	=1.0060	10.000	10.933	11,644	12.172	41.	12.993	13.092	13,132	13.123	13.077	13+003	12.702	12.642
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٩	#/B ##X	211-855	0.000	187-691	175.438	158.508	135,950	118.800	104.639	2.	62.423	'n	65.601	58.642	52.457	46.934	I SP =	0	c	280.434	277.514	272.909	267.131	260.587	253,591	246.388	239-155	232-013	27.5-13.5	211.21.2	205-188	198.646	191.831	164.269	175-163	165.520	0 0 0 0 0 0 0	105.201	95.823	84.587	74.53	66.591	56	52.816		6 T D = 2 *		≡ dSI	646.336	268.881
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HDOT	EVAT	7.852	9.588	10.603	11.576	12,765	13.870	14.969	16.051	17.107	13.130	20.056	ETAT	1 176	9 1 2 3	9.00€	9.972	10.994	12.048	13.11*	14.14	16.25	17.257	18-211	19.1Fn	20.066	ETA1 :	6,945	7.627	8.460	9.381	10.36.	12.466	13.438	14.45	15.453	16.432	~	X [(26-1179	기 :	ETA!	505-9	7.155	1.940	8.814
11	216.846	5000.	5000				5000	2000		Ţ		2	218.194	Ι.					5000.			5000		5000		5600.	219.465	5000	5	- 1	KC I	- 1		1				un i	ا'			220.684	5000			1
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_	142	153 3724.	ı	1		1		253 3599.		57 5655		36		9854 170			;		i	UZ 5698.				.0		30 3579				ı	9902 3712	3705		70 3665			- [33 3562 39 3545				ļ	13 5693.	
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a	15P	170.655	178.398	177.810	175.561	172,259	168.313	104.02	16641	1.00	146.578	142.464	ISP	164.639	170.699	172.930	172.566	170.516	167.391	1103.011	155.164	150.837	146.576	142.491	138.475	134.581	ISP	158.772	165.099	167,575	167.427	165.069	154.995	154.995	150.823	146.613	142.470	138.489	110000	127.047		ISP	152.939	159.517	162,229	162,293
THETA	= .9822	17.895	15.177	14.506		7		13.004		14.500		15.000	= .9830	17.756	્ક	14.975	14.244	13.884	369-61 3 × × × × × × × × × × × × × × × × × × ×	13.600	1.5.9.50	14.013	14.231	14.469	14.728	13.000	= .99 ₹#	17.590	15.389	14.756	•	4 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	13.375	13.427	13,554	13.735	13.951	14-188	0 1 2 7 1	15.000		5 486 =	17.421	15.677	14.513	13,793
VMAG	1	5659.					5713	ļ				5859	ETAF	5708.			5644	5560	:		1	5825	1	- :		29/16	ETAF	57.7		- 1		1	5755	į	- 1		- [5935	2007	6042		FTAF	5807.	5751.	5728.	5728.
2	4395.8	1.851	1.82	1.82	101	10.24	1.87	000		6	8	95	4423.2	1.973	. A 5	. A.	e :	er e	1.877	- C	06	•	6	95	6 0	7	6.044	0	1.87	١	1.84.3	ויי' ענ נ∶ם	1.834	5	읾	4			5	2.026		4473.7	6	σ	~	1.881
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HOOL	9-751	10.720	11.723	12.721	13.714	15.634	15.543	17.470	18.316	19.186	20.065		- 1	560 * 3	4.694	7.427	A-253	. 9.147	10.05	1000	12.06	13.969	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	10	16.580	17.446	18.294	9.15	20.065	1413	11 = 5 & 5	5-190	5.876	•	P.503	50 × 60	10.312	11.240	12.164	13.076	14.827	15.676	16.509	17.336	16.176	9.0	20+065	= 1 ¥ 1 3
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CH	.09653	109514	-09339	09144	. 0 4 4 5 4	.08525	·	.08128	67975	.07745	.07555	DCB. 000	7.00	.0A961	.09307	• 0 4 5 5	6	. 09395	*0760*	0000	08710	. 0 A 5 0 7	.08306	.68111	. 07918	.07729	.07543	.07357	.07163	222.211	~	0	.09146	017	910	• 08986 • 08986	• 08828	14920	- 08455	86280.	07872	.07684	.07500	.07319	.07137	.06948	.06737	223.729
a.	160.625	157,896	154,381	25.051	140.45	138.370	134,492	130.674	156.965	123.333	116.707	48	i	146.313	153.737	156.686	156.965	504.641	146-592	145.005	141.984	138.030	134.115	130.347	126.632	. 123.022	119.485	9.95	112,299	= dSI	140.286	147.362	150.563	151.075	145.816	14/.411	144.284	16/-041	10/00/21	129-411	125.765	122,166	118.666	115.233	111.904	2	104.325	= d\$l
THETA	13.364	13,151	13.090	200	13.440	13.655	13,390	14.148	14.418	14.701	15.000	7386 =	1	,	5.437	4.251	. 204	13.069	12-01-0	12.819	12.933	13.115	13.324	13.564	15.821	14.091	14.374	٠	15.000	- 9860	15.966	15.145	13,930	13-153	12.739	17.4.51	17.403	266-21	13 7 16	12-948	13.182	13.439	13.799	13.992	14.292	14.619	15.000	.98€
٥ ٧ × ٧	5742.	5755	5794.	270	58.99	5936.	5972	6008.	6043.	£013°	6115.	FAF	,	5859	68 00 e	5775	-2175		7. 0.7	1		1	- 1	11.	6046.	0	117	6153.	6141.	ETAF	5919.	5855	58.7	5823.	28.54	ירו ארו היים מים	5000	• • • • • • • • • • • • • • • • • • • •	4 9 2	6020	6055	.0609	6125.	6160.	6196.	N C	6276.	EIAF =
3	1.99.8	1.898	÷ 6		1.35.9	1.976	目	2.639	듹	ů	500	4497.9			֭֭֡֡֓֞֜֓֓֝֓֓֓֓֓֓֓֓֓֓֓֓֡֓֓֓֡֓֓֓֡֓֡֓֓֓֓֡֓֓֡֓֡֓֡֓֡			1	. 6	96	96		8	. 31	20.	2.044	9	•	7.007	4522.R	1.748	.6•	9	1.924	7	100		1.000	1 998	2.015	2.032	2.048	2.055	•	σ .	-	-	4550.0
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ב	5586.	56162	5644	9.5	5738	5748.	5	£.	453	80	5906	F10 :			591	760	6110	5,44.9	_	5722				5.043.	871	66	n, c	ר מ	ž.	E UI	5651.	25	، ام		5717	- 1/		5806.	۔ اہ		•		5951.	5977.	000	7 5	0	F10 =
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		7	1.280	-	1.439	4	3	1.581	S	.67	1.714	12°		0		-		400	. 33	38	4	65.		5	*	œ.	7	. ;	V.	f. H	1.657	1.103	1.157	1.614	1.531	1000	1.447	1.503	1.557	1.509					₹ 0		•	, , ,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,
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		• 9607	1206.	1512	1692.	1882.	2080.	2281.	2482.	2878	3068.	3257.	3443.	3628	381/	424	4525.	1166. =		1070	1196.	1344.	1508.	• 689	2057	2247.	2436.	2803.	2983.	3161.	3339.	3321	3932.	4504	4571.	8866° =	R21.	895	1000	1127.	1272.	1424.	1594	1939.	2113.	2286.
NOOT	. 82.4.4	\$0.50g	250	6.972	1.764	8-604	27407	10000	12.110	12.967	13.793	14.612	15.416	17-010	17.891	18.864	20.064	ETAT	100	459.4	5.485	6.135	7.531	1000	9.260	10.088	10.911	12.504	13.282	14.049	14.814	16.417	17.351	18.506	20.053	ETA1	3.748	4.067	4.523	5.078	5.707	6.391	7-856	8.604	9.361	10.105
E.	2000	2000	5000	5000	5000	5000		000	8000	5000	50000	5000	2000		000	5000	2000	556.049	5000	5000	.0000	5000.	3000	5000	5000	5000	2000	5000	.0005	5000.	5000	5000	5000	5000	• 0000	228.165	5000	5000	5000.	5000.	5000.	5000	5000	5000.	5000.	5000.
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-	3570.	3602	3617	3620.	3616.	3607	3579	3563	3547	3530.	3513.	3496.	26.25	3444	3426.	3404.	3378.	154	3527	3558.	3576.	3581	3569	3557	3543.	3527.	3311	3477	3460.	3443.	3409.	3390.	3369.	3343.	2201	ISPID	3453.	3495.	3515.	3523.	3522.	3515.	3490	3475.	3459.	3445
RHO	.08216	.08595	.08773	.08813	-08762	10800	. 08330	.08146	.07957	.07768	.07585	.07403	07070	.06872	.06688	.06483	.06237	725-527	.07688	.08088	.08287	.08345	.08212	.08076	.07916	.07743	07.78	.07209	.07035	• 06864	.06526	.06350	.06153	.05916	90000	227-820	.06963	.07389	.07613	56570.	.07681	.07604	.07342		07019	• 06852
a	132.384	139.738	143.226	144.010	143.003	140-832	134.589	131.029	127.391	123.769	120.270	116.812	110-140	106.837	103.411	99.629	95.102	ISP =	122.242	129.909	133,745	134.868	132.309	129.684	126.607	123.290	116-466	113.161	109.88R	106-695	100.422	97.166	93.569	84.263	21.00	ISP =	108.541	116.546	120,805	122.362	122.099	120.630	115.668	112.687	109.589	106.469
THETA	16.631	14.763	113.511	12.718	12-247	11.923	11.958	12.074	12.247	12.459	12.694	164-21	13.506	13.806	14-135	14.517	13.000	- 9877	16.139	14.207	12.907	12.081	11.333	11.243	11.273	11.387	27.11	12.008	12.266	12.539	13.12A	13.458	13.843	15.000		.9888	15.339	13.315	11.947	11.077	10.557	10.284	10.213	10.326	10.501	10.717
VHAG	5992.	5924.	5892	5885.	5894	5940	5971.	6004	6039.	6074.	6109.	61 70	6213	6248.	6286.	6328.	• n + c +	ETAF	60 H9.	6615.	5979.	2968	5992	6017.	6047.	6079.	6147	6182.	6216.	6250	6319.	6356.	6398.	6518.		ETAF :						5105. 6128.		6186.		95259
E	2002	1.970	1.956	1.953	1.957	16	1.992	2.008	2.024	2.041	860.2	2.001	2.108	2.125	2.144	2.165	76197	1582.4		2.013	•	1661	2.00.2	6	2.028	2.043		. 093	-109	170	.150	.179	200	263		624.1	1		- 1		ı	2.067	1	1		- 1
>	1715.	1509.	-	1296.	1250	1227.	1237.	1256.	1281.	1310.	1306	1413	151	1491.	1535.	1586.	•1091		1691	1476.	1335.	1249	178	1173.	182.	1225.	254.	1286.	1321.	198	1435.	614	31	63			648.	416.	264.		000	084.				
>	5741.	5728.	5729.	5741	5784	812	5841.	5872.	5902.	5931.	2750.	5015	6041.	6068.	96	6126.	• 70 70	F10 =	5849.	5831.	3828.	5855	5875.	5902.	5930.	5989	6018.		_	6127.	_	6182.	212	6296.		F10 =						6031. 1			6115. 1	.
_	•656	-702	.755	.810	9 6 6 6	116.	1.031	90	.133	181	-1	316	1.359	1.403	1.450	1.502	T	4526.1	i	677	701	- 1	891	945	666		.149	.195	~	329	.373	20	473	628		4572.1	605	642	694	7 4 6		906	.953	*00	c m	
*	1.124	7	1.223	1.280	1000	1.460	•	1.577	1.633		780	•		93	. 98	2.043	-	11	~	1.258	7 1	?	1.492	1.554	1.616		1.791	•			• 1	7	2.169	. ~		"	1.343	1.000	1.495	1.553	1.623	• • :	•	1.818	1.940	
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HDOT		10.828	11.550	12.263	13.704	14.478	15.357	16.446	17.920) I	ETAT :	2.836	3.045	5.587	P M M M	4.906	5.516	6-153	6.806	7.463	A 7 5 0	0.00	10.041	10.6R2	11.34	12.041	12.841	5.19	17-170	20.056	ETAI	- [2.204	• i •	2.953	.37	3.862	4.387	• 1		6.694	7.250	7.840	d	9.002	9.602	10.980
F		2000	5000	5000	5000	5000	5000.	5000	5000	•	30.901	5000	5000	2000	5000	000	5000.	5000	5000	0000	2000	2000	5000	5000	5000.	5000.	5000	5000	5000	5000	.34.582		5000.	2000	000	5000	5000.	5000	5000	• • • • • • • • • • • • • • • • • • • •	0000		5000	5000	5000	2000	5000
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<u> </u>		5426.	5409.	3375	3357.	3339.	3317.	3291	3255.		18P10	3352.	3401.	34.28	3441	3435	3426.	3413.	5599.	3383.	1361	3330	3318.	3301.	3283.	3265.	3243	1181	3128.	3053.	UIdSI		3265.	33630	3369.	3373.	3369.	3361.	3349	2226	3361	3289	۱۸	3256.	3239.	3222.	3182
вно		* U65 89	• 06526	.06266	.06048	.05882	.05697	4/4C0.	.05183		230.882	.06003	0 (01/90*	.06836	.06785	.06691	.06570	.06433	.06287	05130	448504	• 05698	.05554	•05409	.05257	.05088	1440	.04249	.03761	234.765		2/250*	71000	.06150	.06186	.06155	.06090	95	00000	42100	.05455	.05320	.05187	.05054	.04920	.04622
 a- 		5 (100.418	94.569	91.663	88.545	85.308	41.311	69,195		= dSI	90.839	0	103.825	106.164	105.214		101.226	48.695	96.016	90.620	87.953	85,335	82.752	. 80.160	77.464	74.479	66.283	•	1.83	ISP =	•	0 1 - 1 - 1	0 -	93.522	94.167	93.601	.92.231	90.348	901.00	20.00	0 0 0 0	78.588	76.226	73.888	71.537	66.375
THETA		10.956	11.219	11.787	12.095	12.430	12.821	14.514	15.993	וי	- 9899	14.018	11.862	66501		.62	•	8.553	• 1	8.857	9-115	609	4.897	10.199	10.517	10.864	11.266	12.471	'n	2	= .9908	- 1	12.830	5 6	8.022	7.442	7-144	7.044	7.0 A1	4.513	0100	7.915	A-203	8.504	8.818	9.149	9.925
V *AG		6256	6319.	6326	6420.	6457.	6498.	6569	6617.		ETAF	6430.	6334.	6241.	6256	6246.	6285.	6310.	6559	6369.	6434	64.65	6498	6530.	6564.	6599.	6640.	6757	6853	.0669	ETAF	- 11	5246	6427	6398.	6391.	6398.	6414.	64.16.	0000	24.00	6553	6585	6616.	6648.	6681.	6755
! S.	ľ	2.144	-1	2.194	2	• 2	2	٦,			4680.R	.21	∵:		12	7	-	-	٦)،	2.136	۱۰	,	: \?	.2	2	2.306	2.327	31"	2.446	S.	4755.4		? (216	.201	.197	.200	•208	2	•	2 6		1		2.332	~	2.391
>		1145.	1229.	1265.	1345.	1396.	244	200	1738			55	302	e (c	959	0	932.	93A.	456	981.	1043	1079	1117.	1156.	1198.	1244.	1297	500	, A	60			1467	000	893	828.	196.	786.	793.	0 4 4 0	0.70	905	939.	978.		ماد	1164.
>		6171.	198	6223	278	306	6336.		6421.		F10	~	199	2 .	6180.	195	6216.	240	6266	293	6348	6375	6401.	~	6454.	6481.	6512.	1	664	6752.	F10]	-1040	6.148	6336.	6337.	6348.	6365.	6387	2716	4747	6491.	6517.	6543.	6569.	6546.	6654
>		-	1.195	1.240	1.329	1.377	1.431	1.498	1.589		4633.6	.563	.593	20	731	.00	.835	•886	156.	. 1887	1.082	1.128	1.173	1.218	1.263	7	1.367		. 6	8.	4711.5	1	120.	į.		1	- 1				!	.021	.067		.158	- 1	1.310
**		1.998	2.055	2,111	2.221	2.280	2.347	2.429	2.540	:	F :	22	1.562	3 5	1.740	1.809	1.879	1.949	2.019	2.087	2.217	2.280	2.342	2.403	2.465	2.531	2.697	2.921	3.002	92.	11 L	٠,	1.000	1.764	α. •	Œ	1.966	2.041		2,265		2.406	2.474	2.541	2.608		123
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HDOT	11.899	14.998	17.706		ETA1		1,808	•	2.645	• 05	3.513	•1	5.027	5.552	6.070	6+596	7.123	7.656	8.797	9 4 4 75	10.326	11.492	13.217	0	20.052		ETAI	1.345	M) v	1.5.1	10 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	2.418	2.814	3+24 4-6-42	4.156	4.628	5.098	5.575	6.055	1000	7.594	6.221	10 000	11.713
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RHO	.04432	i		•03048	237.551	4669	05147	02440	16660				• 05374	1		i i			• ;	9 Y	1,59	03828	_				239.651	.04158	•04643	이	01100	.05184	นกเ	.05063	기약	04748	.04635	• 04519	+0+0.	04289	54040	.03912	91	• 03228
a.	63.119	53.198	45.767	40.275	ISP =	67.184	75.528	80.710	83.4408	94.100	83.039	81.453	79.545	76 307	73-130	70.953	68.803	699•99	64.519	52.277	56.801	52.941	7.69	↑0.878	35.956	2001	ISP =	- σ	66.748	₩	D - N - 4 - 1	و اه	75-373	74.040	70.514	68.557	ما	-			** B • 6 9 B		100	48.080
THETA	10.448	12.220	13.749	15.000	1066.	11.694	S	7.661	약	5.729	5.636	5.687	5.836	e i	6.591	6.895	7.212	7.541	7.886	8.260	0 0 0	9.971	11.053	12.614	13.887	15.000	0066* =	4	a.	6.359	5.282	4.366	4.286	4-357	4.527	4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4	5.346	5.670	6.005	6.351	6.713	7.552	8.133	9.986
VMAG	1	.l.			ETAF =	6743.	6626.	6557.	6522	10 10 10 10 10 10 10 10 10 10 10 10 10 1	6527.	6547.	6572	6600	6667	6689,	672P.	6751.	67.83	6.918	0000	6971	7065.	-1	30	74 05	ETAF	6878.	6750.	6673.	6633	5617	.8.99	9	5669	66.34	6752	6782	6812.	68434	6874	6947.	6995	7059
£	2.418	511	552	.661	4812.2	2.584	2.320	2.283	2.265	862.6	2.267	2.278	2.291	2.306	2.571				2.406	2.426	2446	0.510	2.570	65	0.1	2.786	485A.C	2.66.0	2.3P8	2.346	2 - 32 4	2 2 2 3	2.321	2.331	20163	2.00	2.389	2.405	2.402	2000	2.4.F.	2.590	2.520	2,567
>	234	1330	. 60	967		7.47	1070		753.		641.	•	· ~	6969	72k.	803		امدا	931.	979	900	207	nπ	572	754	ᆐ	н		646	739.	611.	4 4 5 4 5 4 5 4 5 4 5 5 5 5 5 5 5 5 5 5	495°	505	526.	556	629.	670	713.	75.70	90 P	913.	987.	1242
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 		1-475		2.022	4.767.4	000	500	.543	.581	• 625 113	722	.172	.821	.870	.917	1010	1.050	1 - 1 0 0	1.147	1.197	1-253	1.524	1.567	1.786	1.983	2-173	4.409.5	4	- th													197		1.368
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u.	3382	1	4378		2665	247	250.	276.	320-	379.	450.	531.	619	713.	. 912.	913.	1013.	-2111	1222.	1329.	1441.	1562.	1702.	1880.	2126.	2496.	3052.	3541.	*000t	4435.	4655	٠,	= •9977	190		- LB & C	0000	350.		400	581	999	760.	863	946.	1041-	1140.	1243.	1355.	1465.	1651.	1 581
TOOM:	14.114	160215	18-180	250000	ETAI =		199	1.164	1.349	1.598	1.899		2.619	3.018	30,000	3.858	4.232	4.717	5-155	603	6.067	6.570	7.151	7.885	p.899	10.414	12.680	14.673	16.543	18.325	20.047	- 1		195		~ 40	*00	1021	10.0	400	2.00	9.0	200	100 100 100 100 100 100 100 100 100 100	3,978	4.376	4.788	5.215	6.679	6.217	468.5	7.845
ŢŢ	5 30.	5000	5000	20005	241.429	•	5000	.000	5000	, 000 a	5000	000	5000	5000	5000	5000	5000	5000	5000	+500S	5000	5000	5000.	5000	2000	5000	5000	000	5000	5000.	5000		242.782	.0008		5000		5000	-0005		0000		5000		000	\$000g	5000	5000	£ 000.	\$000	5000.	5000
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RHO	. 02830	02531	1	* 020 B3	2414197		81750	.04207	09525	*0.*0	4000	46740	75	30.40	06521	12010	0413	20040	70040	0000	0.020	6.5817	47 47 0	10000	22020	1,200	02512	0 1 0 7 0 °	02101	01910	01748	•	242.225	.03332		.04149	.04341	.04435	.04461	04440	• 04 587	.04312	04225	04130	2000	104010	0.2630	403740	90706	0.000	.03244	€03049
d	36.850	\$2.22	28.526	25-502	186 = .		•	59.296	<, □	67.795	65.208	C + 4 * 6 3	The Bar	200	777653		0000	000	27.22	300	22.623	55-789	21.500		10.0	200	32,366	20000	25. 75.	000000	20.663		ISP =	44.823		58.314	61.564	65-173	63.618	63,255	62.345	61.078	59.592	57.984	56.345	54.662	22-919	51.234	93.286	65.79	43.397	40.288
THETA	11.581	2000	14.005	15.000	06 46.		9-5-52	5.892	52074	3.949	32.326	. O 20	2.972	3.069	20000	5.00	75446	4.166	215-4	4.869	5.235	5.616	6.024	2.4.9	7.074	7.862	9.005	10.635	11.932		000	•	1788. =	F. 6.05	1	\$-763	2.592	1.957	1.687	199*1	1.794	2.051	2.332	2.672	3.028	5.402	5.784	4.173	6.573	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	6-101	61916
V P. A.G	-		⇃.	16.0	1 4 5 5	1	2003	6865.	67.20 e	6754.	6714	6711.	67.19	6775	672be	67.81.	6808	6856.	6865	5845	6925.	6356.	6989.	7027	7975	7140	7232.	7565.	70.70	1007	1522	• 00: 4	ETAF			68.80.	6289	6 R 04 .	6797	6 8 0 2 ·	6817.	6836.	6866.	6886.	6913.	6941.	6970.	6669	70.50	7065	1	72.12.
3			2	406€		22116	2.55.5	l	l	47.9	15.0		42	2,379	2.33	2.405	25.420	436	3	2.470	2:407	2.505	2,525	2. 5. E.	2.576	2.615	20674	2.140	2.0.2	2 · · · · · ·	١.	7 1 1 1	4921.6		100.	2.461		20418		2.417	2+425	2.436	2.450	2.454	2.4.0	2.496	2.513	2.551	24549	A 556	2000	2.651
>	1	# # # D # #	1698	1961			1162.	624.	60%	464.	3892	355.	3484	561.	1254	# 1 B	1554	* 87.	2900	585.	6.52.	6.81.	154.	795.	971.	#76.	1132	1559.	1549+	1715.	1865.	2101.			1 000	26.34		25.5		197	21.3	2.5.2	239	\$21.	XC5.	.23.	450.	-605	561.	616.	650	965
Þ		7136.	ď۷	3 7	:	1	. 200	į.	6754	6719.		6701.	521Ac	6725.	67552	6768.	67.93	6.18.	58342	6876.	6896.	425	5951 e	382	7021	7072.		1239.	邮	372	3	7466.	0.5	1	. 1 . 0 .		0000	V C	* 4	6	44.4	13	Ů,	6.57.50		6929.		1	_			7159.
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×		K C D	20000	100	Í	- E E	•	4 5 2 2	3-16-1	2.222	7.085	2.177	2.454	2.552	2.69.2	2.7.1	2.616	500° 6	986	1.369	1.152	3.5.5	\$ 54.5	3.4.25	A. A. A.	1,727	2.43.	. 156	4.585	* 66 * 1	54.932	585	11		1. 1.	1	27.	*	*			200				4.146		1	-	2	3	3,457
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-	2229.	3220	3657.	4074	4864	= .9963	144.	84.	176.	216.	268.	396.	. 04.8	630	712.	798°	976.	1071.	1295.	1449.	1992	2490	2933.	3749	4133	45.06.	1.9947	105.		121	154.	197.	307	372.	514.	588.	666.	
MOOT	9.267	13.299	15.080	16.781	20.037	ETA1 :	.594	414	.729	899	10116	1.660	1.972	2.644	2.990	0.04×	4.089	4.482	5.407	140.9	8-761	10.286	12.085	15.412	16.990	18.534 20.035	ETAI :	•430		495	.634	. 815	1.283	1.557	2.154	2.466	2.189	30105
i.	5000.	5000	5000.		5000	243.707	5000.	5000	5000	5000	5000	5000	5000	5000	5000.	5000	5000	50000	5000	5000	5000	5000.	5000	5000	5000.	5000.	44.191	5000		5000	5000	5000	50005	5000.	5000.	5000	2000	• 0 0 0 0
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-	2874.	2732	2674.	2623.	2533.	15P10	2917.	3063.	3095.	31111	3116	3109.	3099	3073	5059.	3027	3011.	2994	2954	2927.	2836	2758.	2693.	2584.	2537.	2455.	1591	2855.		3049.	3069.	3077	3072	3063.	3038.	3024.	5009	277.00
RHO	.02780	02157	-01940	.01/PB	.01478	242.794	.02992	.03823	.04027	.04134	.04162	.04119	.03974	.03886	.03795	.03605	.03508	013410	.03189	03046	.02503	.02264	.02010	.01634	.01490	.01367	242.905	.02690		.03738	.03858	.03910	0	.03823	.03670	.035.85	03496	00f00*
۵	36.059	26.596	23.421	18.701	16.893	ISP =	39.391	52.859	اف	58.062	58.535	57.807	55.375	53.907	52.396	49.261	47.682	46.078	42.514	40.249	53-330	28.181	24.433	19.055	17.065	15.189	ISP =	34.669		51.448	53.446	54.299	53.788	52.855	50.335	48.941	47.486	46.013
THETA	8.093	11.101	12.262	14.184	15.000	= .9863	7.521	2.455	1.236	.603	184	.567	1.198	1.578	1.970	2.787	3.203	4 - 0 A 2	4.595	5.227	7.284	A.98R	10.350	12.553	13.465	14.285 15.000	.984₽	6.535		261	871	-1.051	- 703	356		.921	1.058	1 • 6 10
VMAG	7305.	7545.	7639.	7799.	7869.	ETAF :	7231.	6972.	3	6885.	6877.	6889.	6907 . 6929.	6954.	6980	7036.	7065.	1 28.	7166.	7213.	7369.	7501.	7608.	77.96.	7862.	7931.	ETAF =	73.37.		.4669	6962.	6947.	6956	69 72	6993. 7017.	7042.	7069	* LF D
1	2.720	2.882	2.949	0.00	3.121	4-0464	2.673	5	2.491	•				100	• 1		2.570		2.632	2.661	2.762	2.851	2.926		3.116	3.170	4950.2	2.741			2.509	500	505	515	541	l l		
>	1028.	1453.	1622.	1911	2037.		946.	299.	കി	72.	.94	68.	145.	192.	240	342.	395.	507	574.	657.	934	1172.	1567.	1692.	m	1957. 2069.		835.		-32.	-106.	-127.	-86.	-43.	59.	113.	169	*(2 2
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a	17.417	17,175	18.463	16,717	16.138	16.897	14.381	11.904	15,444		1SP =	20.803	19.204	17.267	14.804	11,878	9.918	25.50	4.514	61.63	0.688		5,810	16.903		16.719	3	17.805	16.198	15.584	14.937	14.385	13.887	13.424	12.978	75.35	1SP =		20101	10007	10.873	9.058	7.703	6.658	5.831			5.418	16.056
THETA	-4.927	-4.553	-5.528	-3.631	-2.916	-1-540	1,004	004	0000		.9824	0.000	•595	1.389	2.498	4.063	5.292	125.0	7.275	8.092	5. N. B. W.	`	8.693	-5.057		786		-6.324	-4.321	-3.567	-2.795	-2.144	-1.561	-1.025		0000	= .9824		0000	(8)	10707	4.695	5.752	5.675	7.490			7.925	-6.973
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*	110-111	12,152	12.294	13.034	13.368	130.23	****	3	15.048	1	1 1	62.872	7.131	7.487	8.019	. R2	.53	0.29	0.8	1.48	10		11,396	11.996		•		12.561	1 4 5 3 3 5	13.687	14.058	*		ŝ	4	2	H L		7.404	7.778	× 00 ° 5	9.0.0	10.646	•	11.998			100	12, 193
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	RHO.		.01362		l	1.		l	.01090	242.166	,01280	01121	00794	.00612	.00547		.01304		.01297	.01294	01210	.01178	.01112	.01081	0.010.6	242.166	.00976	.03804	.00597	.00525	
		15.983	15.801	17.071	14.984	13.00	13.355	12.476	12.039 11.569	ISP #	14.22	12.12	8.022	5.868	5.129		17.071		14.942	14.658	13.737	13.237	12.348	11.933	11.058	ISP	10.2.2	6.135	5.688	4.890	
	THETA	-6.783	-6.330	-7.331	-4.364	-2.845	-24225	-1.16	573 0-000	.9824	0.000	1.115	3.965	4.560 5.831	6.694		6.915 -8.257		-7-730	-6.330	-4.512	-3.748	-3.068	-1.659	- 655.4 0 - 0 0 0	9824	0.000	554	3,832	5.503	,
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0005		4	5809.	3832.	•	187.898	190.942	1.0162	
0005		*	3809	4832	1.0060	187-679	190.943	1.0162	
8000			18.09	3632	1.0060	167.902	190.945	1.0162	
2000	1.00004		3609	3832	• •	187.904	190.9+6	1+0162	`
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00	•l	7	78 10 c	*A32	1-0660	187,934	190.963	1.0161	
•	10000	ĭň	188	3832	1.0059	187.940	190.965	1.0161	
•	• •	٦	5018-	5833.	1.0059	187.947	190,968	1.0161	
•		, ,	3410	5833	1.0059	187.954	190.972	1.0161	
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2 2	B) (1	4068	3996	.9821	200-688	199.093	.9921	
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	•		4177.	4.091.	4616.	205.040	203.838	.9893	
•	•	-	.222	4135	.9793	208.284	206-024	•9892	
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1.01136	•	-	44.74	4044	9845	220.684	219-463	.9945	
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	40004		4503	4460	.9860	223.109	222-211	9960	
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8.0710					4600	244.042	242,166	.9923	
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ISP(2-D)	242.166	242-156	242,166	242-166	242,166	242,166	242,166	242-166														
1SP(1-0)	040-440	244.042	244.042	244.042	244.042	244.042	244.042	244.042					!									
ETA(F)		4 C C C C	982	40 d d	4994	4000	4000	9824												•		
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9.256	3.162	23.807	.38 .616	0	242 17 66	173/SE402	1.558	2420266
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12.581	3-162	17.751	79.631	•	242-166	1063£ +02	7.802	292.165
13,162	3-162	16.483	8,5,955	0.	242,166	-, £323E+01	8 .242	2.42.166
13,753	3,162	15-375	92 - 146	•	242-166	**************************************	6.472	242-166
14.372	3,162	14.57	98 06 4:3	0	242 - 256	- 7004E+ 10	13.507	24.24166
15.048	3.162	13.443	105.867	0.	24(1+16/	• 3787£031	8034X	292,166
15-841	3.162	12-523	114.7-1	Ü.	24.20.16 5	6 5 3 4 8 € 0 1	2060,	242-166
16.863	3.162	11.559	127-236	ô	2420156	.1468E>02	7.211	242-166
18.312		10.537	150.418	9	9: 3 0 2 1 2	.2189E+02	6.120	2420166
19.025	3,162	10.265	160.030	•0	242.146	e49345401	5.874	242,166

) 718 H)/LB#		
AFETERS	242,166 (18F-SEC)/18P	206.785 (LBF-\$EC)/LBH 125.174 (LBF-\$EC)/LBH 31.562 0EG		
PEPFORMANCE PLANETERS	ISPXH = ISPXH = ISPXN = E	ISPXH E SECTACET = SECTACET = SECTACET		
LL SCARFED MOZZLE P	66005E+04 LB* 17889E+03 _BF	414998E-04 L3F		
SUMMARY OF OVERALL SCARFED MOZZLE MOZ	FYN =4			
SUM	13		1 1 ,	; }

APPENDIX B ABBREVIATED OUTPUT FOR SAMPLE CASE NO. 1

THE PROBLEM UTLL MARIZE THE FLOWING NAMED AND PERFORMED FOR PROPERTY WORLD TO SETERAL OPTIONS. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHARACTERISTICS. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 606 27 HARDTHOOM, FLOW LIGHT STANMING CHINGS SHOCK WAVE. # 7 HARDTHOOM LIGHT STANMING CHINGS SHOW HARDTHOOM, FLOWING CHINGS SHOCK WAVE. # 7 HARDTHOOM LIGHT SHOW THE CHINGS SHOW HARDTHOOM LIGHT SHOW WAVE CHINGS SHOCK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING CHINGS SHOCK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING CHINGS SHOCK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOK WAVE. # 600 11 HARDTHOOM LIGHT SHOW THE STANMING SHOW THE CHINGS SHOW THE CHINGS SHOW THE CHINGS SHOW THE CHINGS SHOW THE STANMING SHOW THE SHOW	MO22LE PERFORBARCE PREDICTION PROGRAM.	
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#60E 7: TRROTATIONAL FLOW ALGOR CERT-RUNNING CHARACTERISTICS. #60E 1: F.GW WITH AN EMBEDDED ALGOR-RUNNING DOLIDUE SHOCK WAVE. #60E 4: F.GW WITH AN EMBEDDED ALGOR-RUNNING DOLIDUE SHOCK WAVE. #60E 4: F.GW WITH AN EMBEDDED ALGOR-RUNNING DOLIDUE SHOCK PARE FOR SHOCK WAVE 18 ASSUMED TO BE ROTATIONAL. ***********************************	MODE 1. TAROTATIONAL FLOW ALONG RIGHT-RUNNING CHARACTERISTICS.	
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ĸ	1.154	4	1.095	1.006		1.083	1.00.1	1.093	7	1,112	7	1.153	20101	1.157	27.10	1,199		12101	1-107	1.100	1.047	1.099	• 0 7 • 7	120	1.150	1.141	1.152	1-154	1114	1.204	1.2.9	1.250		7 1 1 7	1.125	1.120	1.1.8	1-121	77.0		1.155	1.156	1.178	1-190	1 - 29.5	1.230	1.246	, ,,,
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•	AM5	2225.	27.70	-23750	17812.	.23827	.22725	.23557	.23345	-23898	. 228.52	* 22558	. 22271	.23579	299120	.21374	.21.845	.20671	.20216	.19681	94481:															
		446.242		473.169	675.412	475-814	672.559	468.554	463.476	457-627	451.564	444.621	455.847	431-167	624-167	416×955	409.264	490.551	389-685	575.884	356.965							•								
•	118.74	113		1.670	1.56.1	2.898	20226	2.298	2.525	2,369	292.3	2.191	2.056	1.961	1:649	1.659	1.524	1.500	1:024	•615	022-9															
	A W. W.	5566		12.56.	3626.	3826.	36.35	3455	3865.	5915.	3.94.8	5981.	41.16.	4852.	4.5.6.8.	4126.	* 1.66.	4213.	1276	4.346.	** 50															
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7		3. jg	3.832.0	613	u	5.889.0	ETAF	F1.0060	3 43E	150.942	ISP19	u	187.898	ETAT	#1.0162
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5		ça Va	\$652.1	110		1.6088	5113	1.0060	3 451	190.945	15910	14	187.902	ETAT	=1.0162
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F 40	6, 14	. 075	1.000	4558.	557.	1.564	44 95.	7.000	349.602 329.525	.18455 .17568	4197.	1000.	5000	17.650	3371 - 3832 -
£ 4		ų,	\$432.2	619	h	3809.3	£ 74 F	#1.5060	1SP E	190.950	ESPID	μ	187.911	ETAI	=1.0162
4. P	23 →	 		4531.	652	2173	4574.	7.889	554.702 516.011	.16965	4166.	1000.	5000	20.070	3852.
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2		-	3832.3	£10		\$609.	ETAF	#1-006G	1SP =	190-955	ISPID	u	187.919	ETAI	=1.0162
0. N	FG 44	K. 0	1.000	4758.	616.	1.512	4736.	\$.420	296,765	.16551 .15817	4106.	1000.	50000	17.913	3825.
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		*1.016X	3086	1625.	1910-11	1693.	CTAT RI-\$162	3484.	CTA! #2-6161	5497. 5635.	ETA! #1.0162	35.44	3849.	ETA! =1-5626						
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	-1527E 4	198-968	.1328\$ +1	t	198-965	.14829 41		*14290 SI	199-968	.13758 31 .13235 31	190-972	23262 33	433305 33	193.800						
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	1031	3869.6	\$ 0.23.E	1.565	3869.4	1.56%	31620.0	1.5584	34125	1.653	\$4185	24622	1,642	3918.7 CTAF # -9925						,
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١.	281.	669	.274.	l	1.512	4295.	5.461	85.4	2	4265.	1000-	5000	15,113	3.8
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÷	. 35.4	1.907		-	1.554	5191.	15.654	34.41	322	S	1000.	5000	18.376	3578
_	42.0	1.275	16234		1.555	5201.	15.000	233.036	.13162	3922	0	5000	20.069	3944.
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•	č	4	53.5	72	1.414	4572.	9.073	335.014	.17811	4167.	1000	5000	13.890	2665.
	ŗ.	é	* 3.8	4016	1.447	4661.	10.110	319.650	.17128	1154.	1000	5000.	14.103	2711.
	1	5	6.5.7	417	1.479	4747.	11.142	504.981	.16470	4102.	1000	5000	14.514	2753.
		9		1019	1.512	4631.	12.172	290.951	.15837	*070*	1000	5000	14.505	2793.
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	ž	9	492	55.2	1.569	5228	17.275	229.016	.12975	3911.	1000	2	15.320	2967
	7	-	016	1668.	1.664	5215.	16.117	230,896	.13361	5916.	1000	50005	9	5301.
			67 61 60	38.7	1.666	5222.	15.400	229.837	.13012	3913.	1006.	5000	œ	364A.
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ø.	. 235	.725	07	294.	1.237	40612	4-125	425.465	.2.737	4336.	1000	2000	10.503	200%
·		. 742	4177 .	\$81.	1.277	4194.	5.208	404.075	.20822	.599.	1000	2000	10.966	2099.
	. 25.2	151	4 7	0	1.515	.000	6-219	364.255	.19968	4263.	1000	2000	11.365	2178.
	176		9	643	15.45	. 6 0 3	7.341	365.782	.19161	4228.	1000	5000.	11.716	2248.
٠.					4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4		A	X48.216	1839	41.4	1060	2000	12.033	2313.
١.						4.5.43	0 4 4 5	441.656	1766.2	4150.	1000	5000	12,323	2372.
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COUTINUATION OF THE IRRUTATIONAL FLOWFIELD ALONG RIGHT-RUNNING CHARACTERISTICS

11	-	2528	2575	56192	2661.	2761.	5055	3382	5714.	1066. =	1427.	1562	1704.	1807	1976	2049.	2117	2234	\$295	2348	2448.	2495.	2032	51.62.	200	*060 *	E .9893	1.5.1	131.	***	1000 1655	1741	1521	1963.	202H.	21.64	2205.	2259	
1, 1, 2, 1, 1, 1, 1, 1, 1, 2, 1, 1, 2, 2, 1, 2, 2, 2, 2, 1, 2, 2, 2, 2, 1, 2, 2, 2, 2, 2, 2, 2, 2, 2, 2, 2, 2, 2,	RDOT	15.075	13.295	13-502	13.700	10.071	15.507	17.020	18.548 20.067	ETAI	7.684	8.283	6.907	9 4 4 4	10.269	10.628	10.957	11.545	11.812	12,064	12.531	12.749	14.293	15.718	170	20.067		6.034	5.8EF	7.55.1	3.593	9.020	: [10.409	10.975	11.241	11.4492	
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_			12.7	13.5	154		26.6.93	156.938	11537	5620.	1000	5000	6.467	1528.
			53.43	. 40	1.796	5532	15.784	186.163	000000	5781.	1000	5000	6.777	1386.
	:	4630	233	¢	1.854	22.25	16.773	175.44	10569	5741.		5000	7.078	1455
		. 7.2 %	2436.	1741.	1.873	5736.	17.756	164.459	.09654	5782.	1200.	5000.	7.570	1521.
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-		3 2 4	5568	1365.	1 . N . 2	5751.	13.699	155.467	.09595	56.62	1000	5000	14.176	3062
-		. 176	5 F 2 B .	4		37.86.	13.620	155.163	.09379	3665.	1000	5006-	15.227	3279.
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#7 ##	•		-	\$25	1.352	6405	4.705	203-635	84161*	*25B	1000	000	2 2 3 3	F 4
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			5045	1049.	1.650	5202	11-650	252.777	415150	5922 .	1000	5000	5.032	1020.
•	•		5172	- 2211	9.69.	5550	239.23	015*812	977210	1681	1000	2000	5.572	1095
			5243.	1273.	1.738	5396.	13.647	205.186	11838	2040	1000	2000	5.702	1165
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e -		* **		1 7 4 1		2 4 5 7	17.538	158-772	. 3556	5679	000	\$000	6.926	1437
Q.	A CANADA SAN AND A CANA		5486.	1.54.7	1.201	5 7 24	442.51	165.099	31260	\$705	1000	2000	129.1	1595
51 62	**		5497.	1046.	1.867	5686.	14.756	167.575	. 10000	3713.	1000	5000	6.455	1780-

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HDOT	9.376	10.355	11.369	12.403	13.433	14.404	15.454	0 3 4 4	18.276	19.173	20.060	ETAI	.751	1.303	1,811	2.281	32122	3.530	3.906	•	4.61	5.281	5.59R	5.907	0 . Z . Z	7.149	7.935	8.810	10.724	11.718	12.717	13.708	14.680	15.527	17.433	18.311	7	91	ETAI	*	1.040	1.508	1.950	2.757	3-149
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!	3712.	3705.	10	3680.	3665.	5648.	3631.	2010	3546	1860	3545	ISP	4322.	4265.	4213.	4165.	4118	4030	3986.	3944.	3902.	3819	3778.	3738.	3697	3682	3697	3693.	3686.	3662	3647	3630.	3513.	3596.	3562	3545	3528.	3510.	ISP1				4148.		
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d	167,427	165.569	62	158,995	154.995	150.823	146.618	014541	138-489	1150461	2 2	ISP =	416.955	5.08	J)	333.984	~110	, ,	10	240.891	225-923	198.646	186.194	174.456	163.385	152.95	162,229	162-293	160.626	101 640	150.525	146.486	142.404	138.370	130.674	126,966		2	ISP =	409-264	377.232	350.090	326.105	280.640	266-359
THETA	14.047	13-634	1	13.375	13.427	13.554	13.735	104.401	14.188	֓֡֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓	15.000	= .9845	1.699	6	•	1.	7 117		42	10.457	11.475	12.484	14.478	լտ	16.446	17.421.	14.518	13.790	13-364	13.12	13.13A	13,262	m	'n,	13.840	•	14.701	- 3	- 9853	1.524	2.749	3.961	5.015	7-159	8-219
VMAG	5685	5701	5725	5.7	5789.	ויש	5862.	2677	5935	ها م	6042.	ETAF :	4126.	Iσ	*	ln i	~ 0	4935		5148.	5249.	5547	55.37	5629.	5719.	5807	57.7R	5726.	57.2.	5765	5828	5863.	5899.	5936.	216	6000 6000	o (c	6115.	ETAF	4166.	4338.	4487.	4623.	- DC / 4	4985
Œ	.863	•	1 880	•	1.909	1.925	1.942	1.77	1.976		2.026	4473-7	1.253	1.313	1.367	1.416	9	1 . CO CO CO CO CO CO CO CO CO CO CO CO CO	2	1.636	-67	1.718		1.838	1.877	1.917	1-882	1 8 8 4 4		1.898	1.926	1.942	•	1.976	2 993	2000	Ö	2.060	4497.9	1.267	1.328	1.583	1.433	1 8 4 8 1	1.571
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	4515		55.28	5599	5631.	5663.	5694	5725.	5754.	5/82.	5810. 5837.	اما	١.	4290	4432+	4560°	4576	* C W G 4	4975	5062.	5144.	5221.	5,561.	5425.	5485.	5540.	454 K	5563.	5586.	.:		5707	5738	5768.	5797	10 / D =	A R R C C	5906			333	77	10	723	4932
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	Signa	3-167	4-211			2010		2010			6-247	9.141	10-076	11.634	10.00	13.902	14-824	15.709	16.583	17.443	18.283	20.059	ETAI	345	.790	1.214	1,623	2.016	24395	4 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	3.458	3-792	4-117	4-745	5+049	F 40 40	6.184	6-8-9	7.650	8.4.97	9.388	11.233	12-158	13,069	13.962	14.820
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-	29.07	1926	2067	2637	4144	716	77.7		3668	3671	36720	3667.	3636.	3643	4611	33.90	5577.	3560.	3543	3526.	3509.	3475	ISPID	4295.	2 3 5	4180.	123.	4081.	4034	3944	3900.	38574	3814	3730.	3638.	3645.	3634		3649	3544.	3634.	36.06.	3590	3573.	3556.	3539.
2	13421	13167	11446	T T T	76 77 1 0	11135	44.4		10000	69493	.09469	.09395	+03564	76060°	08710	0850	. 03306	. 08111	-07918	.07729	10 4 G 10	.07163	222-211	.20671	. ~	.15082	e17018	.16048	415153	13536	.12800	12105	1047	.10231	• 09668	.09132	08984	0914	• 09172	• 09100	.06986	08647	.08455	.08258	• 08062	.07872
.	243-315	233-557	216-752	0.4.0.2	175-071	166.039	157.156	146.91	155.757	156.686	156.965	155.493	152.892	149 590 448 890	٠, •	138,030	130.115	150.397	126-632	3.02	C N		ISP =	400-551	368-316	341.138	3170192	an .	257 606	241.083	225.363	210-756	177.055	172.234	0	150.290	147.362	150-563	151.075	U.	147-4:1	пч	· •	in a	•:	125.765
T ME. T.A.	9-235	10:269	11294	120239	36.20.2	14.064	70.21	2.5.6	15.437	16-251	13.00	13.064	12,842	12,775	12.010	13-115	13.529	13,564	13.821	14.091	14.574	15.000	.9860	1.509	2.546	3.699	9,012	5.896	0000	0 0 0 0	10.060	~	12.072	14.049	15.02B	16.000	15.146	13.930	7	12.709	12-477	120000		12.735	12.948	15.182
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>	5026		5196.	25/26	3347	CA & C		45.67		5597	56:3.	5635.	25620	5691.	975	5784	5814.	5545.	5671.	5899.	2425	5960	613	4212.	383	.828	46534	6777.	9727	5003	5172	52562	10000 10000 10000	5.78.	\$543.	3604	2652	\$656	5670.	25910	5717.	2,75	5806.	5857.	5867.	5635
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4	118.566	115.233	111.804	106,249	104.525	= dS1	189.866	457.332	X X A C C X X	300 m	204.796	265.289	247.369	230.802	215.416	201.081	187.691	175-163	163.424	nį •	142-163	1.9.758		144,010	143.003	140.832	137-927	154 - 155	127.391	123,769	120.270	116.812	113-445	106.837	103.411	9.62	95-102	15P =	47K-A04	TA 2. 470	515.303	291.571	270.375	251.158	233.569	217.362	188.419	175.438
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THE 1'S.	12.250	13.262	14.22.5	15.186	16-139		ď.	12.081	11.589	11.53.3	11.245	11.273	11.387		1.	•	12.266	12,539	12,025	15-178				15.000	# 888 #	0.000	1.256	2 . 3 7.1	3.455	4.515	S.	ъ,	S	8.59.0	¥ .576	10.557	11,528	13.490	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	666.	100004		11-077	1 10	10.284	ю	10.213	10.326	10.501	114.03	10.956	2,00
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	986	1:129	00001	76.	10983	2.207	2-346	2.634	2-956	3.382	3.865	4.390	4:946	5.519	b. 102	7.264	7.843	8.422	9.005	9.605	10.248	10.983	11.902	13-157	15-002	1 6 6 0 2	20.056	ETAI :	0.000	•026	560*	194	750	612	.776	.954	1-158	1.525	12728	1.814	2.010	2-294	2.650	3.062	3.518	4.512	5.032	5.557
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	129.225	118.00	169-209	10001	84-642	77.740	86.097	91.043	93.522	94.167	93.601	92.231	90.348	83.156	65.601	2000	78.588	76-226	73.838	71.537	69.088	66.375	63-119	58.925	53.298	197.65	40.275	1SP =	196.875	178.954	163.174	149.010	156-542	114-262	104-639	95.823	67-734	200000000000000000000000000000000000000	67-186	20.00	80.710	83.408	84.355	84.100	83.039	79.545	77.458	76 267
	7.164	8+129		150031	11.00	12-850		9-003	8-022	7.442	7.144	7.044	7-081	7.213	7.410	1.60.1	200.8		8.838	9.149	9.509	9.925	10.448	110164	12.220	13.749	15.000	₹066. 3	0.000	1.061	2.09	3-106	700	6-057	7+017	7.968	8-910		40.5	960-6	7.661	6.631	6.030	\$122	5.636	5.836	6.051	90.
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	2.016	2.065	2-113	2.161	2.256	40.0		2.216	2.203	2.197	2.200	2.208	2.220	2.254	2.249	2 - 26 5	2 0 0 B	2,4,0	22.20		2 . 36 9	2.391	2.418	2.056	2.511		2,661	4812.2	1.754		1.678	1.932	0 m 0 m	2.087	2,137	2-147	2.236	2.286	200	0010	2.263	•	2.258	2,259	2.267	2.291	2.506	
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E HO	12878	11443	.05536	.65115	0.09723	78740	0 47 1 8	.04207	.04525	•04104	98740	32/30	+04102	.04618	.04521	-04417	215	76040	03989	.03879	•05765	63634	* P C P C	1 2000	.02612	.02230	.02101	01610.	.01748	242.225	.07001	.06439	.05928		146.54	.04269	-03932	.03621	• 05532	404149	04341	-04435	.04461	04440	10520	.04225
•	115,729	197.10	82.423	74.55	66.150	# 10 0 P	51.114	59.296	64.703	67.795	69-208	66.871	67.760	66.303	64.640	62-666	000019	57.00 57.00	55.623	55.789	51.868	49.733	4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4	43.046	53.471	29.184	25.766	22.980	20.653	15P E	109.248	96.609	69.469	44.00	66.591	. 60.349	54.677	49.518	44.823	584.574	61.364	65-175	63.618	63-255	X4 U- 17	54.592
THETA	2.006	1	906.4	5-655	6.789	7.11	0 2 2	6.69.0	5-674	5.949	000	2.475	8 · 0 69	3.266	5.533	3.50	9		5-235	5.616	6.024	6.492	6 / D 6 /	2000	10.635	11.952	13.097	103013	15.000	.9877	000-0	-995	1-974	404	4.6.56	5.770	6.696	7.614	8.525	32763	2.592	1.957	1.687	1.794	7.0.2	2.532
VRAG	6218		6535	6633	6729	6523	7000	66.65	67.80	67.54	6714		67.25	6756.	6781.	6800	90.20		6925	6956	6969	7027	.0.0	253	7365	7474.	7569.	7655	7750.	ETAF E	6223.	63.87.	6447	65524 65524	25.13	6848.	.1.69	7031.	1120.	68.88	6829	6604	6797.	6802.	4. 8. 4. c.	6860.
E	20116	1	2.27	20.724	6.576	27.07	2 4 4 4		2.465	2.379	2.368	25.5	2.379	2.391	2-405	2.600		2.470	2.487	2.505	5 2 2 5	2.546	4/00/	5.2.4	2.760	2.452	2.898	2.459	5.015	4921.6	2.11.5	2.170	2.225	2000	Z = Z	2.143	20005	2.550	2.604	3.66	2.432	2.418	2.414	2.424		2.450
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MDOT	3.209	3.593	3.987	4.3AB	161.4	2000	766.3	000	0000	260	٠.	11.41.0	000	∞اد	18.457	6	1 4 4 0		000.0	.013	640	106	5 7 4	978	. 64	409	•	•625	.740	.910	•	n ()	0/901	25.41	2.555	3.001	3.359	3.724	660.♠	264.4		•	6.052	6.936	8-272	10.6296	13.795	15-422	17.000	18.545	20.047
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КНО	1	04032	18650	3830	03729	03625	61660	02343	0.5244	64000	02780	02424	- :	01740	00710	18		242-134	906	05423	!	577	4206	3864	03220	3260	2662	0.5823		.04134	.04172	.04162	.04119	.04054	10000	00000	5 J C	203605	.03508	.03410	.03306	.03189	.03046	. 02860	.02603	.02264	.02010	*0100	01490	0136	.01263
-		56.345		52.979	51.294	49.586	47.74	66.00	45.597	40.28H	26.059	30.591	26.396	23.42		. 0		100	49.107	80.411	72.616	65.601	59.269	53,543	48.337	43.656	. 39.591	52.859	56.264	9	8.7	-	57.807	56.708	55.575		26.370	49.26	47.682	46.078			40.249	37,317	35.330	28.181	24.450	21.462	17.065	15.309	13.997
THETA	- 2	8-B28		5.784	4.173	4.575	5.000 0.000 0.000	•	4	• 1	8.093	6	•	12.262	9	15.000	ŀ	. 9863	00000		1.944	2.896	5.838	4 - 77!	5.695	6.612	7.521	3.666	1.236	603	.365	.384	.567	•852	1.198	5,00	016.	1 × 1 × 1	* 20 S	0 6 4 6	♣.0 A.2	20.00	5.227	6.072	7.284	8.988	ο,		12-555	7	15.000
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	27-225	19.8.4	16.486	10.00	12.5	11.986	169.6		9.015	24.649	24.007	23.116	22.443		160-410		55.585	52.657	9.988	100	43.676	1.023	38.974	56.795	54.287	27.225	22.589	19.044	16.486	12.837	11.485	+09*6	9.015	24.649	24.007	99-400	22.443	47.551	62.165	56.410	52.238	45.670	* 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	41.243	38.994	
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		184		1	.6047	9779	6220) -	6266.	7504.	7524	255	7575		27.43		69250	6976.	7023	7069	7154	1.76	7240.	7268	3.6	7527.	7672	7787.	75.25	8047	8116.	8224.	8260.	7504	7524	1000	7575.	7068.	6819	6912.	6983.	70.3	7145	71 92.	7239.	
	293.2	3.036	5.134	3.263	3-266	0.20.0	3.41.9	· · · · · · · · · · · · · · · · · · ·	3.453	2.853	2.667	2.887	2.90		289.6		2.487	2.517	2.545	20072	2.625	2.651	2.678	2.709	2.101	2.869	2.973	3.038	421.5	3.266	3.325	8.419	3.453	2.853	2.867	2002	2.903	2.572	2.427	2.479	2,571	2.557	2.619	2.648	2.67	
	1000	1385	1443.	1596.	1728	1447	1814.	•	1880.	-199.	-162.	29	, . , o		0	•	-269-	-94.		œ E	0 -	34.	411.	10 •	376.	860.		1285.	644	1728	1849.	INIA.	1880.	-199.	-241-	v 🏲		0	-284	26.	:53-	o ⊳	1 1 6	134.	.05	
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	1.94	1.514	1.759	2.000	20.00	4.6.5			2.673	2.678	25.455	5.627	5.162		0000		151.	.203	4254	100	. 4 3 3 3	1990	.527	595	1991	4982	1.261	1.514		2 2 0	2.479	2.718	2.878	2.878	2.933	2000	3.162	0.090	100	.105	163	612	.351	. 389	. 53	
	5.846	6.056	7.501	7.727	6-142	6.247		•	1 .			• •	4,000	٠.	2000		4.065	9.144	210	97	6.60	6.4.4	4.917	190	1000	2. E. S.	6.384	6.856	- A	8:142	8.549	6.918	4.195	9.196	955	0 70 0	9.992	091.	4.276	4.413	925.	650	10 P C - 4	5.318	5.147	
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d	\$0.813	26.761	21.903	16.053	14.070	12.471	11-150	10.087	21.2.0	0 4 7		:	- '	25.201	ISP =	447.00	62.452	55.861	51.553	48-173	45,302	40.267	37.749	34.943	51.582	22.160	1000	16.369	14.345	12.715		10.087	A5.A.9	ار	, Ki	24.535	÷		å	46.153	44.505	41.862	39.596	15.576	33.662	51.653
THETA	•	6.907	80 4 4 4 4 4 4	7000	11-674	12,497	13-240	13.875	04.	14-14:	-146	•460	•674	0.000	.9824	0	2.599	, .	3.330	3-672	4.016	4.742	5.166	5.692	6.400	1000	0 C	1:4119	~~	12.831	13.559	13.875	14-149	24-1	30.44	04.60	•674	.913	00000	00000	.343	.878	1.394	1.6530	2.740	3.235
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	0	• 1		3.07	5.218	3.241	5.340	2.193		7 1 9 0			• 86 •	A12	47.2	,	2.424	2.485	2.528	2.565	2.598	2,661	2.695	2.737	2.790	2 - 8 - 6	2 3 7 3	3.138	\$ 207	5.271	5.330	3.393		7000	7 2 2 4 C	2.856	2.864	2.872	2 . 8 . 2	2.588	2.610	2.640	2.670	2.693	2.737	2.789
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	28//09	28.48	2.24	1.578	4000	2010	7	5.9%	P. C. C.	777.7		15.025.0			10.419	-4-120	-1.655		-5-222	67.20	22.454	*1. TT5	-1.0054	518	ប្រសារ	= .9824	00000	4.40	-957	1.532	0 1 2 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	5.446	6.55	7.290	8-217	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0			9.563	-4.052	-4.547	*4.217	-3.8.55	14 × 5 × 5 × 5 × 5 × 5 × 5 × 5 × 5 × 5 ×	-2.333	-1.627	
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APPENDIX E OUTPUT FOR SAMPLE CASE NO. 9

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ς,	4	5.16228	10.619	4947.	4860.	.9824	244.042	242-166	9923	
.	-	3.16223	21.267	4947.	4850.	•9824	244.042	242,166	*9923	
٠.	٠,	5.25141	16.379	4945.	4860	.9831	•	242,156	.9930	•
_	٠.	3.25082	16-777	4940.	4853.	9843	243-699	242,300	.9943	.
•	.	\$ 10 mm	16-144	49374	4864	• 9851	243-565	242.366	.9951	
1	0	3.35474	15.462	*934	4865	9869	245.411	242.410	99959	
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APPENDIX F OUTPUT FOR SAMPLE CASE NO. 10 F-1

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.00146 | •00976 | .00337 | .00138 |
| | d | 42.552 | 6.844 | \$7,204 | 6.324 | 5.177 | 53,182 | 5-650
8-255 | 29.714 | 5.683 | 2.815 | 25.508 | 5-110 | | 4.752 | 2.083 | 20-803 | 1.749 | 17,705 | 3.964 | 1.409 | 14.223 | 3.467 | 10.272 | 2.862 | 626 |
| · | THETA | 0.000 | 13-121 | 000+0 | 12-601 | 15.600 | 0.000 | 12.142 | 900.0 | 11.695 | 15.000 | 0.000 | 11.262 | | 10.805 | 15.000 | 0.000 | 10.508
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3.756 | 4.754 | 0,00 | 5.795 | | _ հատ ա <i>-</i> | 5.857 | 3.596 | 164.2 | 9.460 | 5.212 | 3.963
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| | 3 - | 5.988 | \$4.825 | 5+820 | 5.002 | 5.689 | 0.000 | 2.978 | 0.036 | 2.955 | 2.63.5 | 0.000 | 2.925 | . 62 | 2.03.9 | 6.503 | 0 4 | 2.889 | 000 | 2,753 | 7.295 | 00 | 2.706 | 40 | AL AL | 6.1.2 |
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APPENDIX G.

OUTPUT FOR SAMPLE CASE NO. 11

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| 14.20 1.178 1184 1.185 1.176 1.185 | - i • | .545 | 1+146 | 5119 | 1372, | 1,698 | 5300* | 15,000 | 218.612 | .12480 | 5831. | 1000 | 5000 | 20+069 | 404 |
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| 1.202 1.203 5255, 1408, 1.757 5441, 15440 1994031 11154 3821, 1000, 5000, 0.007 1.008 1.009 11157 1.009 11157 1.009 11157 1.009 11157 1.009 | | +664 | | 5184. | 1389. | | 5367. | NO. | 7 | .12027 | 3852. | 1000. | 5000. | 20.069 | 4091. |
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| 62. | 4526. | 572 | 4634 | ¥712.
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|------------------|------------------|------------------|---|--|----------------|---------------------------------------|--|--|-----------------|----------------|--|--|--|---------------------------------------|--|--|---|--|--|
| -275 | 20.063 | 20.069 | 20.069 | 20-069 | 20.069 | 20.069
0.238 | 20,069 | 20.069 | 20.069 | 20.069 | 20.069 | 20.089 | 20.069 | 26.069 | 20.069 | 20.05 | •258 | 20.069 | •258 |
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0-660 | 15.680 | 15.000
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| 2.988 | 2.191 | 3.200 | 2.556 | 2.525
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| 0 | 1651. | 16.7.
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| 0.00.0 | 1.567 | 1.624 | 1,123 | 1.875 | 2.022 | 2-175 | 2.530 | 0.000 | 2.670 | 2.85 | 5.847 | 5.539 | 5.441 | 9.600 | 5.195 | S.868 | 0440 | 5.437 | 8.880 |
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6. 514 | 7.117 | 2.546
A.150 | 241.6 | 11.5.8 | 13.721 | 16.937 | 10.505 | 5.552 | 23.463 | 6.918 | 7.653 | 8.559 | 35.474 | 33.030 | 10.431 | 294.31 | 40.514 | 196.01 | |
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APPENDIX H OUTPUT FOR SAMPLE CASE NO. 13

| WOZZLE PERFORMANCE PREGICTION PREGERAR. |
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| FHIS PROGRAM WILL AMALYZE THE FLOWFIELD AND PERFORMANCE OF PROPULSIVE MOZŽLES FOR SEVERAL OPTIONS. |
| |
| NODE 2. IRROTATIONAL FLOW ALONG LEFT-RUNNING CHARACTERISTICS. |
| NOJE 3. FLOW 417H AM EMBEDDED RIGHT-RUNKING OBLIGUE SHOCK WAVE. |
| MORE 4. FLOW IN A SCARFED HOZZLE EKTEMSION. |
| THE PROGRAM WILL AMALTZE THE PERFORMANCE OF A COMPRESSED PROPULSIVE HOZZLE. ICHP = 1. |
| THE BROGRAM CAN OETECT AND TRACK AN EMBEDDED RIGHT-RUNNING OBLIQUE SHOCK WAVE (RODE = 5 OR 5). THE FLOWFIELD AHEAD OF
THE SHOCK WAYE IS ASSUMED TO BE IRROTATIONAL, AND THE FLOWFIELD DOWNSTREAM OF THE SHOCK WAVE IS ASSUMED TO BE ROTATIONAL. |
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| PROJECT SPECIFICATIONS - |
| AMALTSIS OF A SCARFED MOZZLE WITH AN ATTACHED MIGHT-RUNNING DELIGUE SHOCK WAYT. |
| THE AMALYSIS IS PERFORMED IN CE UNITS CLNF - LBM - IM - FT/SEC+ R) - |
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| ROPERTIES ALONG THE WOZZLE WALL CONTOUR. X | | 1.0150 | 190.981 | 187.978 | 1.0058 | 3833. | 3P 11. | 205-593 | - 1 | . 00292 | 28 | |
| ROPERTIES ALONG THE WOZZLE WALL CONTOUR. 1 | | 1.0160 | 190-975 | 187.962 | 1.0059 | 3833 | 3410. | 224-568 | 1 | .00259
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